



Cosmology from CMB primaries (and lensing) from space and ground-based experiments

Twitter/X: #neucosmos



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22/07/2025



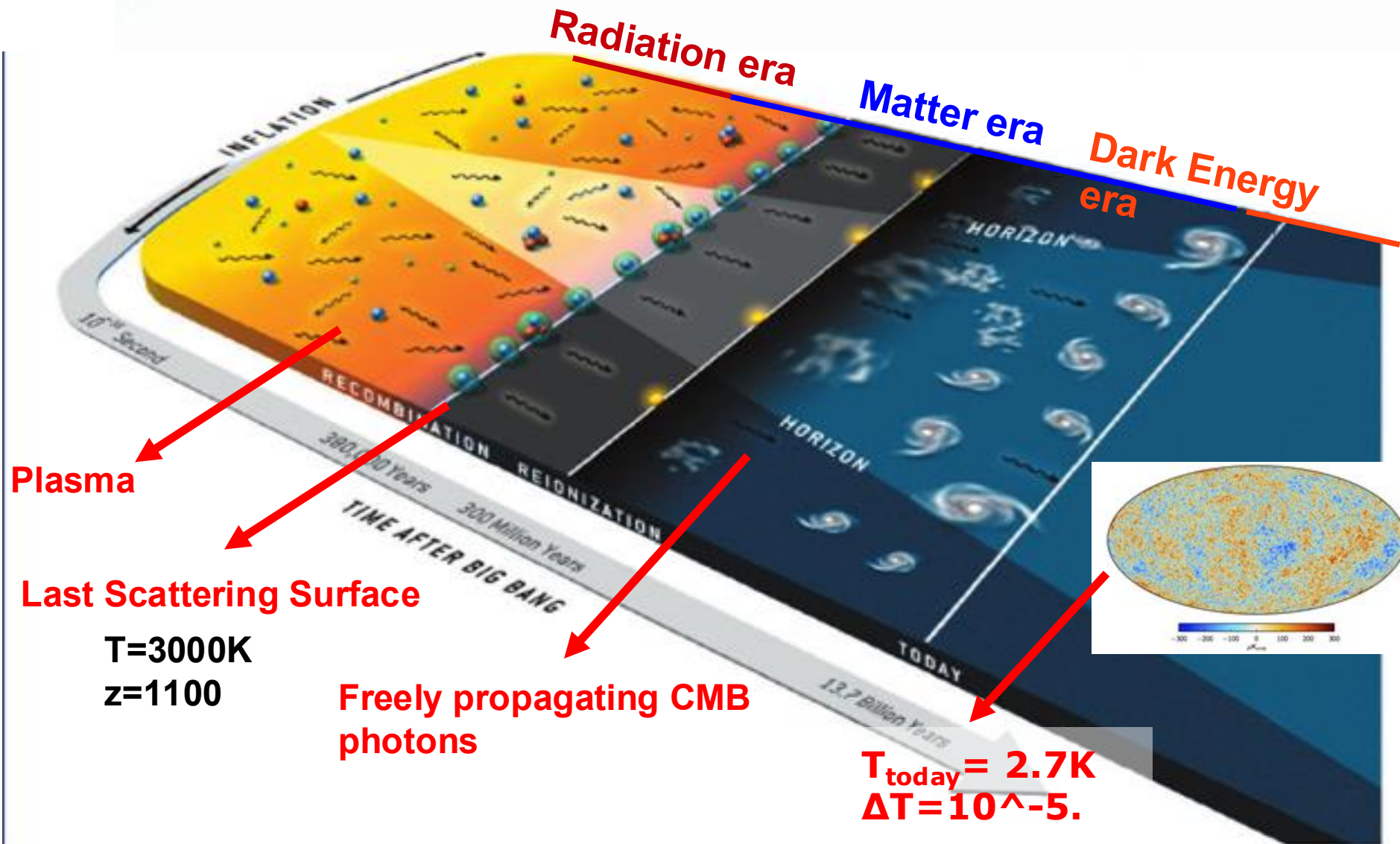
Overview

1. A short introduction
2. The CMB sky and observations
 1. Monopole, dipole, anisotropies, polarization
 2. Ground, balloon and space experiments
3. From observations to cosmological parameters
 1. The CMB angular power spectrum and estimator
 2. Bayes theorem and likelihood
4. Latest results on cosmology (very Planck and SPT-3G oriented)

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CMB



CMB

See Julien's lectures!

Physics of recombination

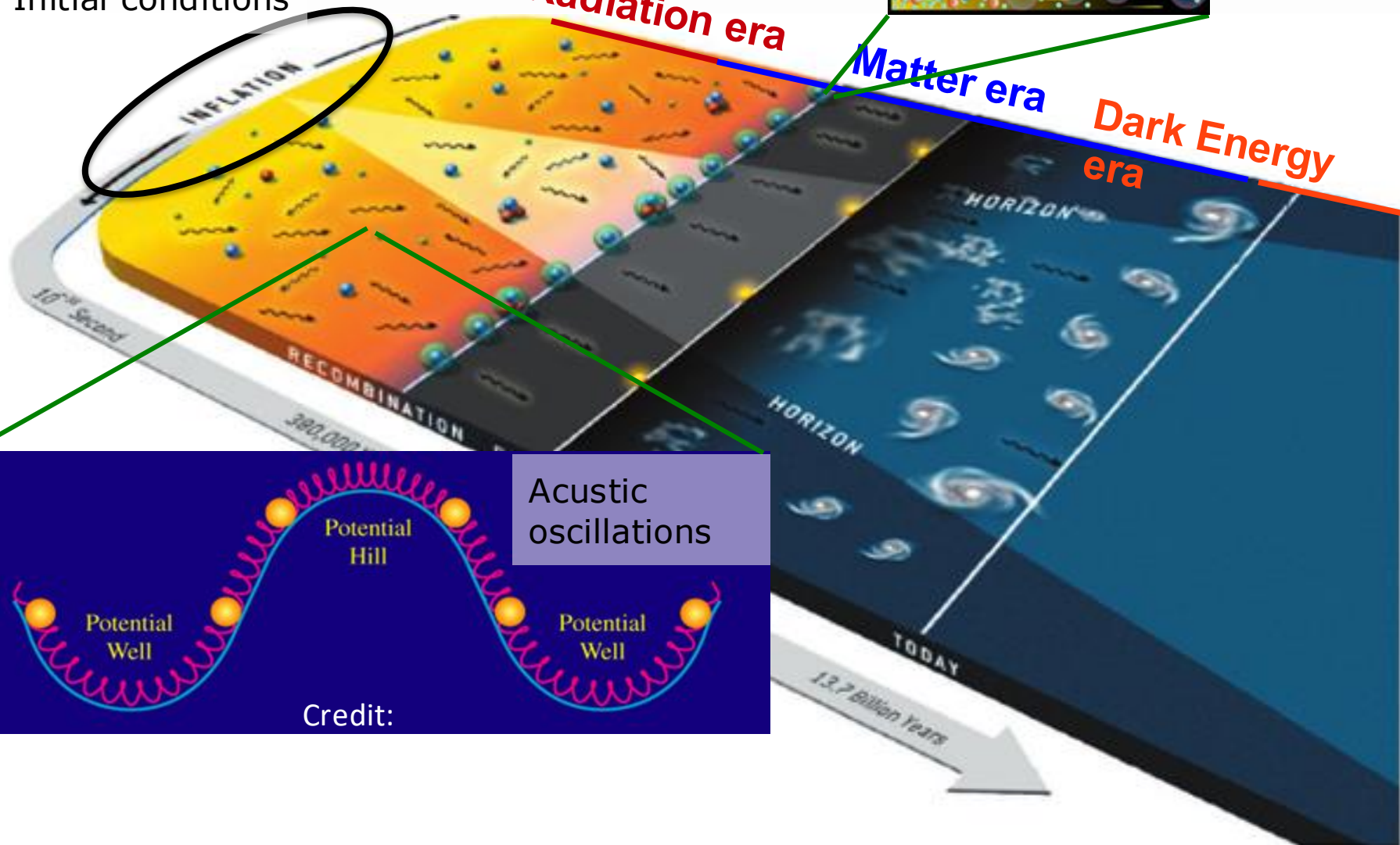


Initial conditions

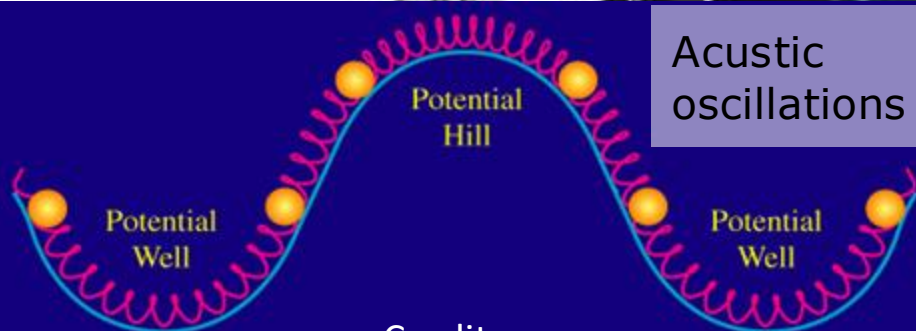
Radiation era

Matter era

Dark Energy era

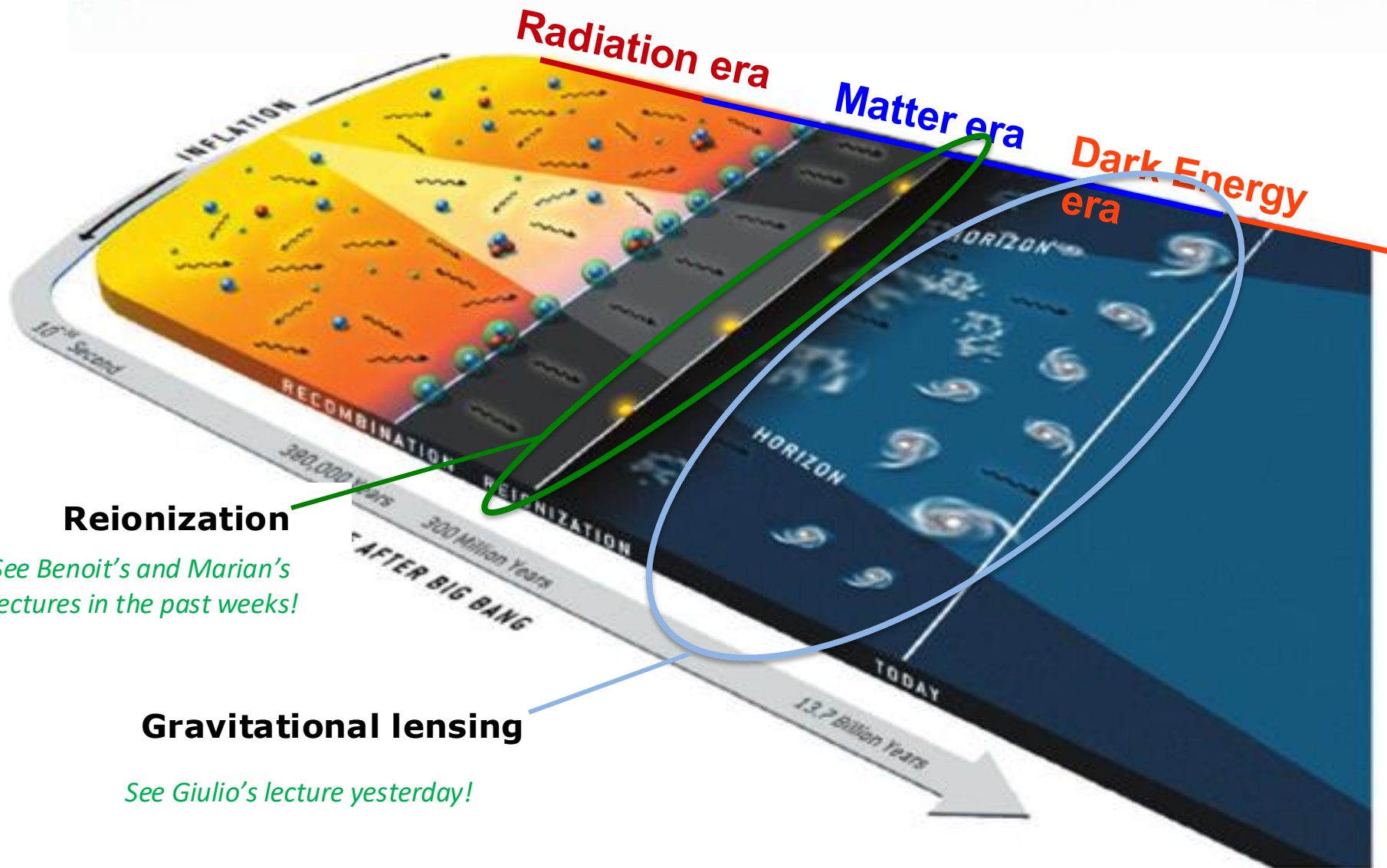


Acoustic oscillations



Credit:

CMB

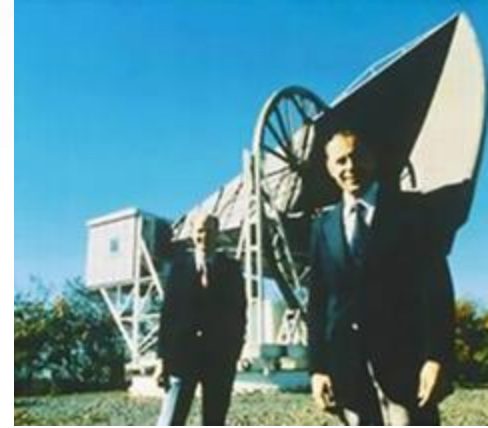


See Benoit's and Marian's lectures in the past weeks!

See Giulio's lecture yesterday!

The Discovery of the CMB

- A story about the importance of theoretical predictions to interpret the data, and of communication in science!
- **1965: Arno Penzias and Robert Wilson** (Nobel in 1978), radio astronomers at Bell Labs in Crawford, New Jersey. Microwave horn radiometer used for **telecommunications** (through balloons).
- Uniform, **unexpected source of noise**. Cleaned birds nests (!) before concluding its cosmological origin.
- Princeton group 60 km away (**Jim Peebles, Robert Dicke, Peter Roll, and David Wilkinson**) working on CMB prediction and detection. J. Peebles (Nobel in 2019) had **unpublished pre-print** about existence of the CMB. A friend, Bernard F. Burke, prof. at MIT, saw the pre-print and told Penzias about it.
- Princeton group confirmed Penzias and Wilson discovery of CMB and published at same time.
- Previous detections in other works, but missed discovery due to missing theoretical interpretation (Andrew McKellar 1940 interpreting obs. from W. Adams 1941; Denisse, Lequeux, Le Roux 1957, Le Roux PhD thesis 1957).



A MEASUREMENT OF EXCESS ANTENNA TEMPERATURE
AT 4080 Mc/s

Astrophysical Journal,
vol. 142, p.419-421

A. A. PENZIAS
R. W. WILSON

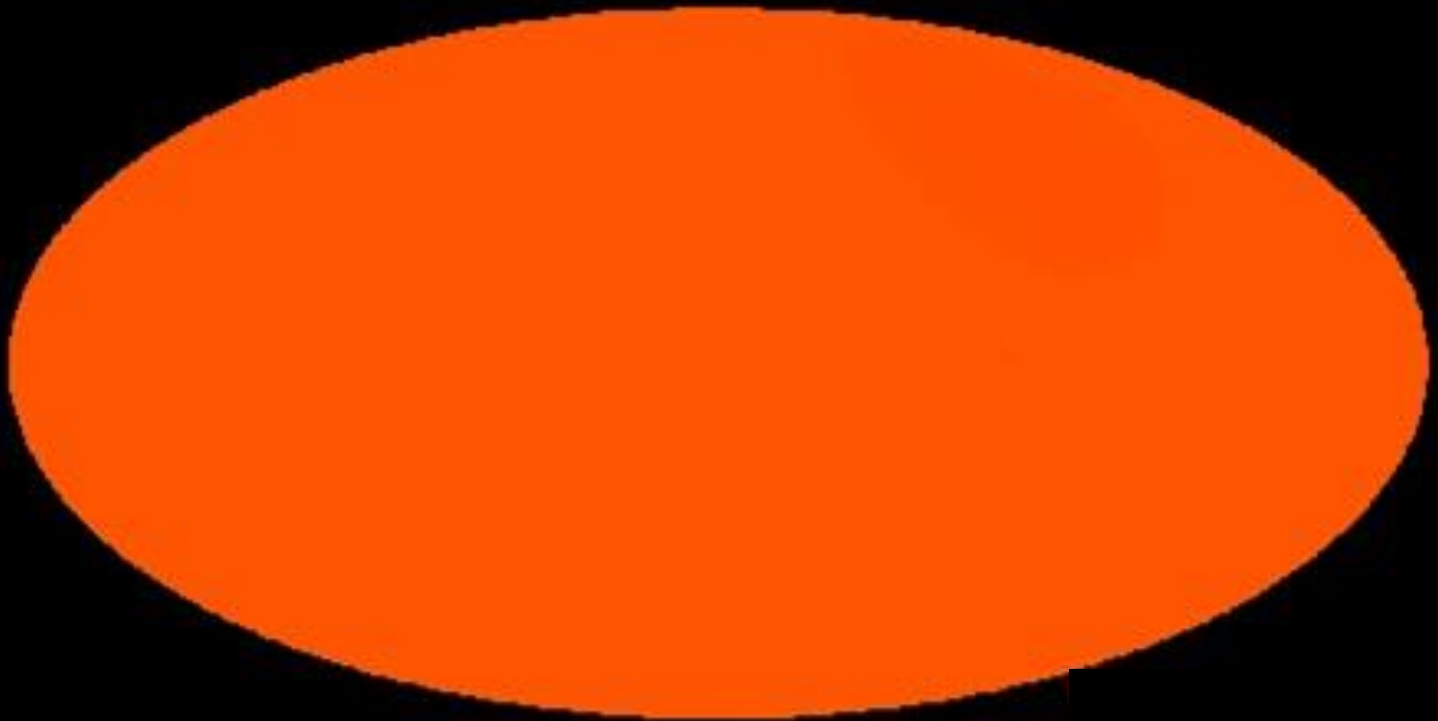
COSMIC BLACK-BODY RADIATION*

Astrophysical Journal,
vol. 142, p.414-419

R. H. DICKE
P. J. E. PEEBLES
P. G. ROLL
D. T. WILKINSON

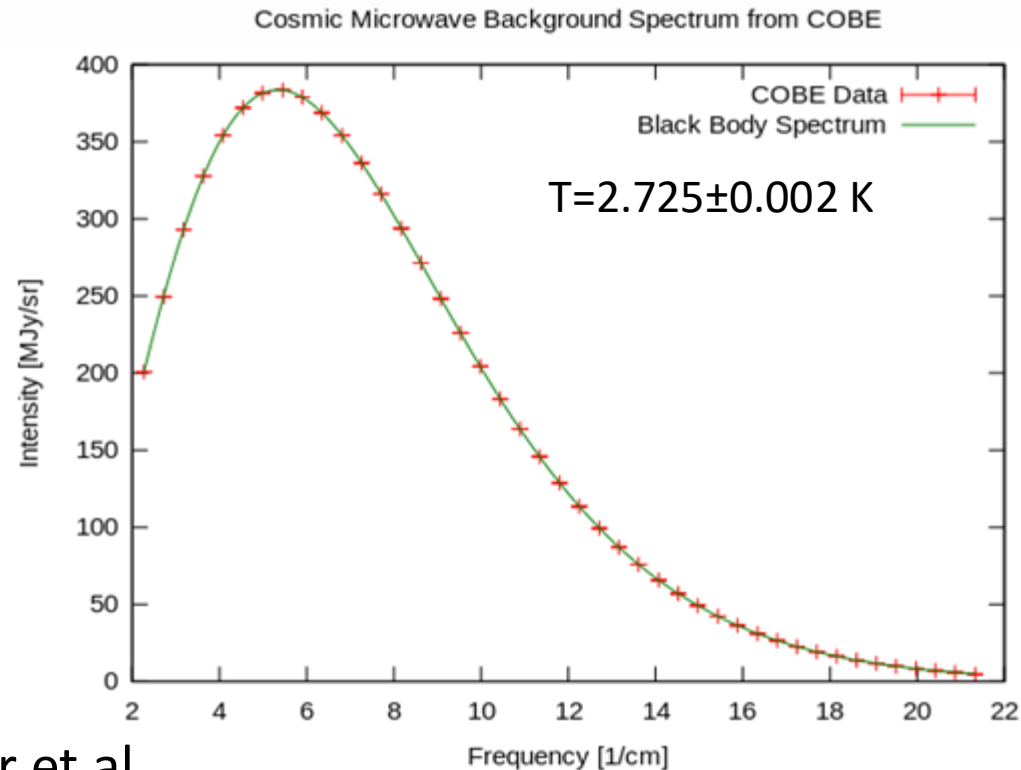
The monopole

- The CMB has a black body spectrum with average temperature of $T=2.725\pm0.002$ K (COBE)



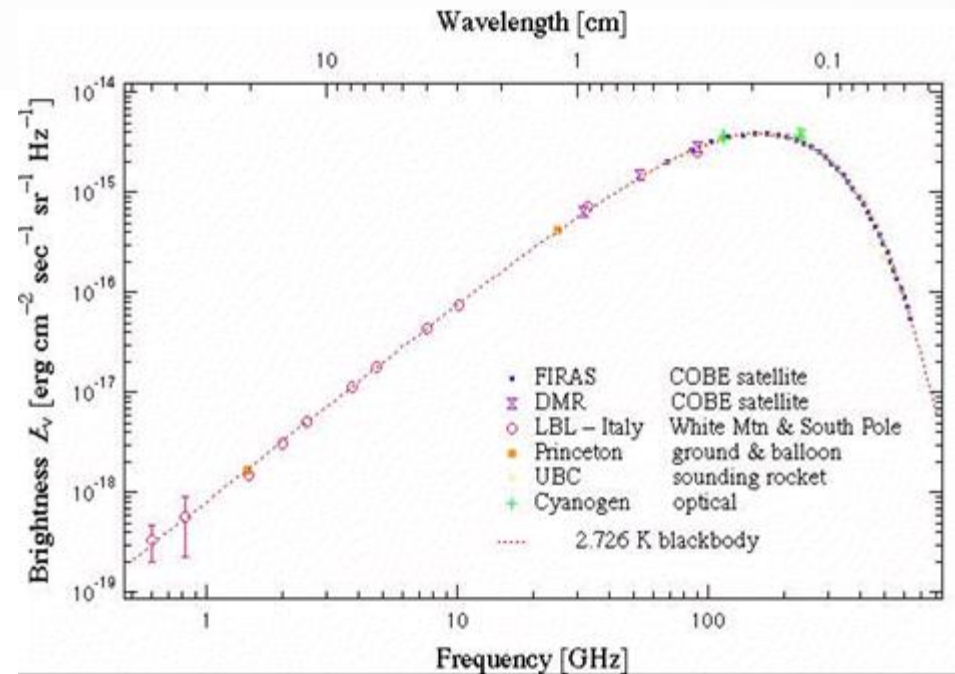
The most accurate measurement to date: COBE

- Launched in **1989**.
- Three instruments:
 - **FIRAS** (BB spectrum) [60-2880GHz], 1yr
 - **DMR** (anisotropies) [31.5, 53, 90GHz], 4yr
 - **DIRBE** (CIB) [infrared]
- FIRAS measurements. Mather et al. (1994, 1996), Fixten 1996
- Peak BB(ν) at ~ 159 GHz.
- Nobel to John Mather (PI of FIRAS) and George Smoot (PI of DMR) in 2006



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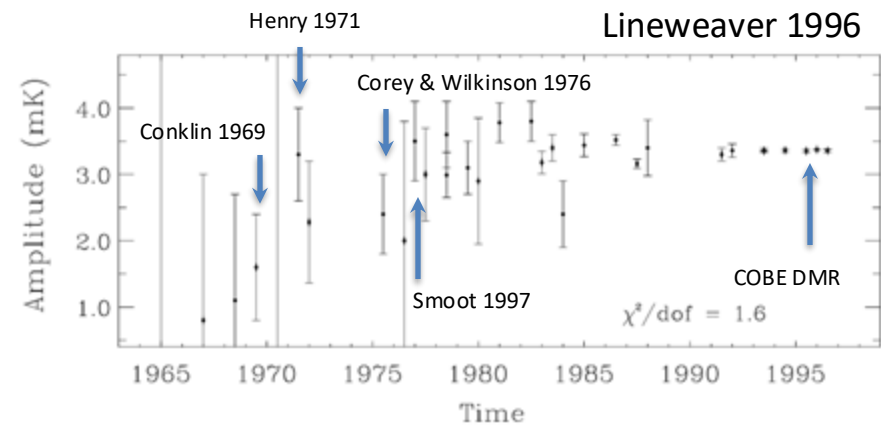
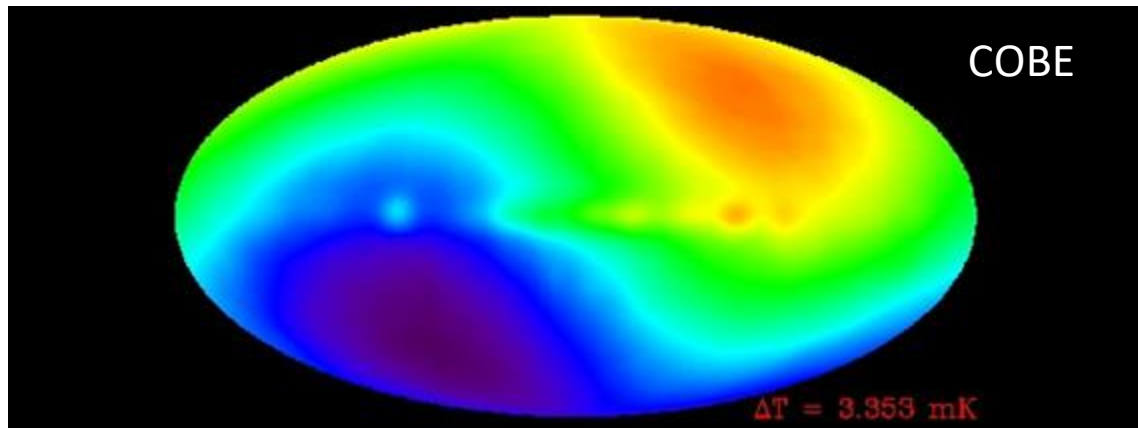


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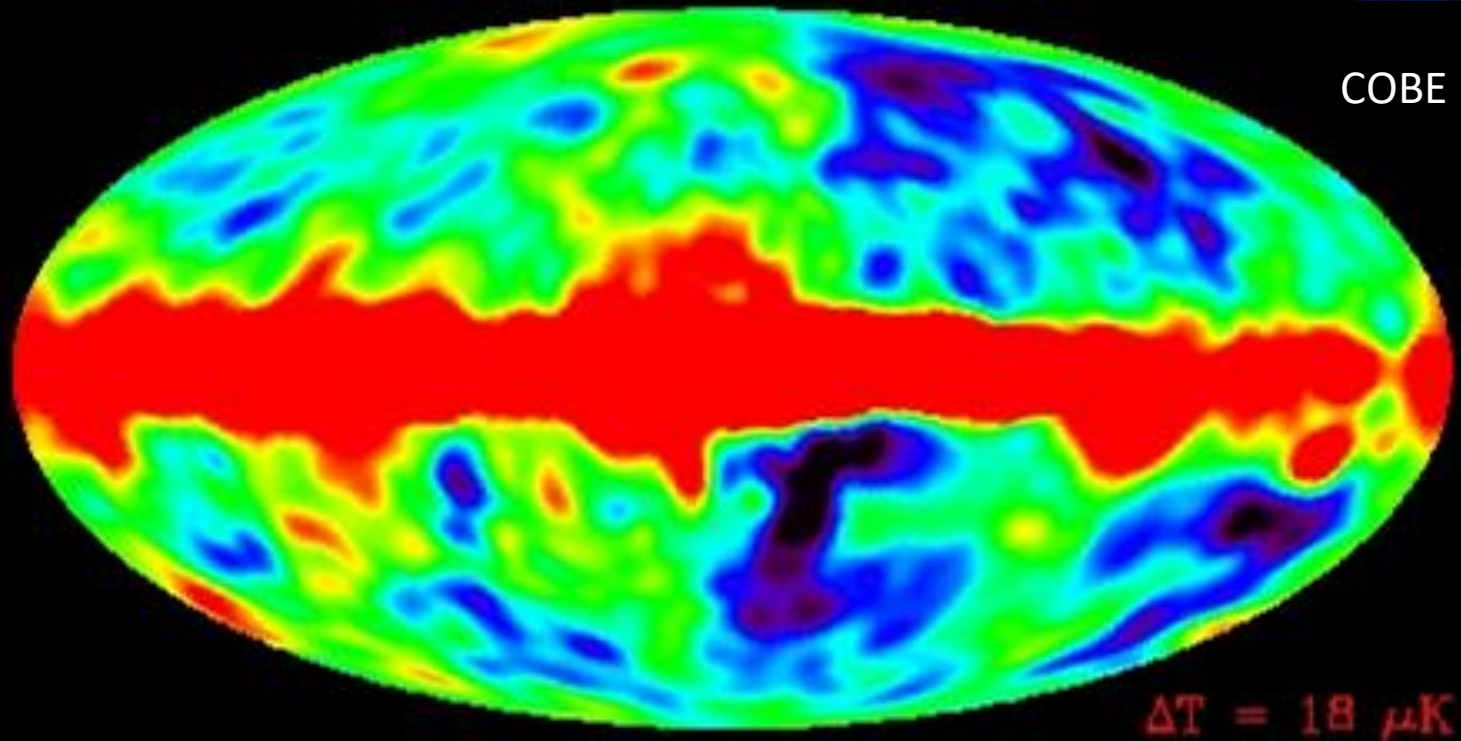
The dipole

- The motion of the sun w.r.t. the CMB reference system produces a dipole of $\Delta T = 3.36208 \pm 0.00099$ mK (Planck 2018) (1000 times smaller than monopole)
- Corresponds to $v = 369.82 \pm 0.11$ km/s.
- Detections shortly after discovery of CMB.
- Velocity of Earth around sun 10 times smaller ~ 30 km/s

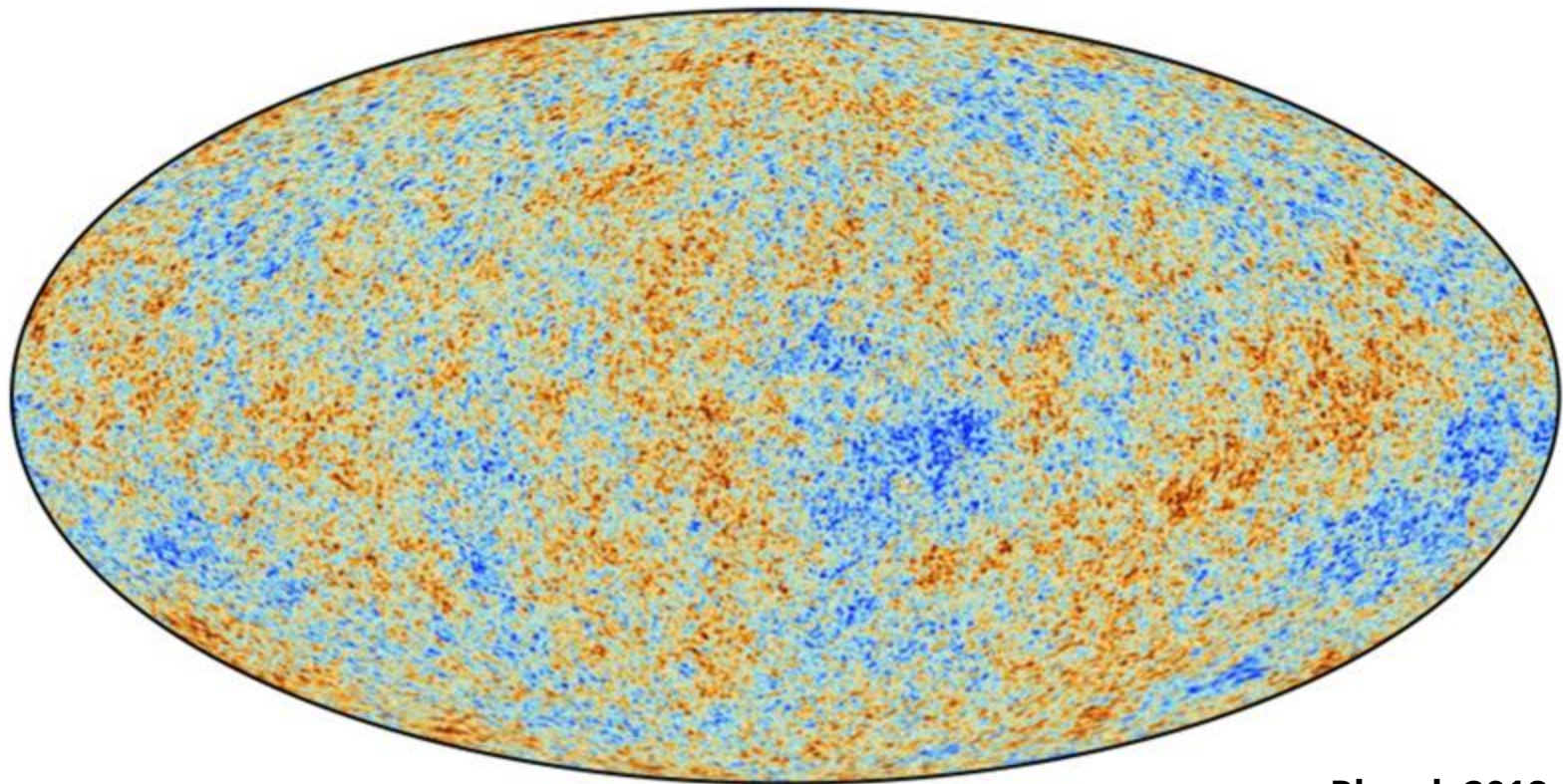


Anisotropies

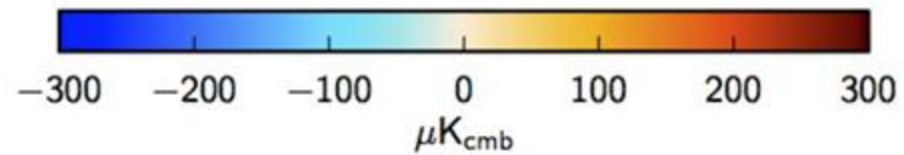
- At the μK level, CMB anisotropies (and foregrounds)!
- First detected by COBE DMR in 1992.



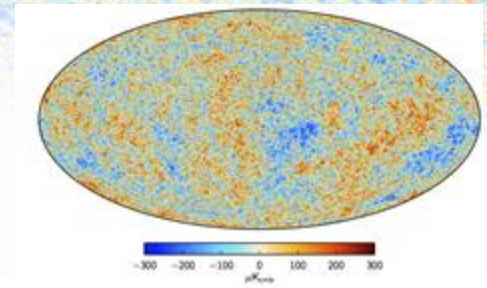
Anisotropies



Planck 2018



A note about units

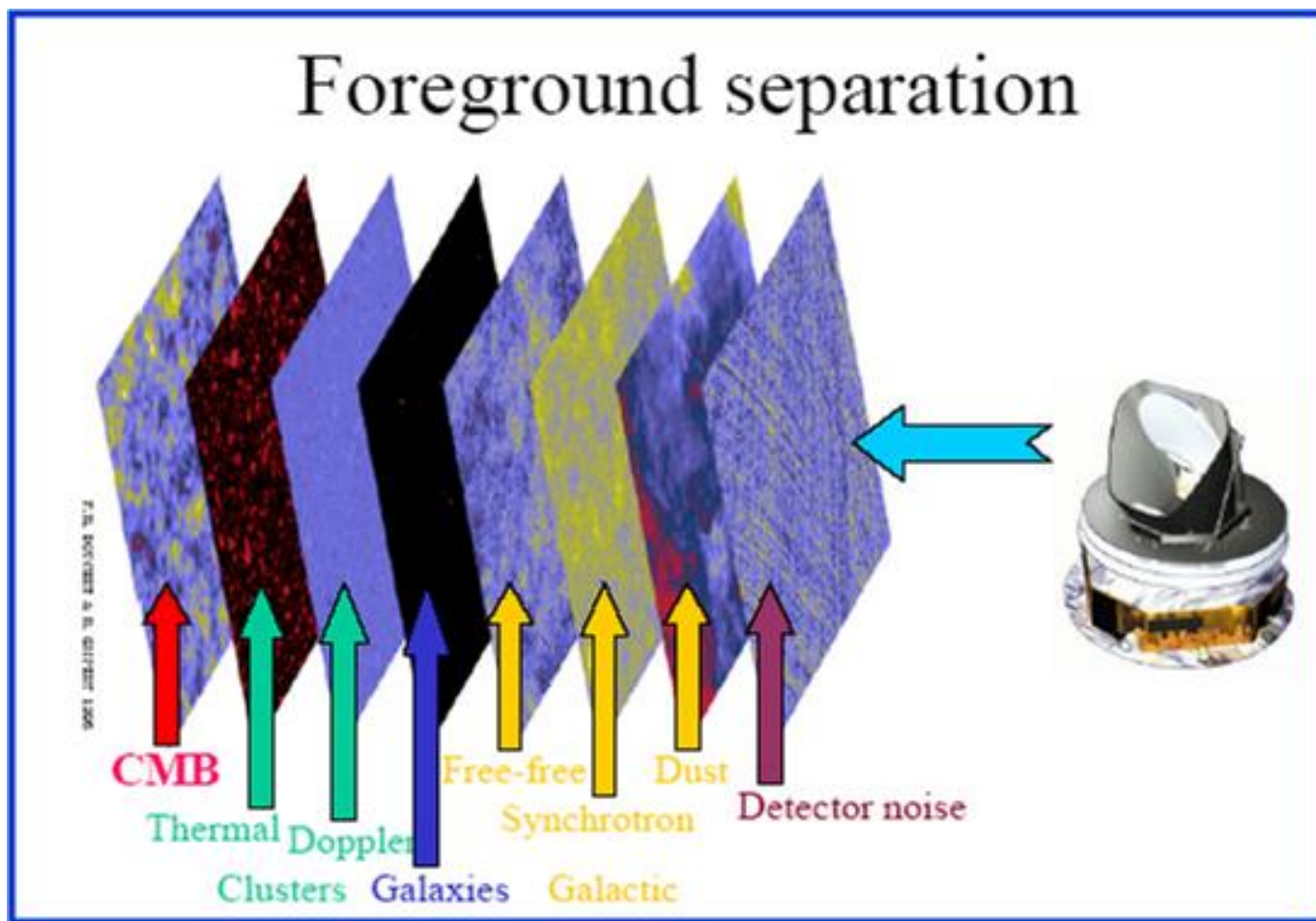


- These maps are in units of thermodynamic temperature.
- Brightness to **thermodynamic temperature** K_{cmb} assuming a black body spectrum:

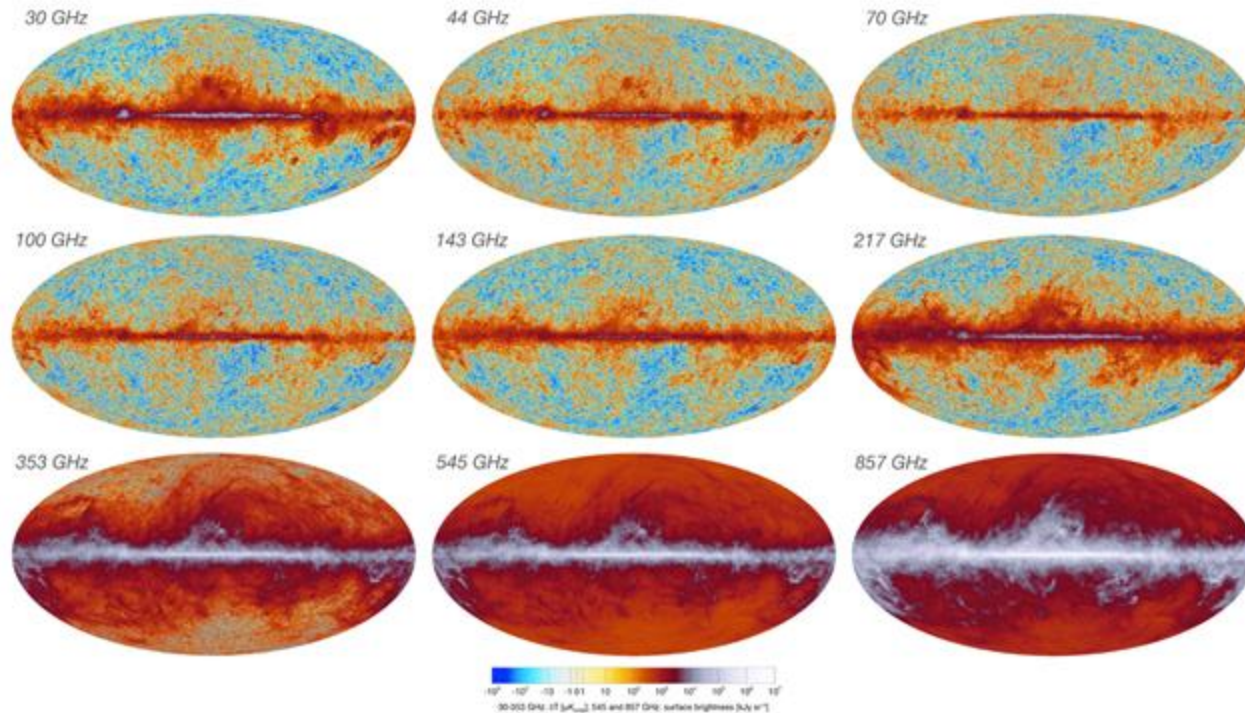
$$BB_{\nu}(T) = \frac{dE}{d\nu d\Omega dA dt} = \frac{2h\nu^3}{c^2} \frac{1}{\exp(h\nu/kT) - 1}$$

- In these units, the CMB has the same temperature at all frequencies, while foregrounds with different emission spectra have different thermodynamic temperatures at different frequencies

Not only CMB...

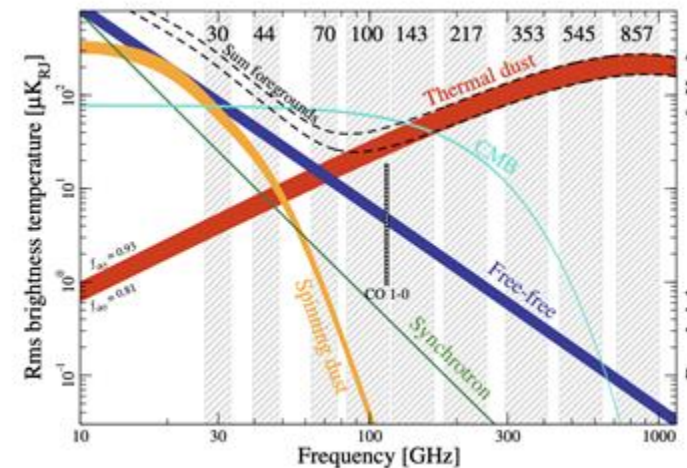


CMB maps at different frequencies



Planck collaboration I. 2019

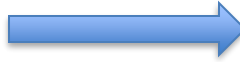
*At 545 and 857 GHz, CMB is weak, calibrated using planets rather than the orbital dipole, units not in K_{cmb} but in surface brightness kJy sr^{-1}



Intensity and polarization in Stokes parameters

For a monochromatic plane wave:

$$E_x = a_x \cos[\omega_0 t - \theta_x(t)] \quad E_y = a_y \cos[\omega_0 t - \theta_y(t)]$$

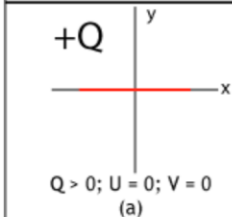
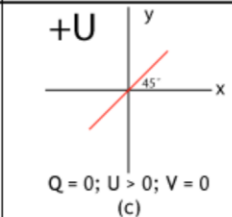
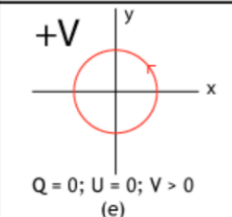
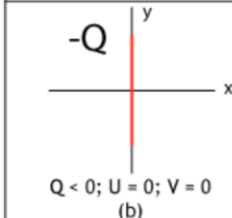
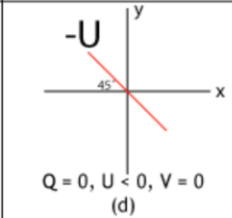
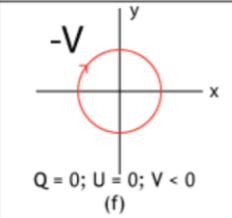


$$\phi = \theta_x - \theta_y$$

$\phi = 0 \Rightarrow$ linear polarization

$\phi = \pi/2, a_x = a_y \Rightarrow$ circular polarization

$$\begin{aligned} I &\equiv \langle a_x^2 \rangle + \langle a_y^2 \rangle \\ Q &\equiv \langle a_x^2 \rangle - \langle a_y^2 \rangle \\ U &\equiv \langle 2a_x a_y \cos(\theta_x - \theta_y) \rangle \\ V &\equiv \langle 2a_x a_y \sin(\theta_x - \theta_y) \rangle \end{aligned}$$

| 100% Q | 100% U | 100% V |
|---|---|---|
|  <p>$Q > 0; U = 0; V = 0$ (a)</p> |  <p>$Q = 0; U > 0; V = 0$ (c)</p> |  <p>$Q = 0; U = 0; V > 0$ (e)</p> |
|  <p>$Q < 0; U = 0; V = 0$ (b)</p> |  <p>$Q = 0; U < 0; V = 0$ (d)</p> |  <p>$Q = 0; U = 0; V < 0$ (f)</p> |

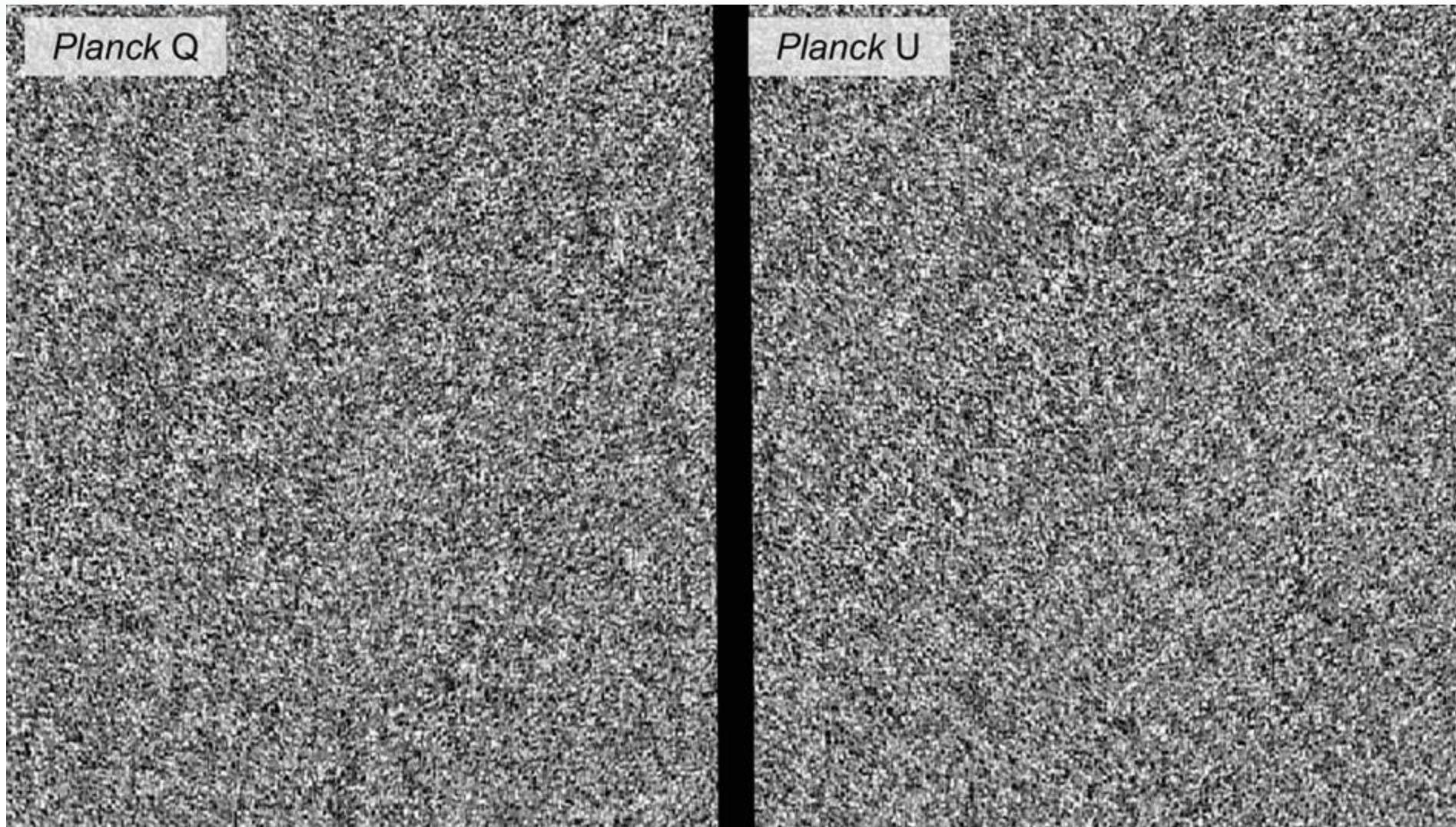
$$I^2 \geq Q^2 + U^2 + V^2$$



- Equal holds for a monochromatic wave or, for superposition of many waves, entirely polarized radiation.
- P = degree of polarization
- In the following, we'll drop V (not produced in standard cosmology model.)

$$p \equiv \frac{\sqrt{Q^2 + U^2 + V^2}}{I}$$

CMB Polarization maps



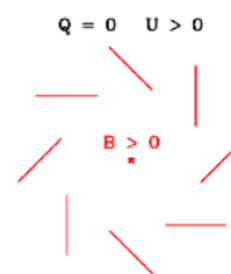
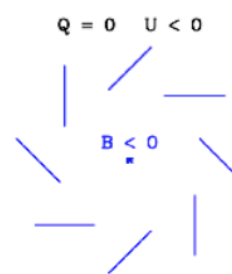
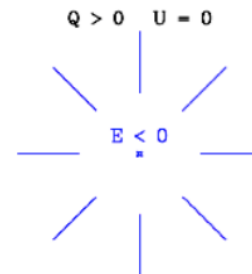
Credit: W. Quan

From Q and U to E and B modes

- Polarization is a headless vector, equal to itself after a 180deg rotation=>Q and U spin 2 fields.

$$Q' \pm iU' = e^{\mp 2i\theta} [Q \pm iU]$$

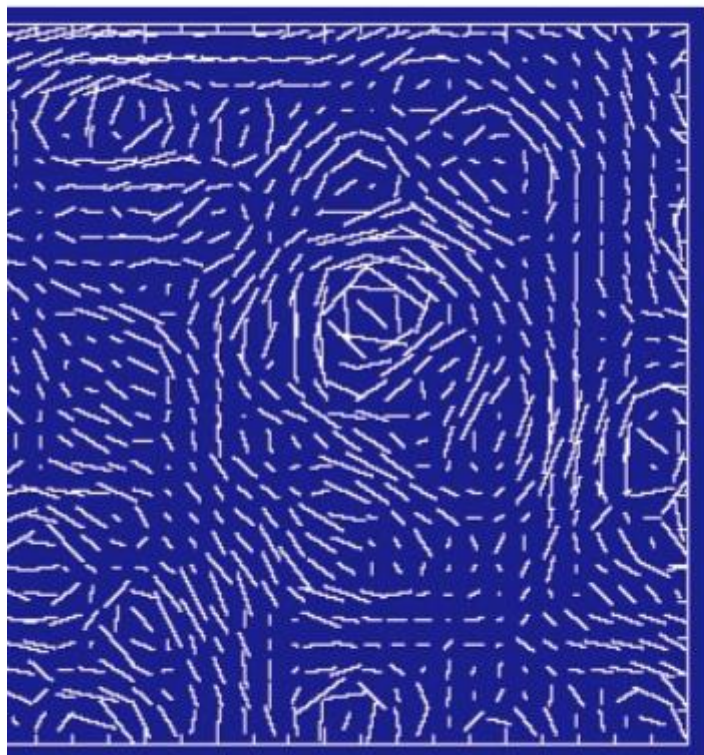
- So Q and U depend on the reference system=> bad to characterize the underlying physics!
- A solution is to characterize polarization not by the characteristics in a point, but a **non-local** average 'pattern' around a point the sky.



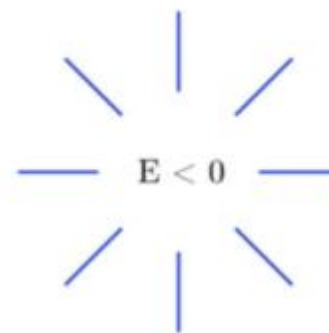
Flat sky approximation:

$$\begin{aligned} E(\theta) &= \int d^2\tilde{\theta}\omega(\tilde{\theta})Q_r(\theta + \tilde{\theta}) \\ &= \int d^2\tilde{\theta}\omega(\tilde{\theta})[Q(\theta + \tilde{\theta})\cos(2\tilde{\psi}) - U(\theta + \tilde{\theta})\sin(2\tilde{\psi})] \\ B(\theta) &= \int d^2\tilde{\theta}\omega(\tilde{\theta})U_r(\theta + \tilde{\theta}) \\ &= \int d^2\tilde{\theta}\omega(\tilde{\theta})[Q(\theta + \tilde{\theta})\sin(2\tilde{\psi}) + U(\theta + \tilde{\theta})\cos(2\tilde{\psi})] \end{aligned}$$

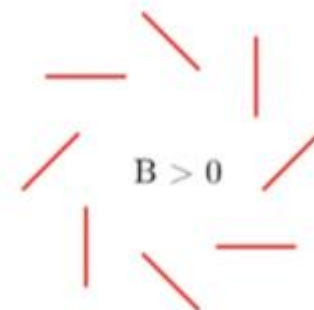
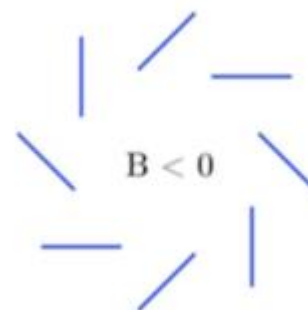
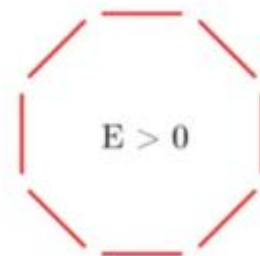
See more details in back-up slides and in Zaldarriaga astro-ph/0106174



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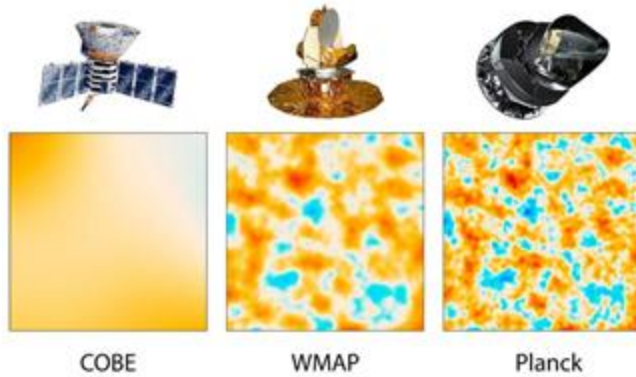


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CMB experiments (very incomplete list)

Satellites



- **COBE** (DMR Smoot et al. **1992**; Bennett et al. 1996)
- **WMAP** (Wilkinson Microwave Anisotropy Probe: Bennett et al. **2003, 2013**)
- **Planck** (Planck I **2013, 2015, 2018**)

Balloons

Boomerang and **Maxima**: first peak. **2001**
Many many others



South
Pole

Ground

DASI: second and third peak 2001, CMB polarization 2002

Many many many others!

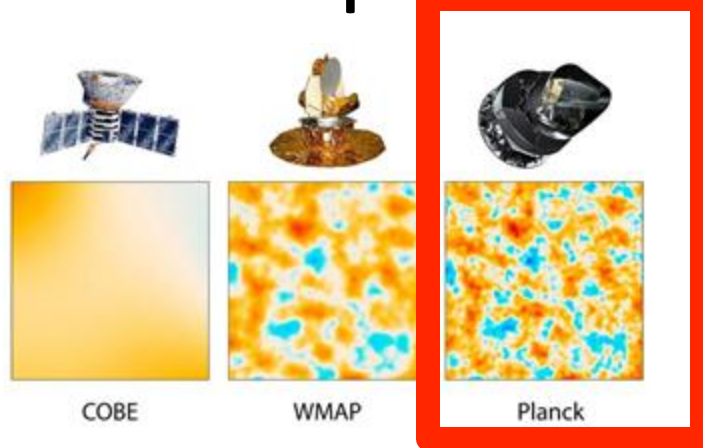
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Atacama
desert
Chile

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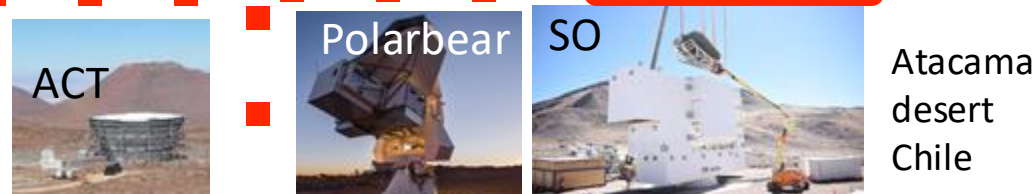
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Ground, (Ballons), Satellites

Satellites

Advantages:

- Full sky observations;
- no atmosphere, large frequency span;
- in L2 Earth and sun aligned and opposite to obs direction so low contamination; stable temperature

Disadvantages:

- Slow to deploy
- small dish, limited resolution (Planck 1.5m, 7' resolution at 150Ghz);
- limited life span (determined by cooling gas/positioning fuel);
- harsh space environment, e.g. cosmic rays; instruments must survive the launch;
- inaccessible for reparations/after launch problems.

Ground

Disadvantages:

- partial sky observations;
- atmosphere limits observation windows;
- ground-pick up.

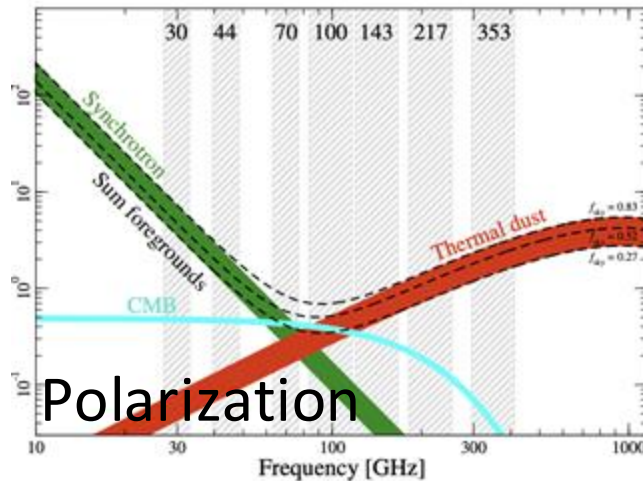
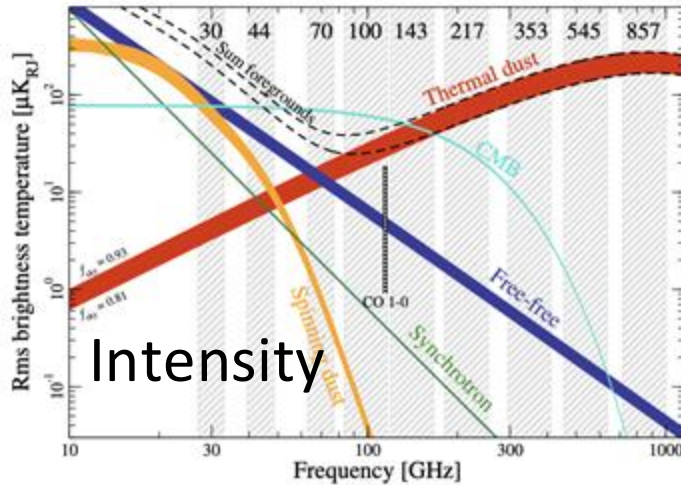
Advantages:

- Faster to deploy
- bigger dish, high resolution (e.g. SPT 10m, ~1' resolution);
- unlimited life span;
- much less cosmic rays;
- accessible for reparations/calibration.

Balloons: partial sky, less atmosphere, less ground pick up, medium to deploy, small dish, short lived, inaccessible for reparation, but payload can be recovered

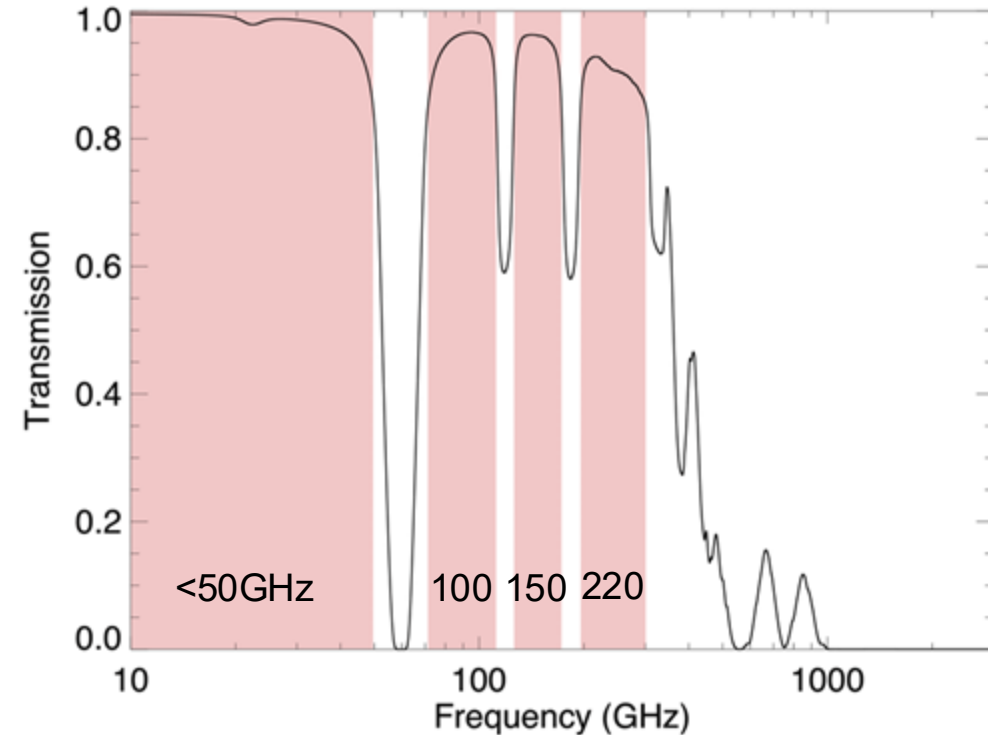
Frequency coverage

Foregrounds vs CMB



Component separation requires many frequency bands

Atmospheric transmission windows



Observations from dry, elevated locations such as South Pole and Atacama desert or balloons or satellite

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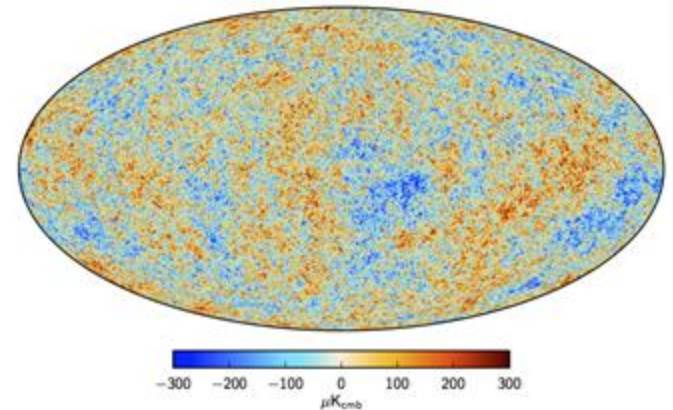
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The angular power spectrum

- The anisotropies are distributed as a **gaussian** random field, so all information is contained in its mean and variance.
- Variance is two point correlation function in real space. Physics correlates temperature in different directions of the sky at $\sim 1^\circ$. Universe is isotropic, the correlation depends only on angular separation, not on the orientation.
- Physical processes put a band limit (limit to the small scale power of the CMB) so useful to decompose it into a complete set of harmonic coefficients.
- Two point correlation in harmonic space is **angular power spectrum**. Isotropy makes each multipole independent from each other.

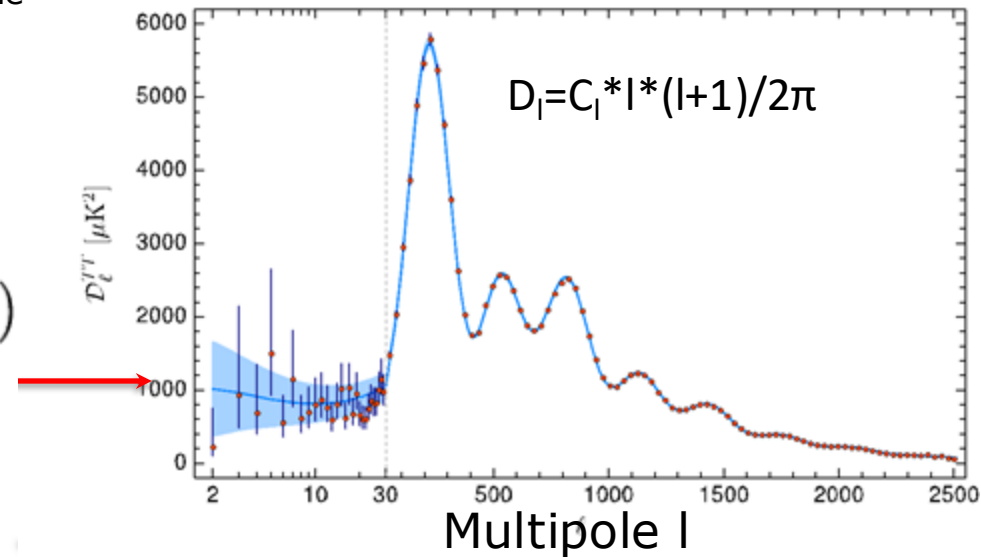


$$\Theta(\vec{x}, \hat{p}, \eta) \equiv \frac{T(\vec{x}, \hat{p}, \eta) - \bar{T}}{\bar{T}}$$

$$\Theta(\vec{x}, \hat{p}, \eta) = \sum_{l=1}^{\infty} \sum_{m=-l}^l a_{lm}(\vec{x}, \eta) Y_{lm}(\hat{p})$$

$$\langle a_{lm} a_{l'm'}^* \rangle = \delta_{ll'} \delta_{mm'} C_l$$

Angular power spectrum



Large scales

Small scales²⁸

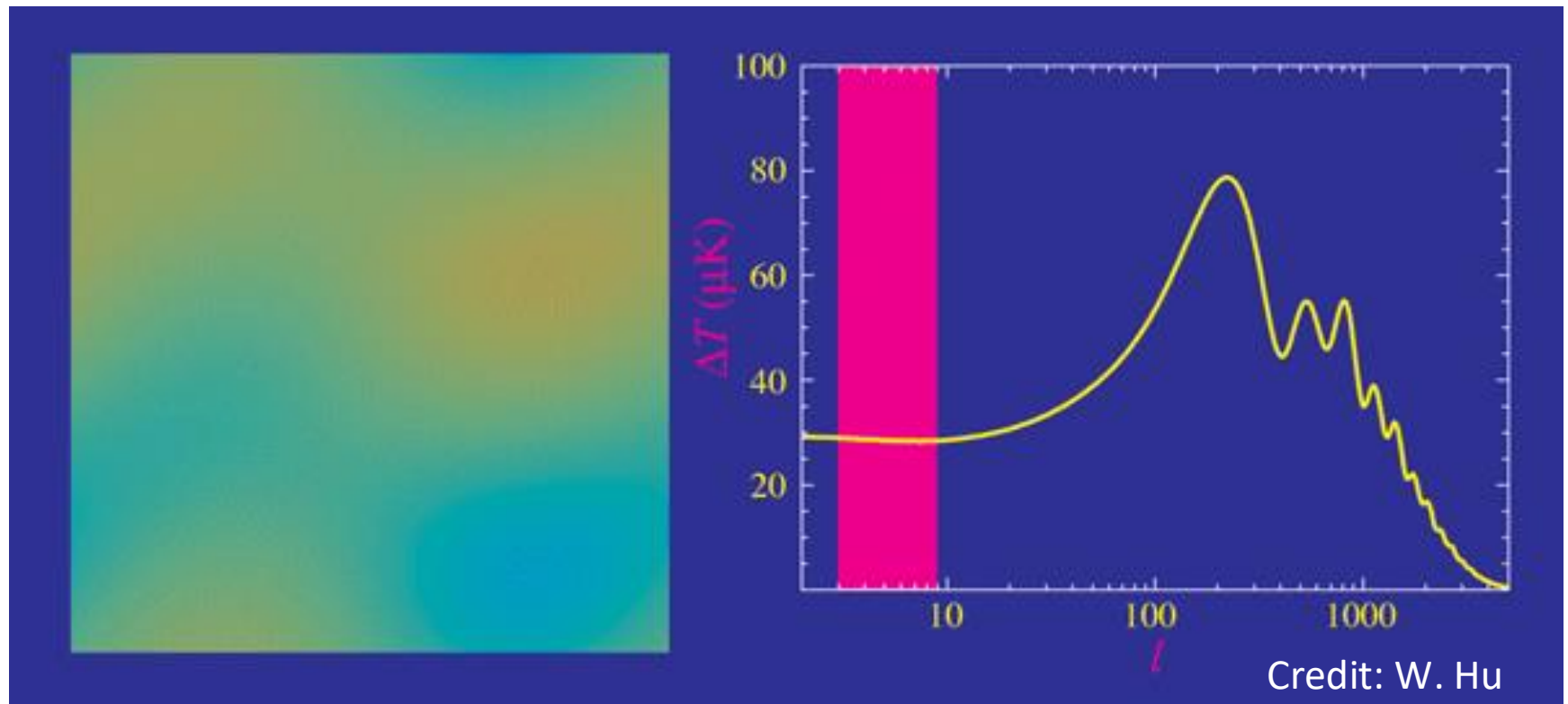
Cl's, Dl's and the 2-point correlation function

- We can relate the angular power spectrum to the 2-point correlation function in real space using the Legendre polynomials and the addition theorem:

$$\sum_m Y_{\ell m}^*(\mathbf{n}_i) Y_{\ell m}(\mathbf{n}_j) = \frac{2\ell + 1}{4\pi} P_\ell(\hat{\mathbf{n}}_i \cdot \hat{\mathbf{n}}_j)$$
$$\langle \Theta_i \Theta_j \rangle = \sum_\ell \frac{2\ell + 1}{4\pi} C_\ell P_\ell(\hat{\mathbf{n}}_i \cdot \hat{\mathbf{n}}_j)$$

- Because of isotropy, the two-point correlation function depends only on the angular separation in the sky θ , not on the orientation of the separation.
- The Dl's are just the contribution to the total temperature variance per logarithmic interval in l

$$\langle \Theta(\hat{\mathbf{n}}) \Theta(\hat{\mathbf{n}}) \rangle_{\ell \pm \Delta\ell/2} \approx \Delta\ell \frac{2\ell + 1}{4\pi} C_\ell \xrightarrow{\Delta\ell \sim \ell \Delta \ln \ell} \frac{\ell^2}{2\pi} C_\ell$$



Each of these maps are drawn from gaussian distributions with 0 mean and variance given by the CI in the corresponding pink band. There is an infinite number of possible realizations.

An estimator for the Cl's

$$\langle a_{lm} a_{l'm'}^* \rangle = \delta_{ll'} \delta_{mm'} C_l.$$

- We only observe one universe=> average over many realizations of the universe not possible.
- Because of isotropy, all the m-modes $a_{\ell m}$ with the same ℓ are drawn from a gaussian with the same theoretical C_ℓ . An estimator of C_ℓ is then:

$$\hat{C}_\ell = \frac{1}{2\ell + 1} \sum_m a_{lm} a_{lm}^*$$

Sample Variance

- The expected value is $\langle C_l \rangle = C_l$
- Since we only have $2l+1$ samples for each l , there is an **intrinsic uncertainty!**

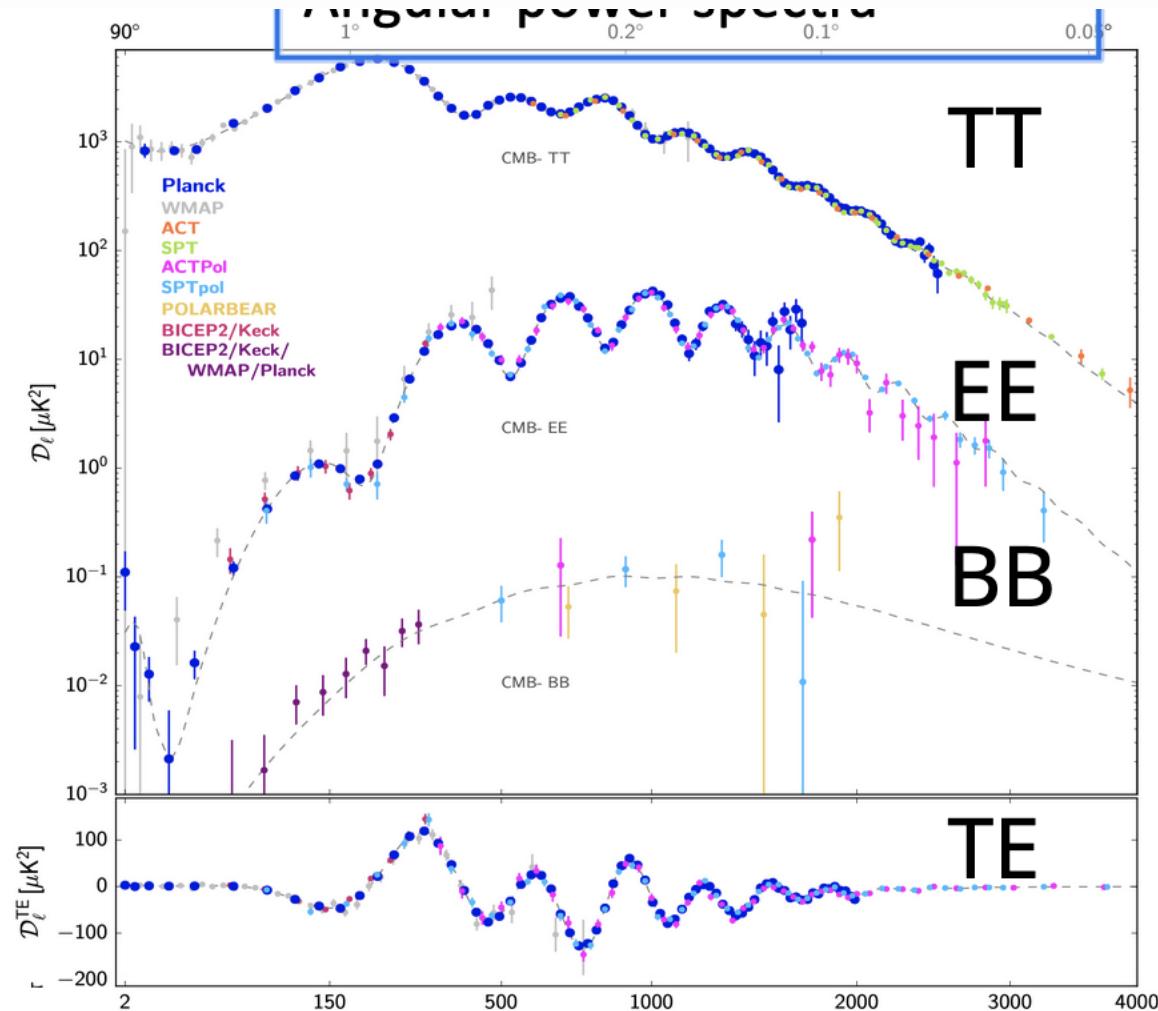
$$\hat{C}_\ell = \frac{1}{2\ell + 1} \sum_m a_{\ell m} a_{\ell m'}^*$$

$$\begin{aligned} \frac{\sigma_{C_\ell}^2}{C_\ell^2} &= \frac{\langle (\hat{C}_\ell - C_\ell)(\hat{C}_\ell - C_\ell) \rangle}{C_\ell^2} = \frac{\langle \hat{C}_\ell \hat{C}_\ell \rangle - C_\ell^2}{C_\ell^2} \\ &= \frac{1}{(2\ell + 1)^2 C_\ell^2} \langle \sum_{mm'} a_{\ell m}^* a_{\ell m} a_{\ell m'}^* a_{\ell m'} \rangle - 1 \\ &= \frac{1}{(2\ell + 1)^2} \sum_{mm'} (\delta_{mm'} + \delta_{m-m'}) = \frac{2}{2\ell + 1} \end{aligned}$$

$$\sigma_{C_\ell}^2 = \frac{2}{2\ell + 1} C_\ell^2$$

For a gaussian field, Wick's theorem says that any N-point (N even) statistics can be written as a function of the 2-point correlation function. Useful relations in the back-up slides.

Temperature and polarization power spectra



$$C_{Tl} = \frac{1}{2l+1} \sum_m \langle a_{T,l,m}^* a_{T,l,m} \rangle$$

$$C_{El} = \frac{1}{2l+1} \sum_m \langle a_{E,l,m}^* a_{E,l,m} \rangle$$

$$C_{Bl} = \frac{1}{2l+1} \sum_m \langle a_{B,l,m}^* a_{B,l,m} \rangle$$

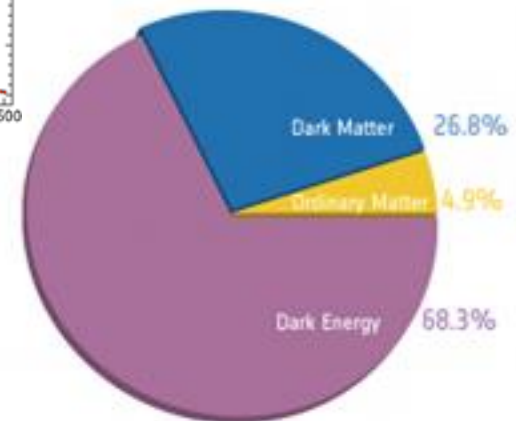
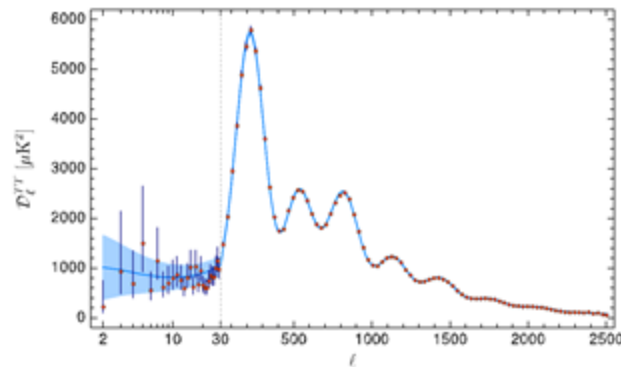
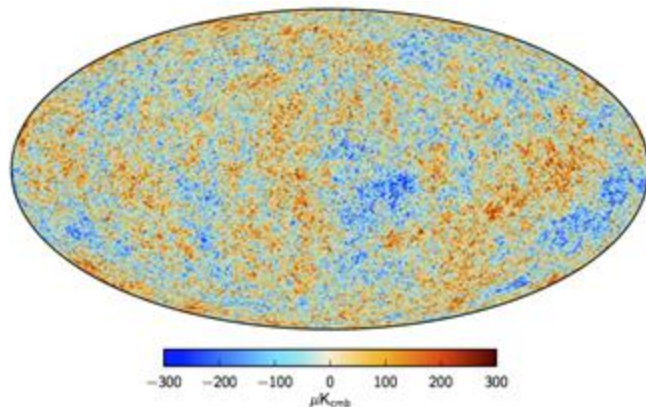
$$C_{Cl} = \frac{1}{2l+1} \sum_m \langle a_{C,l,m}^* a_{C,l,m} \rangle$$

Overview

1. A short introduction
2. The CMB sky and observations
 1. Monopole, dipole, anisotropies, polarization
 2. Ground, balloon and space experiments
3. From observations to cosmological parameters
 1. The CMB angular power spectrum and estimator
 2. **Bayes theorem and likelihood**
4. Latest results on cosmology (very Planck and SPT-3G oriented)

From maps to parameters

We want to measure parameters (cosmological+others) from maps. We want to evaluate the posterior distribution of the parameters θ given the data d , $P(\theta | d)$.





Bayes theorem

- To relate the posterior of the parameters given the data to the probability of the data given the parameters (the likelihood), use Bayes:

$$P(d, \theta) = P(d | \theta)P(\theta) = P(\theta | d)P(d)$$

$$\Rightarrow P(\theta | d) = \frac{P(d | \theta)P(\theta)}{P(d)}$$

Diagram illustrating the components of Bayes' theorem:

- Posterior**: $P(\theta | d)$ (indicated by a blue arrow pointing to the left side of the equation)
- Likelihood**: $P(d | \theta)$ (indicated by a blue arrow pointing to the numerator)
- Prior**: $P(\theta)$ (indicated by a blue arrow pointing to the numerator)
- Denominator**: $P(d)$ (indicated by a blue arrow pointing to the denominator)

So, what is the likelihood for CMB data?



Map-based Likelihood

- CMB maps (\mathbf{m}) have **gaussian** fluctuations with zero **mean** and pixel-space covariance matrix \mathbf{M} .

$$\mathcal{L}(C_\ell) = \mathcal{P}(\mathbf{m}|C_\ell) = \frac{1}{|2\pi\mathbf{M}|^{1/2}} \exp\left(-\frac{1}{2}\mathbf{m}^T \mathbf{M}^{-1} \mathbf{m}\right)$$

\mathbf{m} = **data** vector containing the pixels of the **map**.

$\mathbf{M} = \mathbf{S}(\theta) + \mathbf{N}$ = pixel **covariance** matrix, where \mathbf{S} is the two-point correlation function that depends on cosmological parameters

- In practice, this is only used at large scales/low resolutions. Inversion and determinant of covmat is unfeasible for maps with $O(10^6)$ pixels.

The full-sky likelihood of the C_l

- Instead of using the maps, we can compress the information in the estimator of \hat{C}_l from the maps, and use that as our « data ».
- For an ideal noiseless full-sky experiment, temperature alone.

$$\hat{C}_\ell = \frac{1}{2\ell + 1} \sum_{m=-\ell}^{\ell} |a_{\ell m}|^2.$$

- The sum of the square of $\nu=2\ell+1$ normal $N(0,1)$ variables ($a_{lm}/\sqrt{C_l}$) (with C_l the theoretical C_l) has a χ^2 distribution, i.e.:

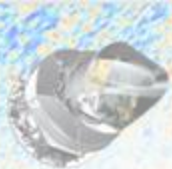
$$\hat{Y} = \sum_{m=-\ell}^{\ell} \left| \frac{a_{lm}}{\sqrt{C_\ell}} \right|^2 = (2\ell + 1) \hat{C}_\ell / C_\ell \rightarrow \chi^2 \text{ distribution with } \nu=2\ell + 1 \text{ d.o.f.}$$

$$\hat{C}_\ell = \hat{Y} C_\ell / (2\ell + 1) \rightarrow \Gamma \text{ distribution with } \nu=2\ell + 1 \text{ d.o.f.}$$

$$\langle \hat{C}_\ell \rangle = C_\ell$$

$$Var(\hat{C}_\ell) = \frac{2(C_\ell)^2}{\nu} = \frac{2(C_\ell)^2}{2\ell + 1} \quad \text{Cosmic Variance}$$

- See Percival and Brown 2006.



The full-sky likelihood of the C_l

- If we have polarization as well, this formalism extends to give a **Wishart** distribution for:

$$\hat{\mathbf{X}}_C = \begin{pmatrix} \hat{C}_l^{\text{TT}} \\ \hat{C}_l^{\text{TE}} \\ \hat{C}_l^{\text{EE}} \end{pmatrix}$$

- For $\nu \rightarrow \infty$, Gaussian distribution (for central limit theorem), with covariance:

$$\mathbf{Y} = \frac{1}{\nu} \begin{pmatrix} 2(C_l^{\text{TT}})^2 & 2C_l^{\text{TT}}C_l^{\text{TE}} & 2(C_l^{\text{TE}})^2 \\ 2C_l^{\text{TT}}C_l^{\text{TE}} & C_l^{\text{TT}}C_l^{\text{EE}} + (C_l^{\text{TE}})^2 & 2C_l^{\text{TE}}C_l^{\text{EE}} \\ 2(C_l^{\text{TE}})^2 & 2C_l^{\text{TE}}C_l^{\text{EE}} & 2(C_l^{\text{EE}})^2 \end{pmatrix}$$

Life is more complicated...

- Masking the sky, noise, beams invalidates the analytic calculation of the likelihood
- Pseudo-Cl in presence of sky cut:

$$\tilde{a}_{\ell m} = \int d\mathbf{n} \Delta T(\mathbf{n}) W(\mathbf{n}) Y_{\ell m}^*(\mathbf{n}),$$

Hivon 2002

$$= \sum_{\ell' m'} a_{\ell' m'} \int d\mathbf{n} Y_{\ell' m'}(\mathbf{n}) W(\mathbf{n}) Y_{\ell m}^*(\mathbf{n})$$

$$= \sum_{\ell' m'} a_{\ell' m'} K_{\ell m \ell' m'}[W],$$

- $\tilde{a}_{\ell m}$ are still gaussian, but not independent since they all depend on the sum of $a_{\ell' m'}$. $\tilde{a}_{\ell_1 m_1}$ and $\tilde{a}_{\ell_2 m_2}$ are correlated.

Life is more complicated...

- Masking the sky, noise, beams invalidates the analytic calculation of the likelihood

$$\tilde{C}_\ell = \frac{1}{2\ell + 1} \sum_{m=-\ell}^{\ell} |\tilde{a}_{\ell m}|^2$$

$$\tilde{a}_{\ell m} = \sum_{\ell' m'} a_{\ell' m'} K_{\ell m \ell' m'} [W]$$

The pseudo Cl are now the sum of the square of gaussians with different variances. Not distributed like a Wishart!



The Gaussian approximation

- For large degrees of freedom $\nu=2l+1$, the distribution of the \hat{C}_l tends to a gaussian distribution (central limit theorem).

$$-\ln \mathcal{L}(\hat{\mathbf{C}}|\mathbf{C}(\theta)) = \frac{1}{2} [\hat{\mathbf{C}} - \mathbf{C}(\theta)]^T \mathbf{C}^{-1} [\hat{\mathbf{C}} - \mathbf{C}(\theta)] + \text{const}$$

data

Model (that depends on the
parameters we want to determine)

C_l covariance matrix
(can be estimated with a fixed fiducial
set of parameters)

It works only at high- l (large dof). Used in
Planck, ACT, SPT.

What goes in the likelihood:

$$-\ln \mathcal{L}(\hat{\mathbf{C}}|\mathbf{C}(\theta)) = \frac{1}{2} [\hat{\mathbf{C}} - \mathbf{C}(\theta)]^T \mathbf{C}^{-1} [\hat{\mathbf{C}} - \mathbf{C}(\theta)] + \text{const}$$

Power spectrum estimation from the maps at different frequencies

$$\hat{\mathbf{C}}^{TT} = (\hat{\mathbf{C}}_{100 \times 100}^{TT}, \hat{\mathbf{C}}_{143 \times 143}^{TT}, \hat{\mathbf{C}}_{143 \times 217}^{TT}, \hat{\mathbf{C}}_{217 \times 217}^{TT})$$

$$\hat{\mathbf{C}}^{EE} = (\hat{\mathbf{C}}_{100 \times 100}^{EE}, \hat{\mathbf{C}}_{100 \times 143}^{EE}, \hat{\mathbf{C}}_{100 \times 217}^{EE}, \hat{\mathbf{C}}_{143 \times 143}^{EE}, \hat{\mathbf{C}}_{143 \times 217}^{EE}, \hat{\mathbf{C}}_{217 \times 217}^{EE})$$

$$\hat{\mathbf{C}}^{TE} = (\hat{\mathbf{C}}_{100 \times 100}^{TE}, \hat{\mathbf{C}}_{100 \times 143}^{TE}, \hat{\mathbf{C}}_{100 \times 217}^{TE}, \hat{\mathbf{C}}_{143 \times 143}^{TE}, \hat{\mathbf{C}}_{143 \times 217}^{TE}, \hat{\mathbf{C}}_{217 \times 217}^{TE}).$$


Model includes:

- the theoretical CMB power spectrum depending on a cosmological model
- the contribution from foregrounds
- Instrumental and systematic effects such as calibration, beams, etc...

Covariance matrix, estimated on simulations and/or analytically

$$\mathbf{C} = \begin{pmatrix} \mathbf{C}^{TTTT} & \mathbf{C}^{TTEE} & \mathbf{C}^{TTTE} \\ \mathbf{C}^{EETT} & \mathbf{C}^{EEEE} & \mathbf{C}^{EETE} \\ \mathbf{C}^{TETT} & \mathbf{C}^{TEEE} & \mathbf{C}^{TETE} \end{pmatrix}$$

$$\mathbf{C}^{TTTT} = \begin{pmatrix} (100 \times 100) \times (100 \times 100) & (100 \times 100) \times (143 \times 143) & (100 \times 100) \times (217 \times 217) & (100 \times 100) \times (143 \times 217) \\ (143 \times 143) \times (100 \times 100) & (143 \times 143) \times (143 \times 143) & (143 \times 143) \times (217 \times 217) & (143 \times 143) \times (143 \times 217) \\ (217 \times 217) \times (100 \times 100) & (217 \times 217) \times (143 \times 143) & (217 \times 217) \times (217 \times 217) & (217 \times 217) \times (143 \times 217) \\ (143 \times 217) \times (100 \times 100) & (143 \times 217) \times (143 \times 143) & (143 \times 217) \times (217 \times 217) & (143 \times 217) \times (143 \times 217) \end{pmatrix}$$



Exploring the likelihood to infer cosmological parameters

$$-\ln \mathcal{L}(\hat{\mathbf{C}}|\mathbf{C}(\theta)) = \frac{1}{2} [\hat{\mathbf{C}} - \mathbf{C}(\theta)]^T \mathbf{C}^{-1} [\hat{\mathbf{C}} - \mathbf{C}(\theta)] + \text{const}$$

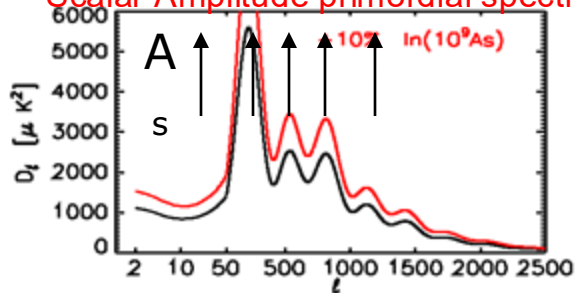
Since the dependence of the Cl on parameters is non-linear and non-trivial, we need Monte Carlo Markov Chains to explore the likelihood to map the posterior distribution of parameters.



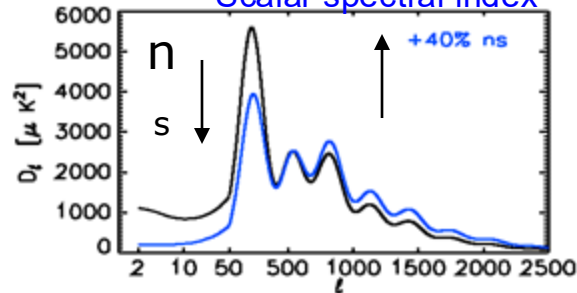
The Λ CDM model

Standard model of cosmology: General relativity to describe gravity, standard model of particles for particle interactions, cosmological constant for dark energy and cold dark matter.

Scalar Amplitude primordial spectrum



Scalar spectral index



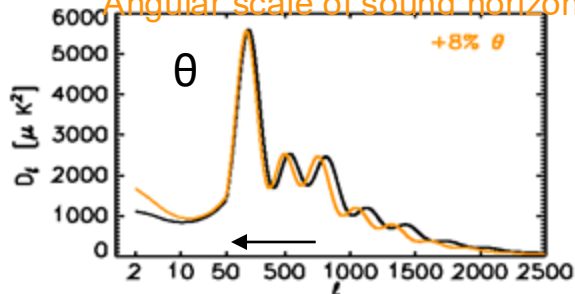
6 parameters:

- Initial conditions A_s , n_s
- Acoustic scale of sound horizon θ
- Reionization τ
- Dark Matter density $\Omega_c h^2$
- Baryon density $\Omega_b h^2$

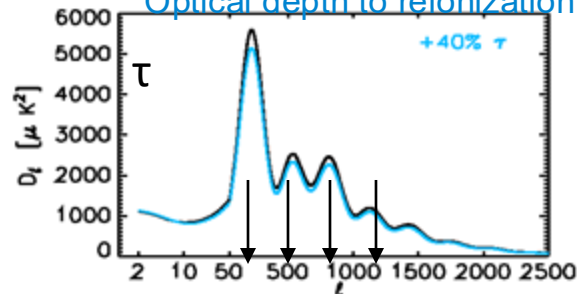
Assumptions:

- Adiabatic initial conditions
- $N_{\text{eff}} = 3.046$
- 1 massive neutrino 0.06eV.
- Tanh reionization ($\Delta z = 0.5$)

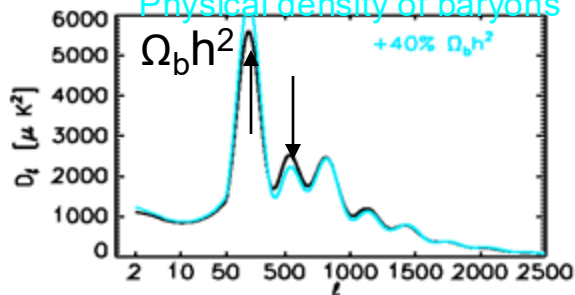
Angular scale of sound horizon



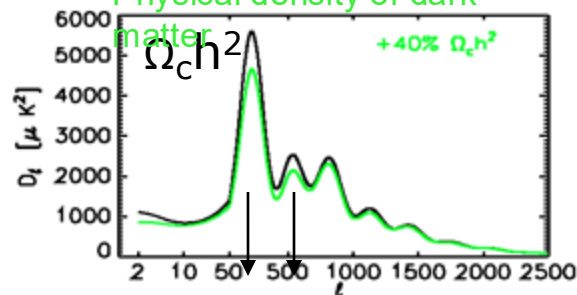
Optical depth to reionization



Physical density of baryons



Physical density of dark matter



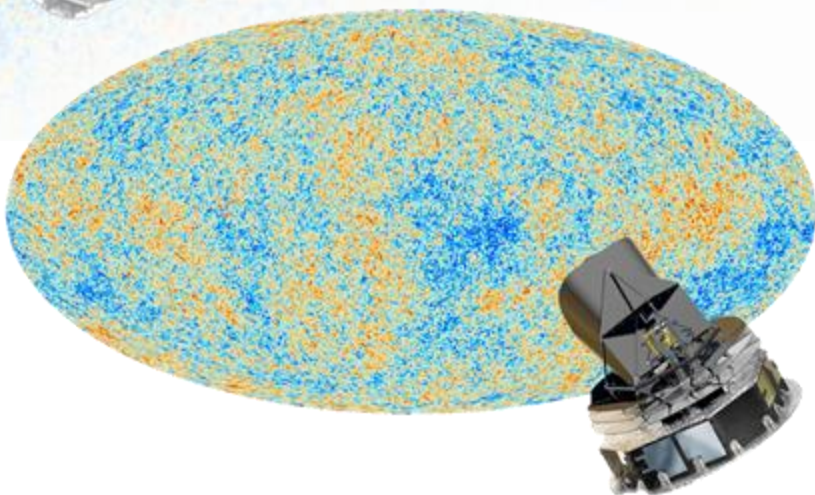
The CMB is a laboratory to constrain cosmology and fundamental physics



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The Planck satellite



3rd generation full sky satellites (COBE, WMAP)
Launched in 2009, operated till 2013.

2 Instruments, 9 frequencies to **disentangle CMB from foregrounds.**

LFI:

- 22 radiometers at **30, 44, 70 Ghz.**

HFI:

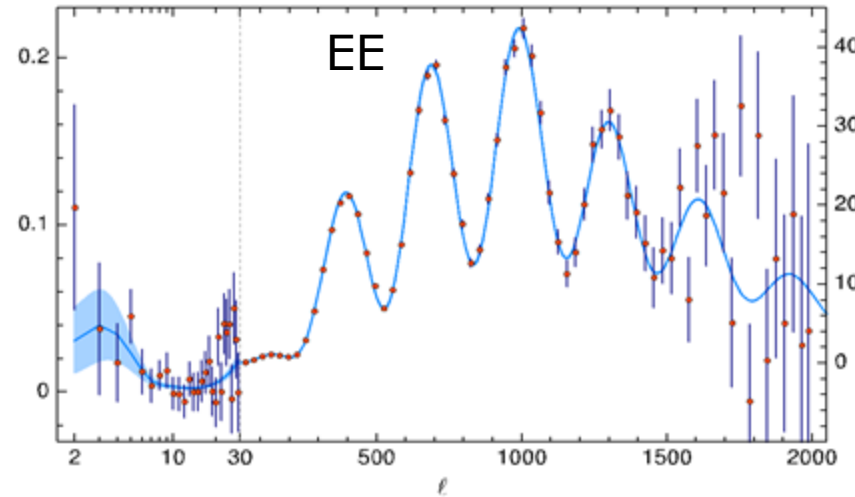
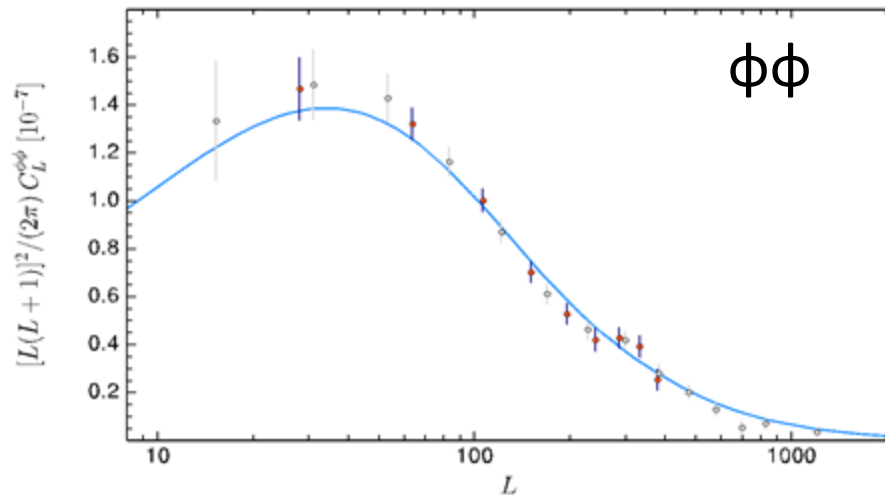
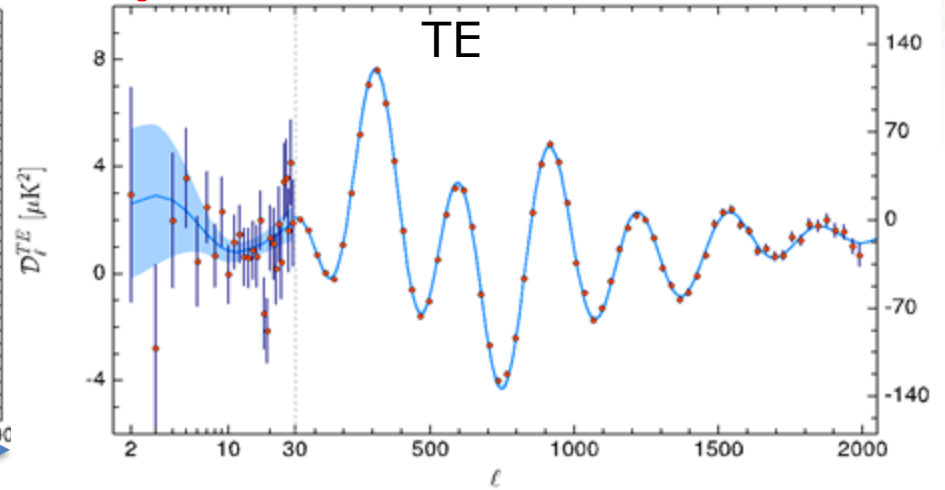
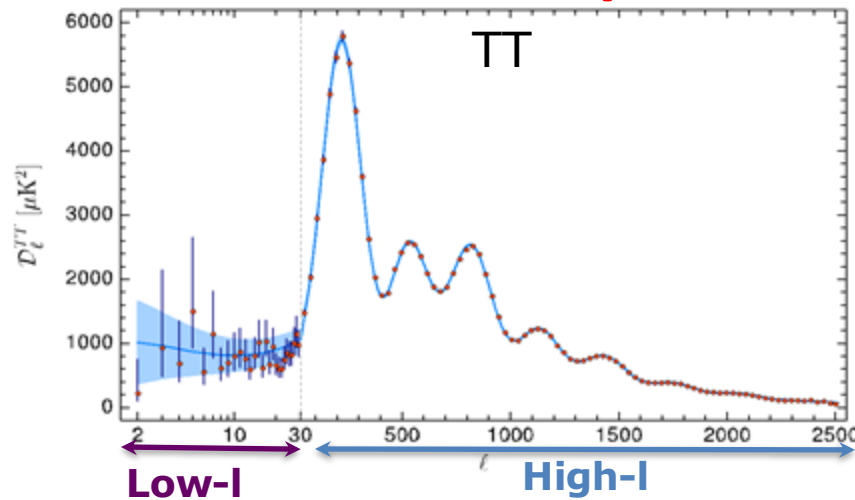
- 50 bolometers (32 polarized) at **100, 143, 217, 353, 545, 857 Ghz.**
- **30-353 Ghz polarized.**

- **1st release 2013: Nominal mission**, 15.5 months, Temperature only (large scale polarization from WMAP).
- **2nd release 2015: Full mission**, 29 months for HFI, 48 months for LFI, Temperature + Polarization, large scale pol. from LFI.
Intermediate results 2016: low- l polarization from HFI
- **3rd release 2018 (PR3): Full mission, improved polarization, low/high- l from HFI.**
Better control of systematics specially in pol., still systematics limited.

Post PR3: new maps PR4 (LFI+HFI map-making, improved low- ell in polarization); new likelihoods and parameters from PR4; new maps Sroll2 (better low- ell polarization); several new estimation of opt. depth to reionization.

No substantially new result compared to PR3.

Planck 2018 power spectra



Planck collaboration 2018 VI.

Λ CDM is an excellent fit to the data. No evidence of preference for classical extensions of Λ CDM from Planck (and non-classical ones: dm annihilation, variation of fund. constants, primordial magnetic fields, variations in recombination, isocurvature, sterile neutrinos, dark energy, modified gravity).....

Baseline Λ CDM results

2018 (Temperature+polarization+CMB lensing)

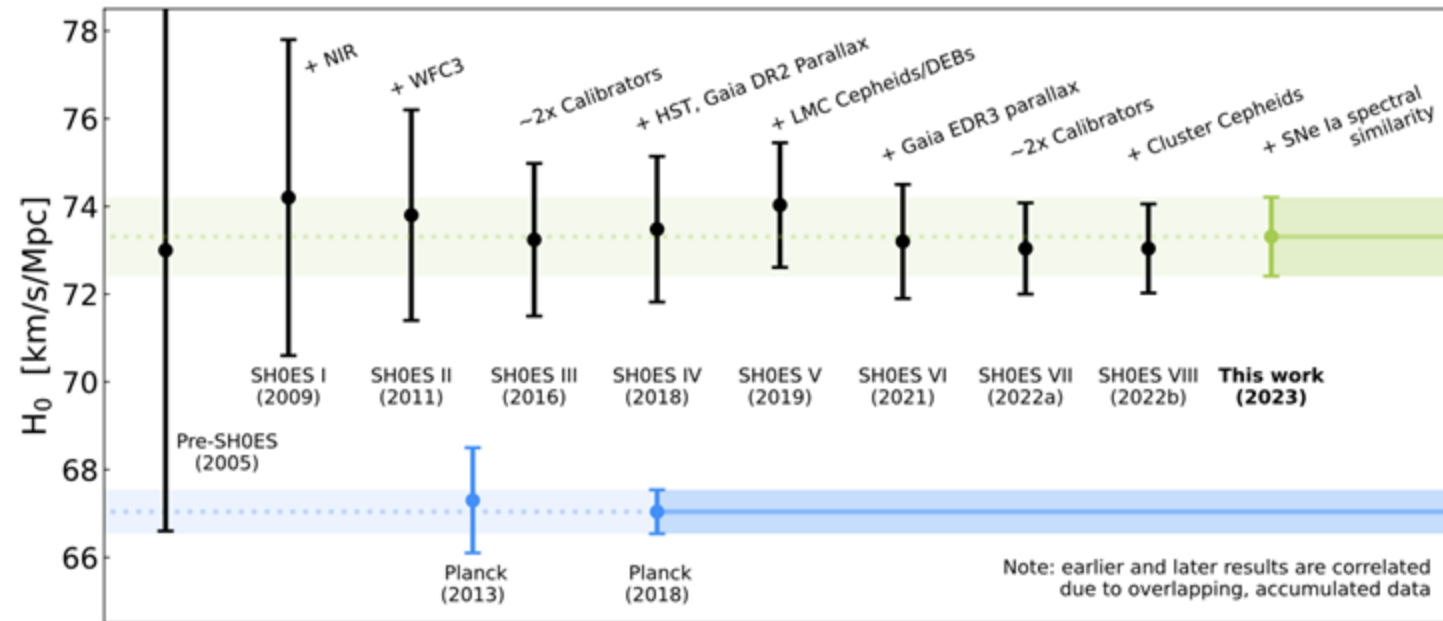
| | Mean | σ | [%] |
|---|---------|----------|------|
| $\Omega_b h^2$ Baryon density | 0.02237 | 0.00015 | 0.7 |
| $\Omega_c h^2$ DM density | 0.1200 | 0.0012 | 1 |
| 100θ Acoustic scale | 1.04092 | 0.00031 | 0.03 |
| τ Reion. Optical depth | 0.0544 | 0.0073 | 13 |
| $\ln(A_s 10^{10})$ Power Spectrum amplitude | 3.044 | 0.014 | 0.7 |
| n_s Scalar spectral index | 0.9649 | 0.0042 | 0.4 |
| H_0 Hubble | 67.36 | 0.54 | 0.8 |
| Ω_m Matter density | 0.3153 | 0.0073 | 2.3 |
| σ_8 Matter perturbation amplitude | 0.8111 | 0.0060 | 0.7 |

- Most of parameters determined at (sub-) percent level!
- **Best** determined parameter is the angular scale of sound horizon θ to **0.03%**.
- τ low and tight, reionization at $z \sim 7$.
- n_s is 8σ away from scale invariance (even in extended models, always $> 3\sigma$)
- **Best (indirect) 0.8% determination of the Hubble** constant to date.



The Hubble tension

The Hubble tension is the difference in the expansion rate of the universe today as measured from Supernovae IA calibrated with Cepheids and from the CMB and other early universe probes. Since the early universe measurements depend on a cosmological model, it could indicate the need of a change in the model, and thus **the discovery of new physics**.



Murakami+ 2023
Breuval+ 2024

5.7 σ tension

Supernovae IA
+ cepheid
SH0ES
 $H_0 = 73.29 \pm 0.90$
km/s/Mpc.

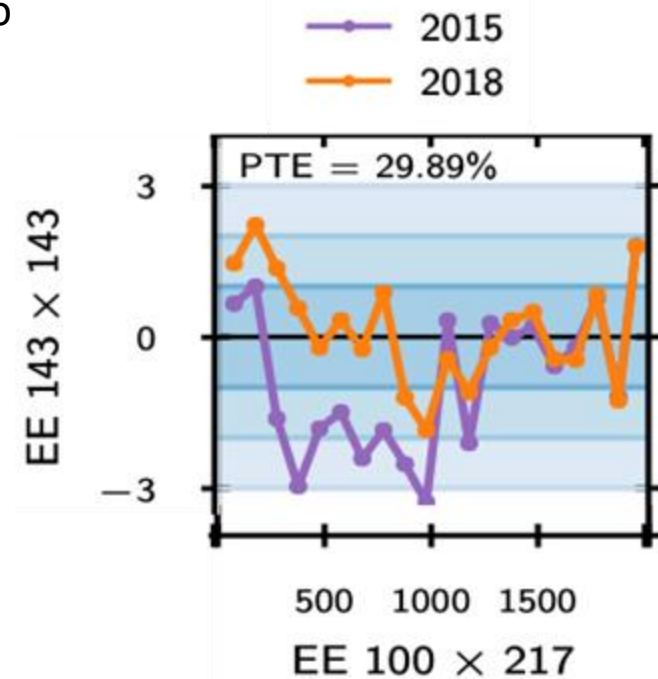
CMB (also BAO)
Planck 2018

$H_0 = 67.36 \pm 0.54$
km/s/Mpc.

The importance of robustness of results

- A **large impact** on the field (tens of thousands of citations)
- **Responsibility to provide the community robust results.**
- Many Tests:
 - **Redundancy** of the data is key in order to be able to do consistency tests at power spectrum or cosmological parameters level from **subsets of the data**, i.e.
 - from different frequency channels, which also corresponds to different detectors.
 - from different map cuts (half mission, versus detector sets)
 - from TT, TE or EE (model dependent).
 - Tested the consistency between a large number of **different analysis choices** on cosmological **parameters** (model dependent).
 - Compared different **analysis pipelines**, which was essential to improve the robustness of the final product.
 - **End-to-end simulations** also allowed us to validate the pipeline.

Difference between CMB-only, frequency power spectra in units of error bars in 2015 and 2018



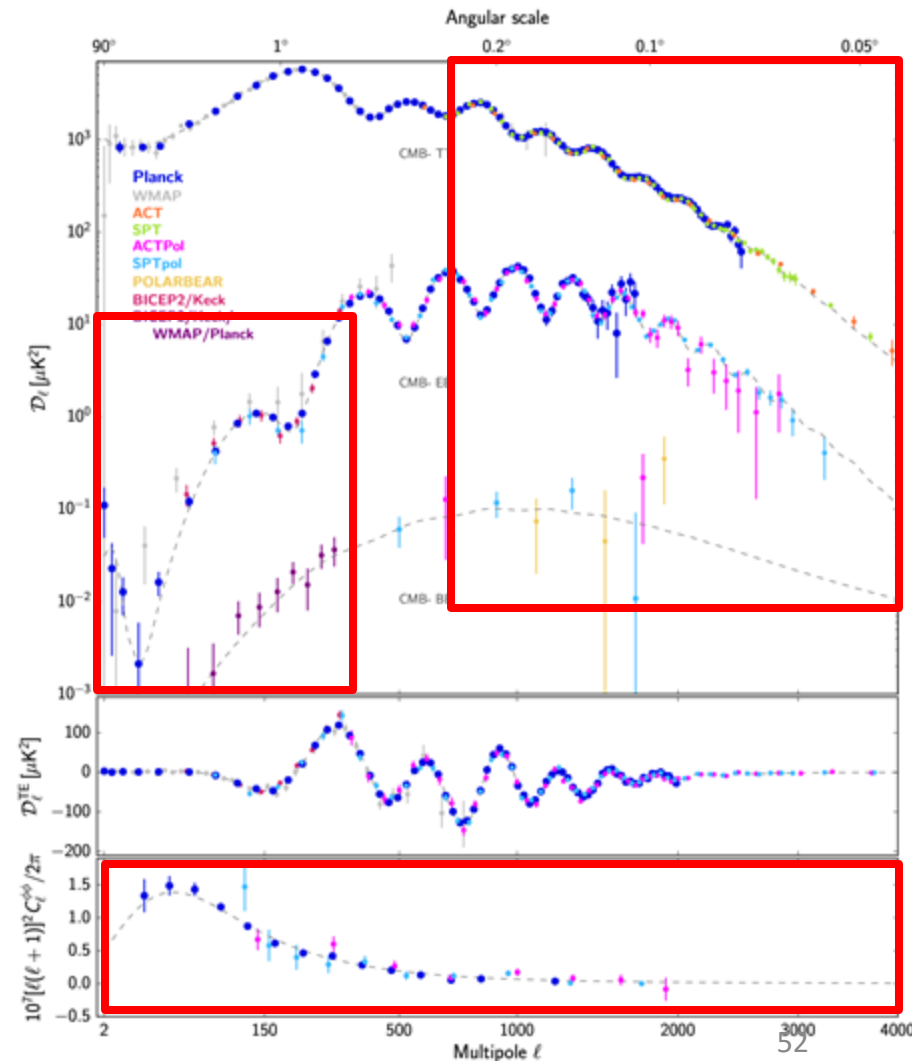
In the second Planck 2015 release, test failed for polarization data. culprit was uncorrected systematics.



Beyond Planck

- Planck results had a tremendous impact on cosmology. They confirmed the Λ CDM model. They were checked with many consistency tests.
- Planck opened new mysteries, such as the Hubble tension, that we will explore with upcoming and future experiments.
- There is still a very large amount of information in CMB to be uncovered. Upcoming experiments have two main goals:
 - at large scales in polarization, to detect primordial gravitational waves and measure reionization.
 - At small angular scales in polarization, to test cosmological models and the properties of the energy content of the universe.

NB: only full sky observations from satellites have access to very large scales ($l \sim 2$).



The quest for primordial gravitational waves

Current best constraints from experiments at the South Pole

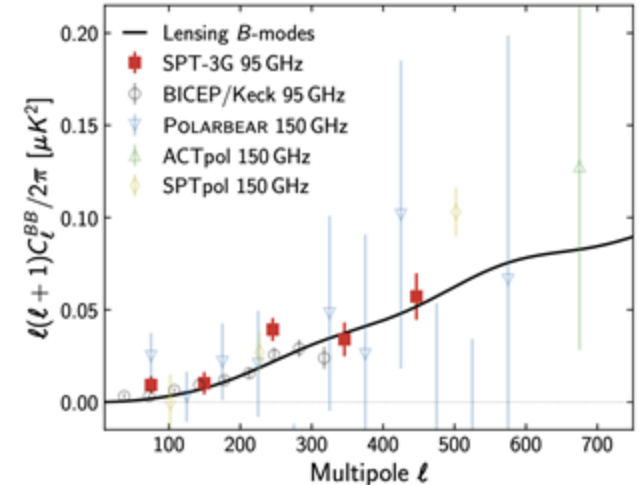
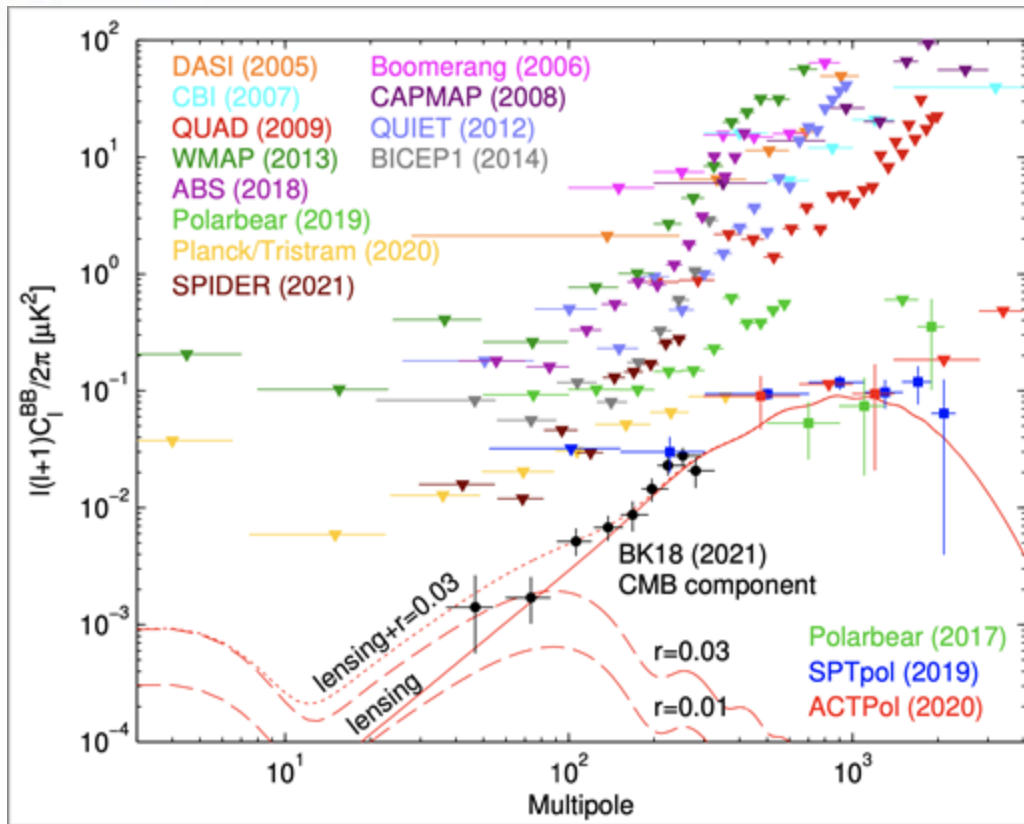


FIG. 7. Ground-based B -mode measurement landscape. Data are from SPTpol [32], ACT [28], POLARBEAR [1], and BICEP/Keck [4].

Best constraints on BB from
BICEP/Keck (BK collaboration 2021)



Second best constraints
from SPT-3G
(Zebrowski+ 2025)

Best CMB high resolution experiments

ACT:

- a 6m telescope observing from Atacama at 98, 150, 220Ghz.
- Recently had last data-release, ACT DR6.
- Observed 16000 square degrees, only 10000 used for cosmology.
- Recently substituted by Simons Observatory, which had first light this year.



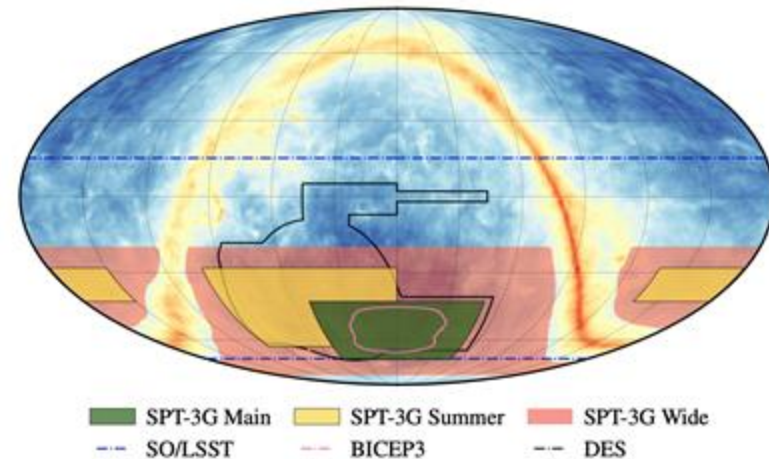
SPT-3G:

- a 10m telescope observing from the Amudsen-Scott station at the South Pole.
- Recently had release of results, SPT-3G D1, based on 1500 square degrees.
- Already observed a total of 10000 square degrees, analysis in progress.



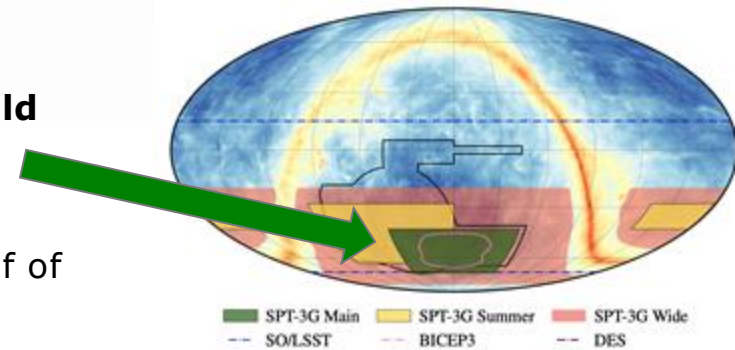
The South Pole Telescope

- SPT is a 10m telescope at the South Pole. It is observing with its third generation camera, SPT-3G, with $\sim 16,000$ detectors at 95, 150, and 220 GHz at high resolution (~ 1 arcmin).
- It has observed **25%** of the sky:
 - **Main** Winter Field **1500 deg²** (5yrs done, 2yr TBD)
Coadded noise **$\sim 17X$** better than Planck.
 - **Summer** fields **2650 deg²** (4yrs done),
coadded noise **4X** better than Planck.
 - **Wide** field **6000 deg²** (1yr)
coadded noise **3X** better than Planck.
- Many scientific goals:
 - **Cosmological constraints from CMB primary anisotropies and CMB lensing**
 - Delensing of the BICEP/Keck field to improve constraints on tensor to scalar ratio r .
 - High-ell TT foregrounds (including kSZ), Cross-correlations with other surveys, High-z galaxies, Clusters of galaxies, Transients etc...



Published power spectrum SPT-3G results from Main field

Results based on the observations of the small but **deep Main field** observed during the Austral winter.



- **4 months of observations "SPT-3G 2018"**: Used only half of the focal plane [published in 2021-2023]:
 - Analysis of **TEEE** (Dutcher,...SG+ 2021, Balkenhol,...SG+ 2021)
 - Analysis of **TTTEE** (Balkenhol,...SG+ 2023)
- **2 years of observations "SPT-3G 2 year Main field"**: Based on ~16 months of observations with the full focal plane.
 - Cosmology from CMB lensing and delensed EE power from polarization with **MUSE** (Ge., Millea,...SG+ 2024)
 - Constraints on Inflationary Gravitational Waves from large scale **BB** (Zebrowski,...SG+ 2025)
 - SPT-3G D1: CMB temperature and polarization power spectra and cosmology from 2019 and 2020 observations of the SPT-3G Main field (Camphuis,...SG+ 2025)

SPT-3G D1: CMB temperature and polarization power spectra and cosmology from 2019 and 2020 observations of the SPT-3G Main field

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G. P. Lynch⁷, Y. Omori^{9,4}, C. Trenafileva¹⁰, A. J. Anderson^{11,4,9}, B. Ansarinejad¹², M. Archipley^{9,4},
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W. L. Holzapfel⁸, J. C. Hood⁴, A. Hryciuk^{3,4}, F. Kéruzoré², L. Knox⁷, M. Korman²⁹, K. Korneelje^{9,4,2},
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A. A. Stark³³, J. Stephen¹⁴, C. Tandoi¹⁷, B. Thorne⁷, C. Umiltà²³, J. D. Vieira^{17,23,10}, A. Vitrier¹,
Y. Wan^{17,10}, N. Whitehorn²⁵, W. L. K. Wu^{5,13}, M. R. Young^{11,4} and J. A. Zebrowski^{4,9,11}

(SPT-3G Collaboration)

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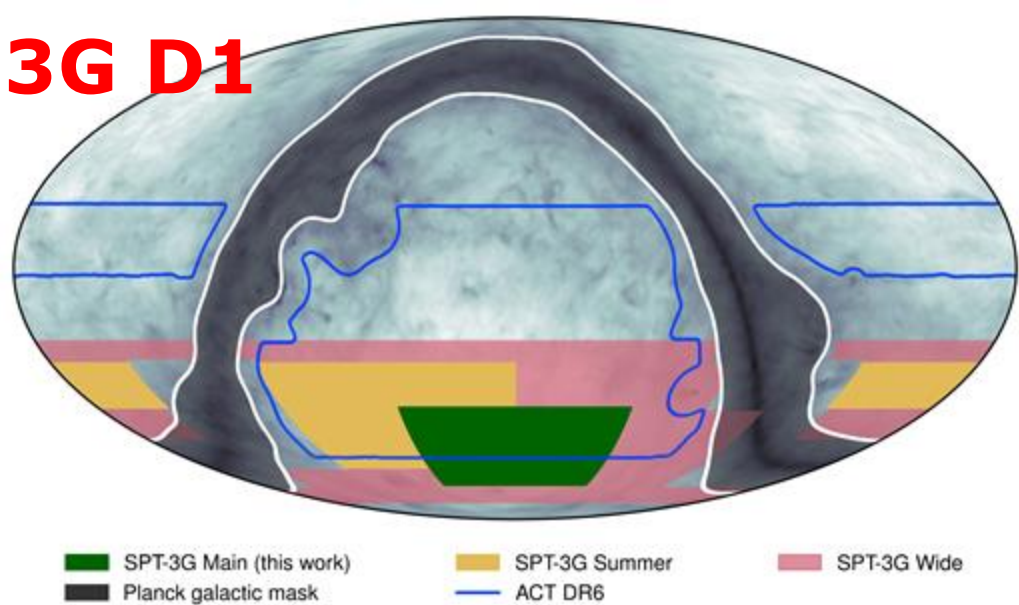
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A New Dataset: SPT-3G D1

- SPT-3G D1: observations taken in the 2019 & 2020 austral winter seasons (March to November) on the SPT-3G Main field
 - Much larger than the 2018 dataset ($\sim 2\times$ detectors, $\sim 4\times$ observing time)
 - Small, deep survey complementing *Planck* and ACT

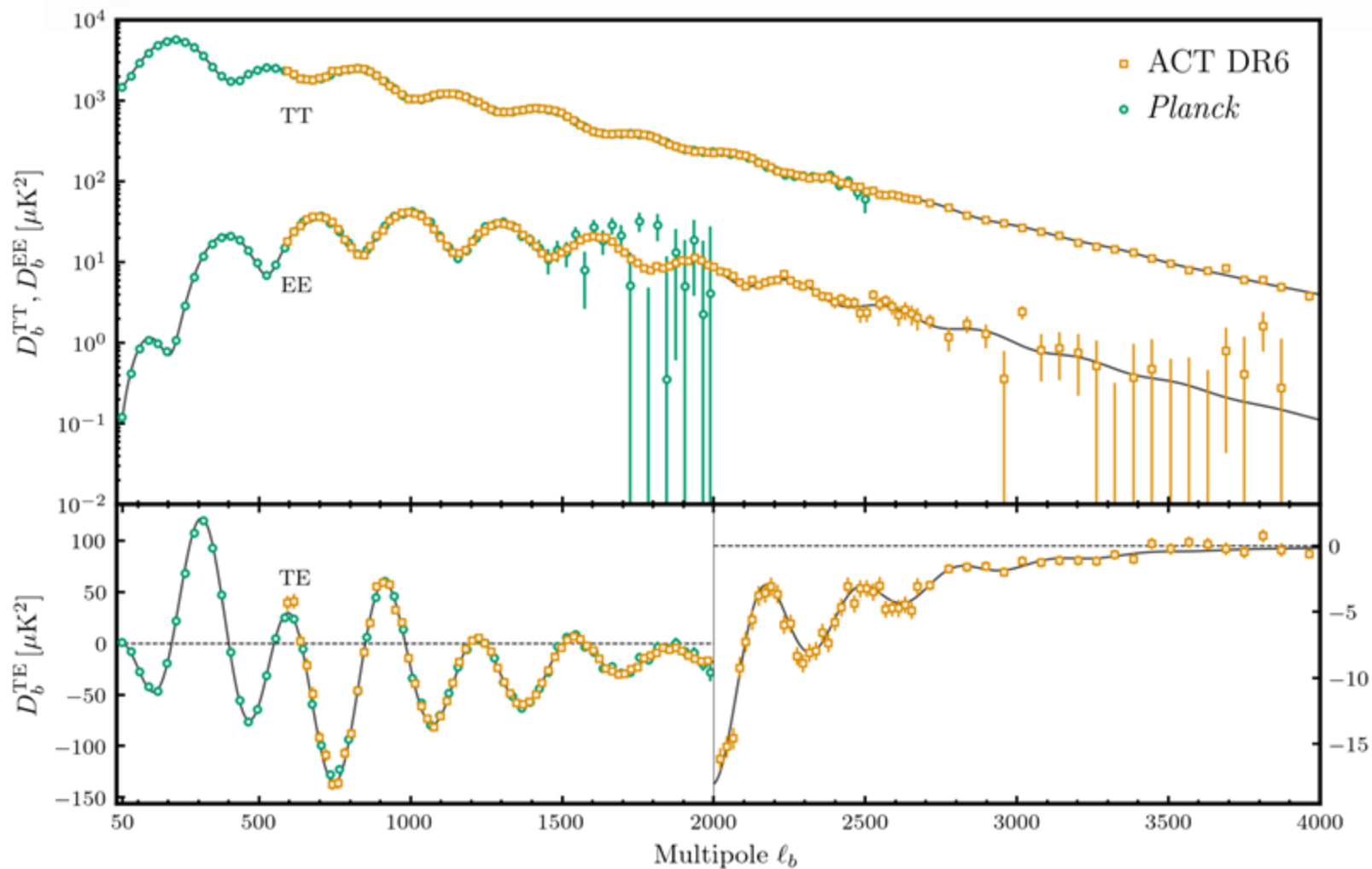
Deepest CMB maps at arcminute resolution for TT/TE/EE measurements

Planck PR3 numbers based on *Planck* 2018 results IV
ACT DR6 numbers from Næss *et al.*, 2025

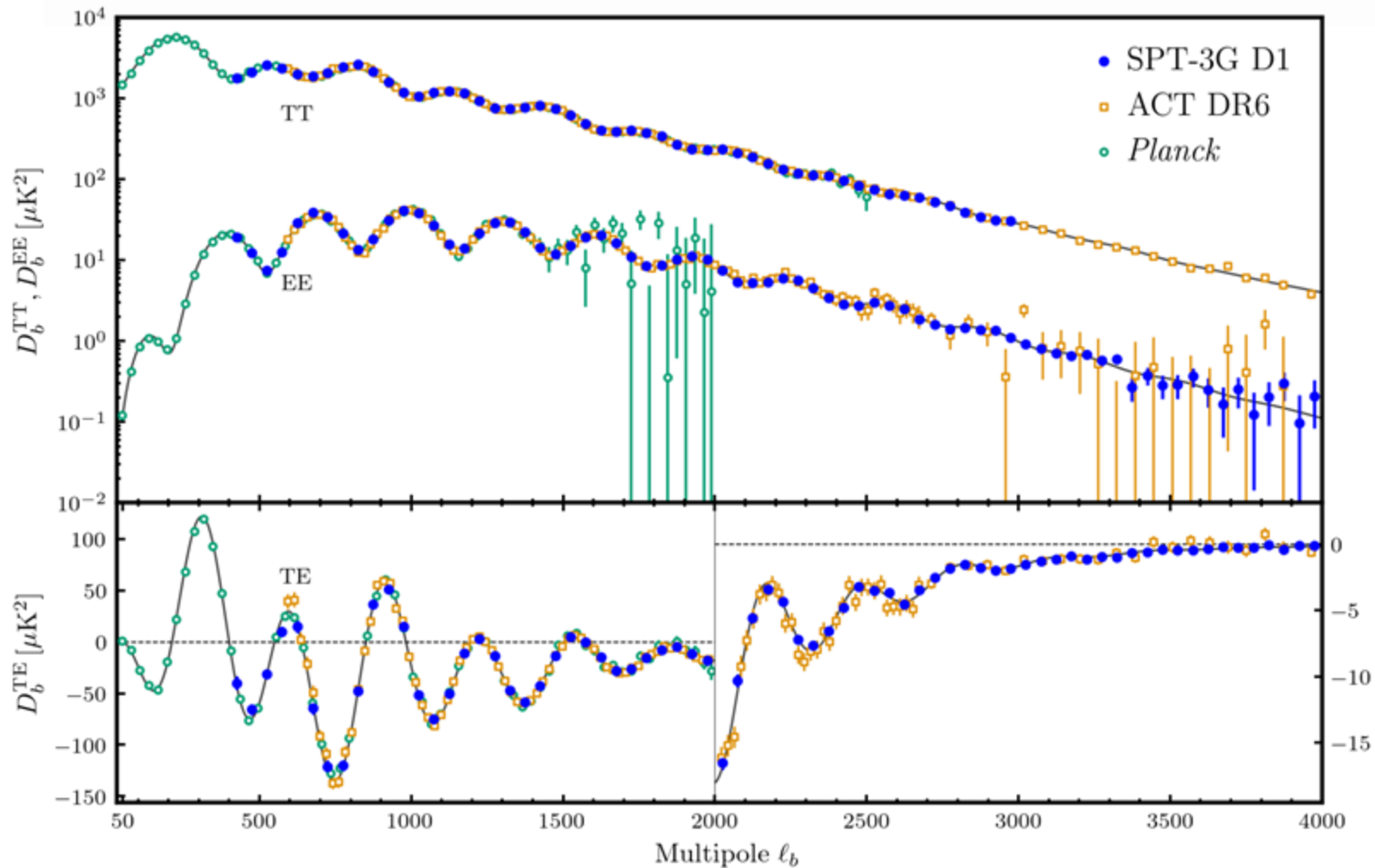


| | Observed sky fraction [%] | Coadded noise level [$\mu\text{K-arcmin}$] |
|-------------------|---------------------------|--|
| <i>Planck</i> PR3 | 100 | 35 |
| ACT DR6 | 45 (25 for cosmology) | 10 |
| SPT-3G D1 | 4 | 3.3 |

CMB angular power spectra



CMB angular power spectra

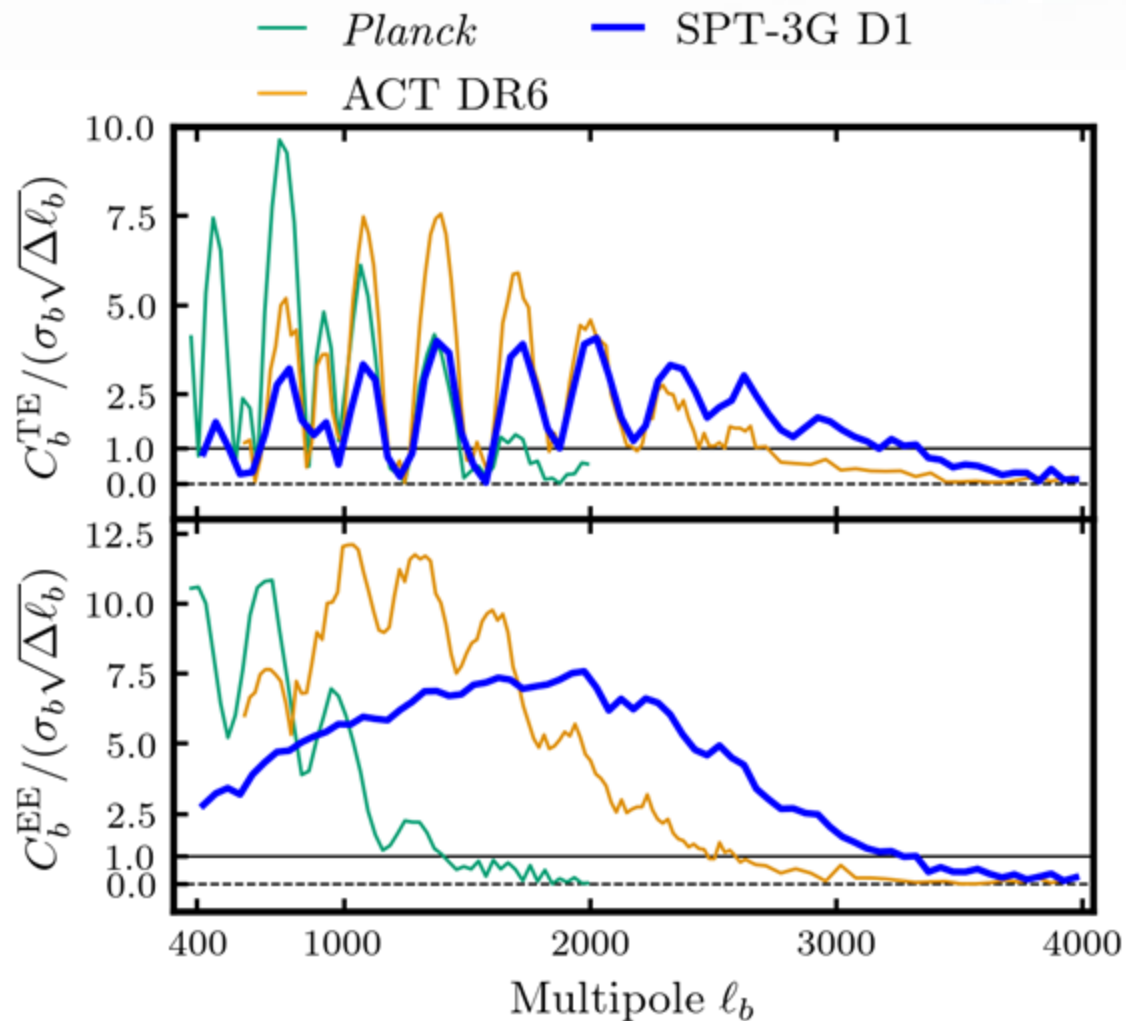


Complementary CMB experiments

| | Observed sky fraction [%] | Coadded noise level [$\mu\text{K-arcmin}$] |
|-------------------|---------------------------|--|
| <i>Planck</i> PR3 | 100 | 35 |
| ACT DR6 | 45 | 10 |
| SPT-3G D1 | 4 | 3.3 |

SPT-3G D1 is the tightest band power measurement:

- In TE at ℓ in [2200,4000]
- In EE at ℓ in [1800,4000]



Likelihood

Foreground and nuisance
model improved over SPT-
3G 2018

<http://ascl.net/1102.026>
<http://doi.org/10.1088/1475-7516/2011/07/034>
<http://doi.org/10.21105/astro.2305.06347>
<http://doi.org/10.1093/mnras/stac064>
<http://doi.org/10.48550/arXiv.2503.13183>

More details in
the paper !



$$-\ln \mathcal{L}(\hat{C}|C^{\text{model}}(\theta)) \propto \frac{1}{2} \left[\hat{C}_b - C_b^{\text{model}}(\theta) \right] \Sigma_{bb'}^{-1} \left[\hat{C}_{b'} - C_{b'}^{\text{model}}(\theta) \right]$$



Semi-analytical
covariance matrix from
Camphuis, ... SG et al,
2023

Differentiable and robust python JAX-
likelihood code Balkenhol,...SG et al. 2024

Analysis validation

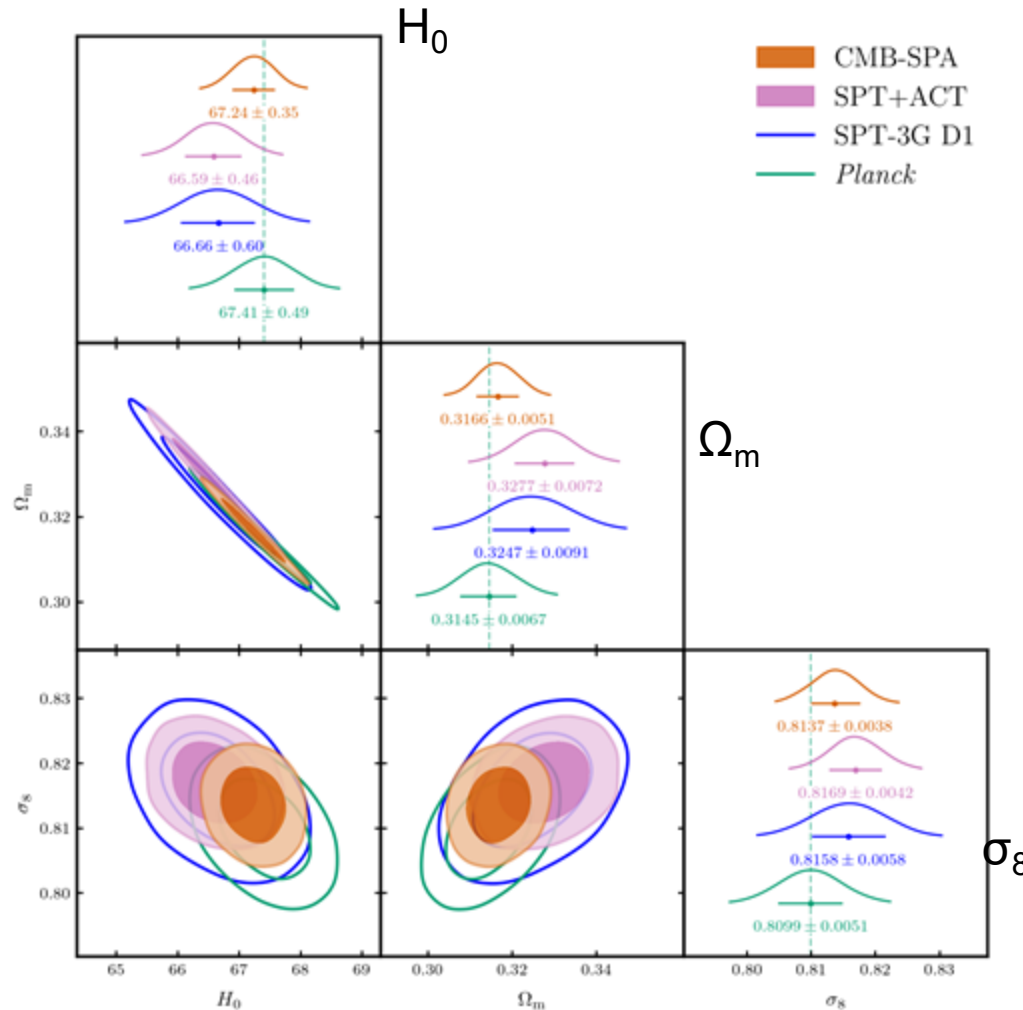
- Validation of the pipeline is done **blind**, without looking at obtained cosmological parameters or comparing to previous experiments. We allow changes after unblinding, but report all of the changes in a transparent way.
- Validation tests include:
 - Null tests at the map level, where we split the data in two according to some criterion (Sun, moon, azimuth, year, scan direction, detector wafers) and then take the difference between the two maps.
 - Differences at the power spectrum level between different spectra at different frequencies, to check that the CMB signal is the same.
 - Differences at the cosmological parameter level from different frequency channels, assuming LCDM.
 - Check that the pipeline is unbiased on simulations.

Data Sets

- **SPT-3G D1:**
 - D1 = Observations of the Main field in 2019-2020.
 - This work, i.e. SPT-3G Main field T&E* data + $\Phi\Phi$ band-powers from [Ge et al \[SPT-3G\], 2024](#).
- **Planck:** Planck 2018 (PR3) [high- ℓ T&E + low- ℓ TT] ([Planck Collaboration et al., 2018](#)) + PR4 $\Phi\Phi$ band-powers ([Carron et al, 2022](#)).
- **SPT+ACT:** SPT-3G D1 + ACT DR6 T&E ([Louis et al \[ACT\], 2025](#)) + ACT DR6 $\Phi\Phi$ band-powers ([Madhavacheril et al \[ACT\], 2023](#); [Qu et al \[ACT\], 2023](#)).
- **CMB-SPA:** SPT-3G D1 + P-ACT([Louis et al \[ACT\], 2025](#)).
- τ_{reio} **prior:** for all the data sets above, we use a prior from *Planck* PR4 ([Akrami et al \[Planck\], 2020](#)) on $\tau_{\text{reio}} = 0.051 \pm 0.006$

* T&E = TT/TE/EE band-powers.

Λ CDM: Most Precise Constraints from CMB to Date



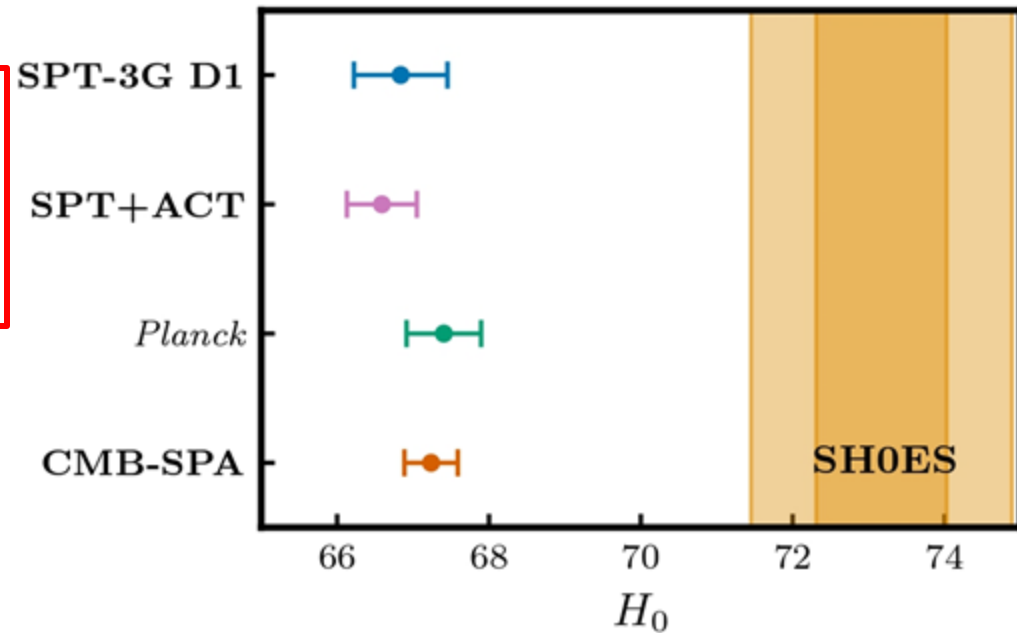
$H_0 = 67.24 \pm 0.35$ km/s/Mpc (CMB-SPA)

$H_0 = 67.41 \pm 0.49$ km/s/Mpc (*Planck*)

- With just 4% of the sky, SPT-3G's constraints on H_0 and σ_8 are comparable to Planck (within 25%) or ACT. Our data agree very well with Λ CDM predictions.
- CMB SPT+ACT finally reach Planck's precision (on some parameters)!
- CMB-SPA yields the most precise determination of Λ CDM parameters from a single probe. All three experiments agree with each other within 1.1σ . CMB science is very robust!

Λ CDM: Hubble Tension with SH0ES

Three independent and complementary experiments confirm the Hubble tension.

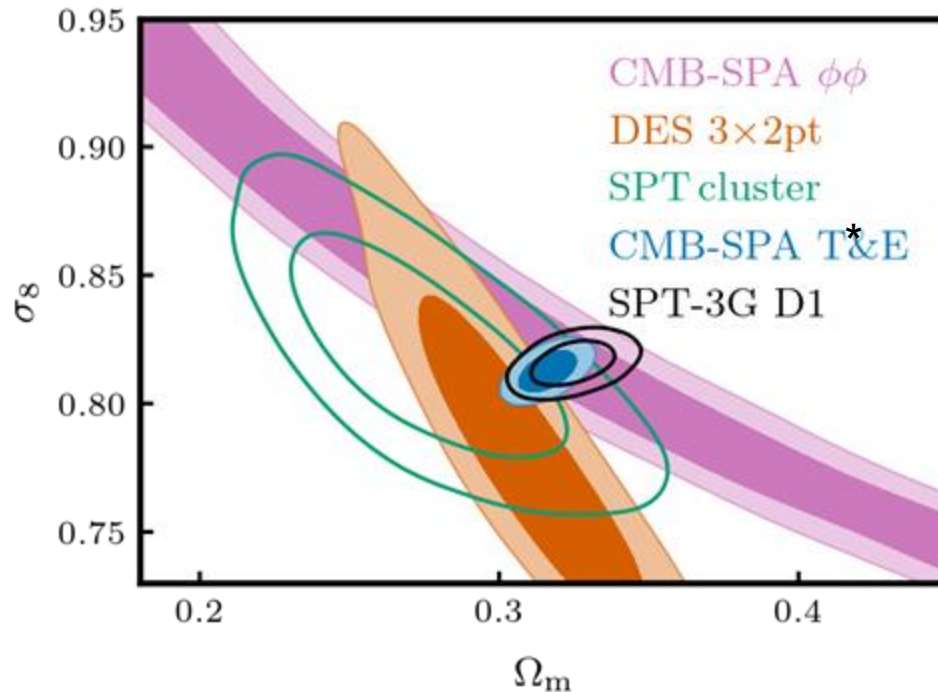


- Hubble Tension at 6.2σ from **SPT-3G** alone.
 $H_0 = 66.66 \pm 0.60$ km/s/Mpc
- **SPT+ACT** and **CMB-SPA** are at 6.8σ and 6.4σ tension, respectively.

[Breuval et al. \[SH0ES\], 2024](#)

$$H_0 = 73.17 \pm 0.86 \text{ km/s/Mpc}$$

Λ CDM: Clustering of Matter



SPT-3G D1 alone, with only 4% of the sky, constrain σ_8 almost as well as *Planck*.

- A variety of probes, spanning a wide range of epochs, are now consistent with each other (including the latest KiDS-legacy cosmic shear results, [KiDS collaboration, 2025](#)).

- $\sigma_8 = 0.8137 \pm 0.0038$
 $\Omega_m = 0.3166 \pm 0.0051$

} CMB-SPA

*SPT cluster = Constraints from SPT identified abundance of clusters ([Bocquet et al.\[SPT,DES\], 2024](#)).

Λ CDM Extensions

- The addition of SPT+ACT to *Planck* reduces the upper limit on Σm_ν by $\sim 30\%$:

$$\Sigma m_\nu < 0.25 \text{ eV (95\% C.L.) (Planck)}$$

$$\Sigma m_\nu < 0.18 \text{ eV (95\% C.L.) (CMB-SPA)}$$

This shows again the constraining power of SPT+ACT.

- We also explored extensions with N_{eff} , Y_P and modified recombination.

$$N_{\text{eff}} = 2.86 \pm 0.19 \text{ (Planck)}$$

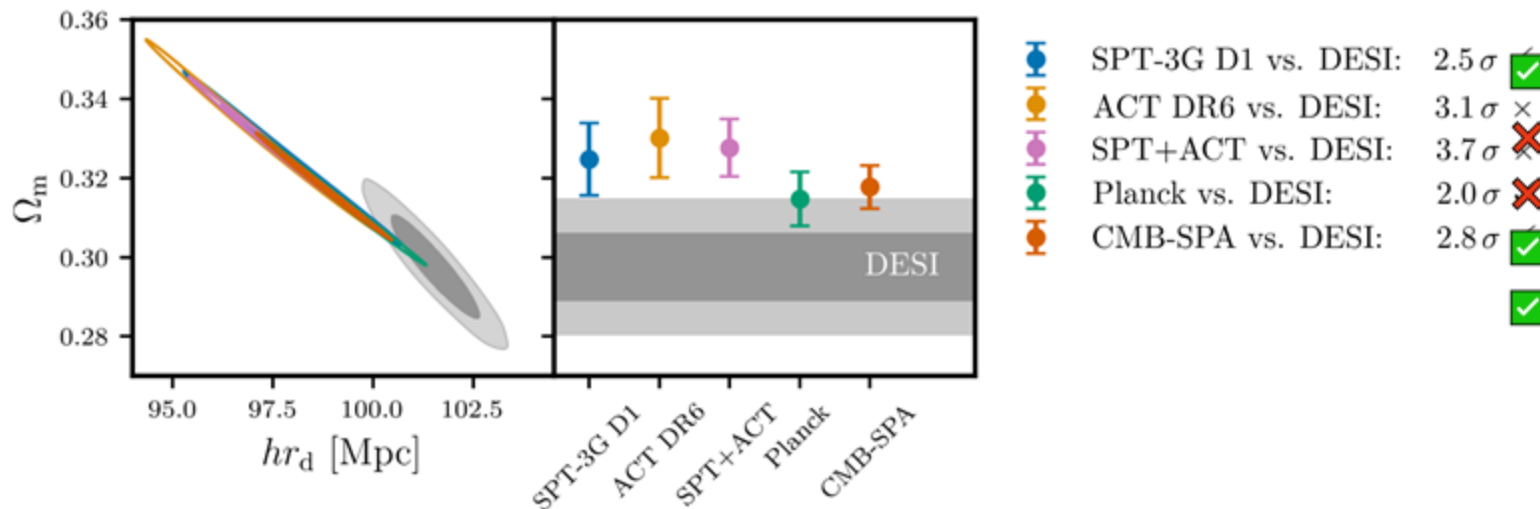
$$N_{\text{eff}} = 2.81 \pm 0.12 \text{ (CMB-SPA)}$$

We do not find any statistically significant deviations from Λ CDM.

A new CMB-BAO tension?

Evaluating the Consistency of CMB vs DESI in Λ CDM

- ACT and SPT+ACT above consistency threshold
- Planck data regularise combined results, **CMB-SPA consistent with DESI**
- Given borderline differences, **joint analyses to be performed with caution**



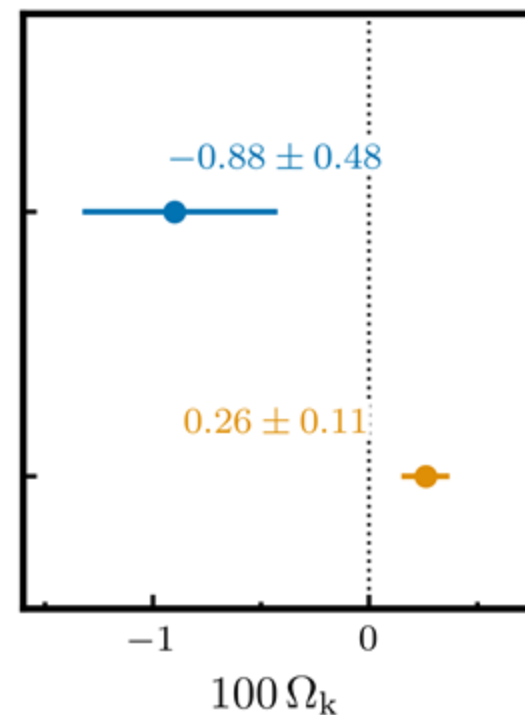
Constraints from CMB and BAO data beyond Λ CDM

Differences between CMB and DESI can be accommodated by 2-3 σ deviations from Λ CDM.

| Model Class | Preference over Λ CDM |
|-------------------------------------|-------------------------------|
| Rescaling of lensing in CMB | 3.1 σ |
| Light relics | <1.5 σ |
| Modified recombination | 2.0 σ |
| Spatial curvature | 2.5 σ |
| Spatial curvature and electron mass | 2.1 σ |
| Neutrino mass | 2.8 σ |
| Dynamical dark energy | 3.2 σ |

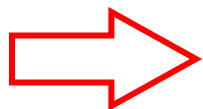
CMB-SPA

CMB-SPA+DESI



Constraints from CMB and BAO data beyond Λ CDM

- With current data, **no definitive evidence for a breakdown of Λ CDM**
 - Evidence is moderate, 3σ -level
 - So far no detection by individual probe
 - Statistical fluctuation or systematic origin not ruled out



More data needed for a stronger judgement (CMB, BAO, others)

Upcoming SPT-3G results

Very soon: Lensing update

- Same observations as in today's results
- $\phi\phi$ from temperature+polarization, quadratic estimator

Soon: Summer

- Ext-4K = Summer + Main (first 2 years)
→ 3 times more sky than in today's results
- TT, TE, EE
- $\phi\phi$

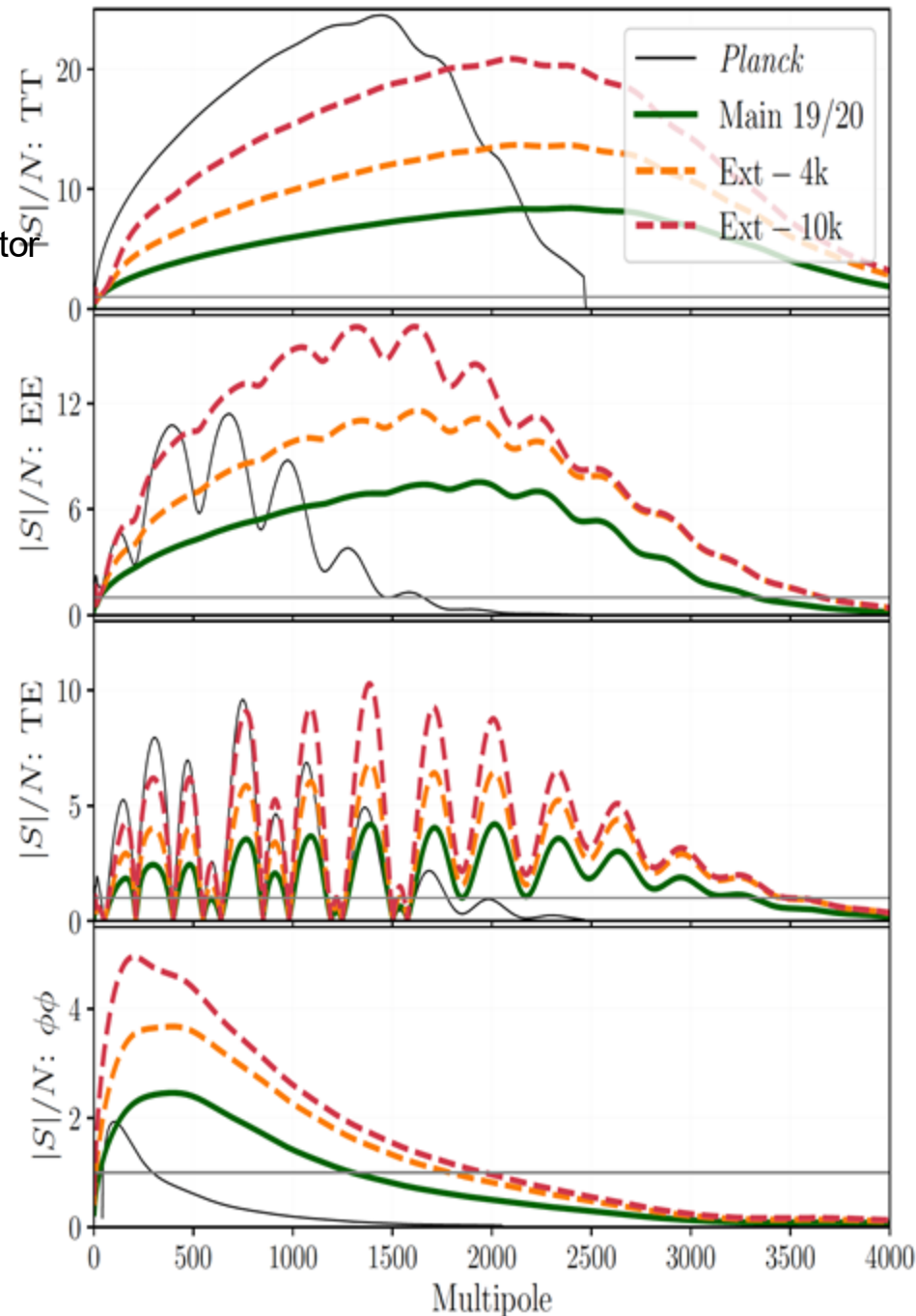
In progress: Wide

- Ext-10K = Wide + Summer + Main
→ 7 times more sky than in today's results
- TT, TE, EE
- $\phi\phi$

Cosmology forecasts: [Prabhu et al. \[SPT-3G\], 2024](#)

Future:

- Main and Summer full depth
- SPT-3G+ camera starting in 2029



Prospects

Soon:

- SPT-3G
- Simons Observatory (SO)

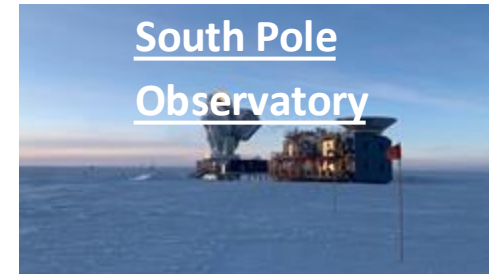
~2028

- SPT-3G+ and its combination with BICEP/Keck, South Pole Observatory (SPO)
- Advanced SO

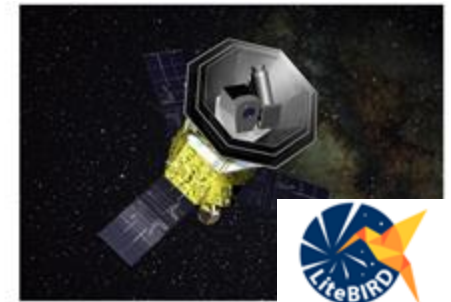
~2035+

- Litebird

Current/upcoming



Future

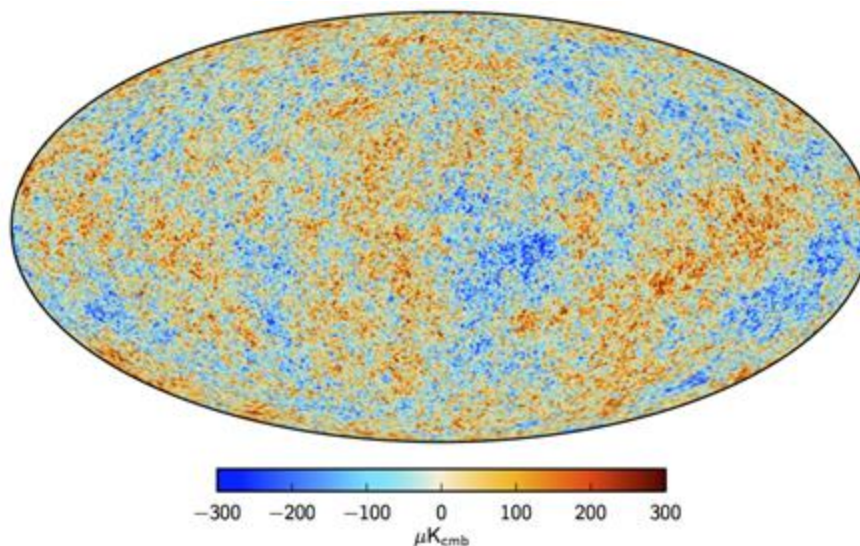


Very recent news of the termination of the S4 project from USA funding agencies. **Change of strategy towards funding upgrades on existing experiments and their combinations, still very strong support to CMB science.**



More details about CMB STATISTICS

- CMB anisotropies are expected to be distributed as a **gaussian random field**.
- We cannot theoretically predict the value of the temperature in the pixels, but only predict their statistical properties.



- A gaussian distribution is fully characterized by a **mean and variance**. All higher odd moments are 0, even moments can be written in terms of the variance (Wick's theorem)

Spherical harmonics

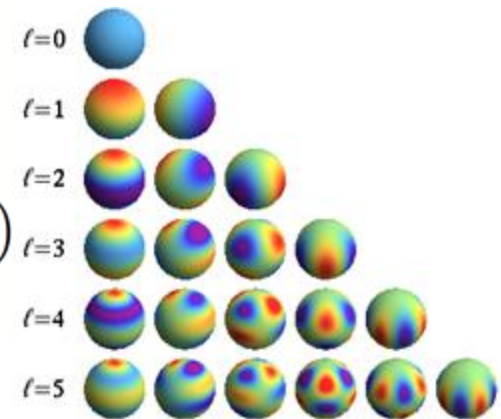
- We can decompose the temperature maps in spherical harmonics.
- SH are a **ortho-normal basis of functions on the sphere**. They are the eigenfunctions of the angular part of the Laplace operator in spherical coordinates.

$$\nabla^2 Y_\ell^m = -[l(l+1)]Y_\ell^m$$

$$\int d\hat{n} Y_\ell^{m*}(\hat{n}) Y_{\ell'}^{m'}(\hat{n}) = \delta_{\ell\ell'} \delta_{mm'}$$

$$\sum_{\ell m} Y_\ell^{m*}(\hat{n}) Y_\ell^m(\hat{n}') = \delta(\phi - \phi') \delta(\cos \theta - \cos \theta')$$

- Complex. Conjugation $Y_\ell^{m*} = (-1)^m Y_\ell^{-m}$



- Characterized **the degree (multipole) ℓ and the order m** .
- $\ell \sim \pi/\theta$** , with the **θ** angular separation in the sky.
- For each ℓ - $\ell \leq m \leq \ell$ There **are $2\ell+1$ m-modes for each ℓ** . For a real field, as the CMB intensity, there are $(2\ell+1)/2$ independent modes because $a_{\ell m} = a_{\ell m}^* (-1)^m a_{\ell -m}$
- The projection on the m-modes depends on the reference system.

Decomposition in spherical harmonics

- Decompose the fractional temperature variation in spherical harmonics

$$\frac{\Delta T}{T}(\hat{n}, \vec{x}, \eta) = \Theta(\hat{n}, \vec{x}, \eta) = \sum_{\ell m} \Theta_{\ell m}(\vec{x}, \eta) Y_{\ell}^m(\hat{n})$$

Line of sight \hat{n}
 Position in the sky ($u, s, x=0$) \vec{x}
 Conformal time ($u, s, \eta=\eta_0$) $\eta \equiv \int dt/a$

Also often called $a_{\ell m}$ in the literature

- Applying the orthogonality of spherical harmonics:

$$\Theta_{\ell m}(\vec{x}, \eta) = \int_{\Omega} d\hat{n} \Theta(\hat{n}, \vec{x}, \eta) Y_{\ell m}^*(\hat{n})$$

- In the simplest models of inflation, $\Theta(\hat{n})$ is a **gaussian random field**. Then, $\Theta_{\ell m}$ are **statistically independent** and **randomly distributed**, each described by a gaussian distribution.

Cl's in theory

- To characterize the statistical properties of a gaussian random field, we can calculate the mean and the variance of the field. For the CMB, the mean of the anisotropies is zero (by definition). The variance can be calculated either as the 2-point correlation function in real space, or equivalently, as the **angular power spectrum in harmonic space**.

$$\langle \Theta_{\ell m} \rangle = 0 \quad \langle \Theta_{\ell m} \Theta_{\ell' m'} \rangle = \delta_{\ell \ell'} \delta_{m m'} C_{\ell}$$

- $\langle \rangle$ are ensemble averages over many realizations of the sky.
- Because of **isotropy**, $\Theta_{\ell m}$ with same ℓ and different m are extracted from gaussian distribution with the same variance C_{ℓ}

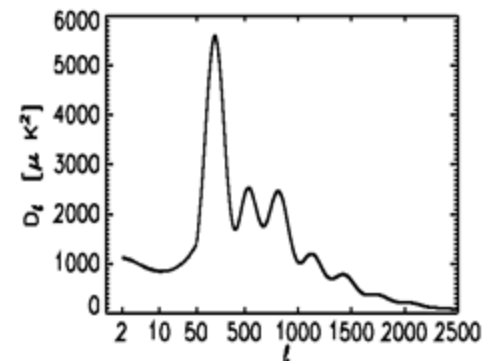
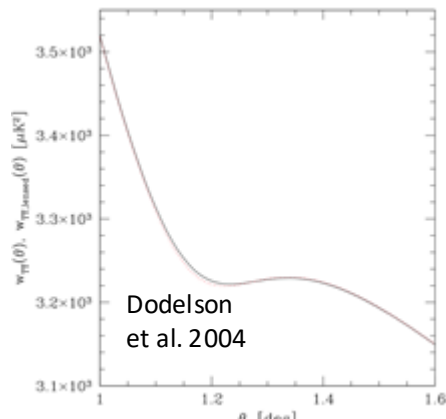
Cl's and 2-point correlation function

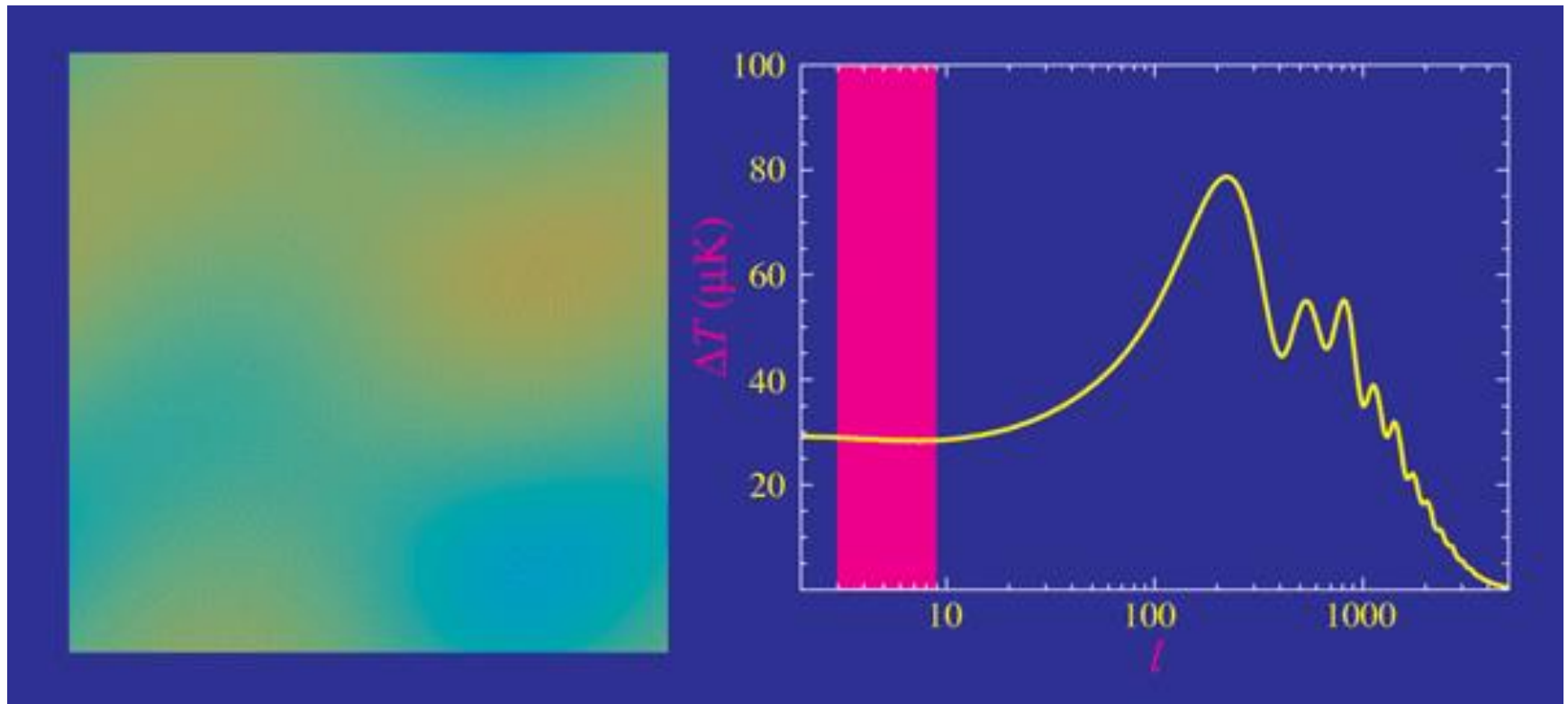
- We can relate the angular power spectrum to the 2-point correlation function in real space using the Legendre polynomials and the addition theorem:

$$\sum_m Y_{\ell m}^*(\mathbf{n}_i) Y_{\ell m}(\mathbf{n}_j) = \frac{2\ell + 1}{4\pi} P_\ell(\hat{\mathbf{n}}_i \cdot \hat{\mathbf{n}}_j)$$

$$\langle \Theta_i \Theta_j \rangle = \sum_\ell \frac{2\ell + 1}{4\pi} C_\ell P_\ell(\hat{\mathbf{n}}_i \cdot \hat{\mathbf{n}}_j)$$

- Because of isotropy, the two-point correlation function depends only on the angular separation in the sky θ , not on the orientation of the separation.





Each of these maps are extracted from gaussian distributions with 0 mean and variance given by the CI in the corresponding pink band. There is an infinite number of possible realizations.

An estimator for the Cl's

$$\langle \Theta_{\ell m} \rangle = 0 \quad \langle \Theta_{\ell m} \Theta_{\ell' m'} \rangle = \delta_{\ell \ell'} \delta_{m m'} C_\ell$$

- We only observe one universe=> average not possible.
- Because of isotropy, all the m-modes $\Theta_{\ell m}$ with the same ℓ have the same theoretical C_ℓ . An estimator of C_ℓ is then:

$$\hat{C}_\ell = \frac{1}{2\ell + 1} \sum_m \Theta_{\ell m} \Theta_{\ell m}^*$$

Cosmic Variance

- The expected value is $\langle \hat{C}_l \rangle = C_l$
- Since we only have $2l+1$ samples for each l , there is an **intrinsic uncertainty**!

$$\hat{C}_l = \frac{1}{2l+1} \sum_m \Theta_{lm} \Theta_{lm}^*$$

$$\begin{aligned} \frac{\sigma_{C_l}^2}{C_l^2} &= \frac{\langle (\hat{C}_l - C_l)(\hat{C}_l - C_l) \rangle}{C_l^2} = \frac{\langle \hat{C}_l \hat{C}_l \rangle - C_l^2}{C_l^2} \\ &= \frac{1}{(2l+1)^2 C_l^2} \langle \sum_{mm'} \Theta_{lm}^* \Theta_{lm} \Theta_{lm'}^* \Theta_{lm'} \rangle - 1 \\ &= \frac{1}{(2l+1)^2} \sum_{mm'} (\delta_{mm'} + \delta_{m-m'}) = \frac{2}{2l+1} \end{aligned}$$

For a gaussian field, Wick's theorem says that any N-point (N even) statistics can be written as a function of the 2-point correlation function

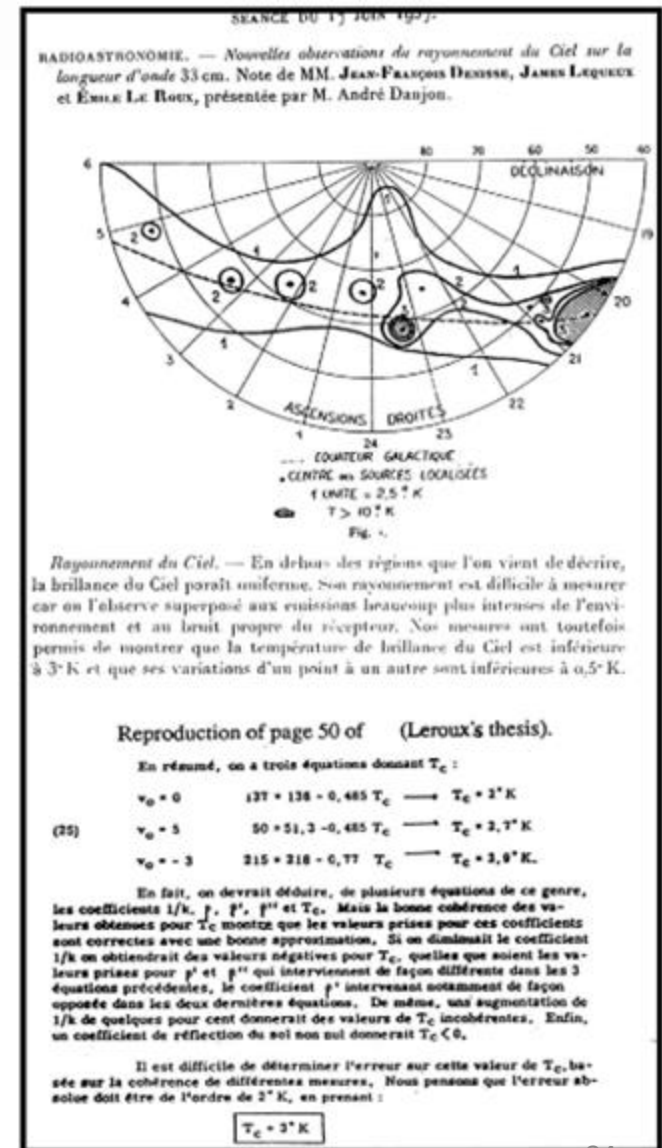
$$\sigma_{C_l}^2 = \frac{2}{2l+1} C_l^2$$



**More details about earlier
detections of the CMB**

Earlier detections?

- 1940 Andrew McKellar observed excited rotational states of CN molecules in interstellar absorption lines. In thermal equilibrium at $T \sim 2.3\text{K}$ (see also W. Adams 1941)
- 1955 Émile Le Roux: survey at $\lambda = 33\text{ cm}$ (Nançay Radio Observatory). Near-isotropic background at $3 \pm 2\text{K}$ (Denisse, Lequeux, Le Roux 1957, Le Roux PhD thesis 1957).





**More details about
polarization**

Stokes parameters

For a monochromatic plane wave:

$$E_x = a_x \cos[\omega_0 t - \theta_x(t)] \quad E_y = a_y \cos[\omega_0 t - \theta_y(t)]$$

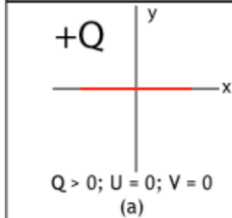
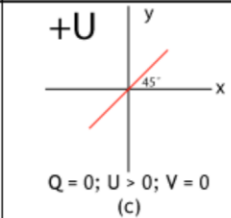
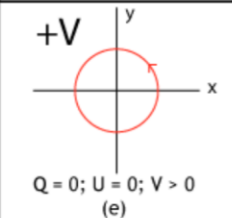
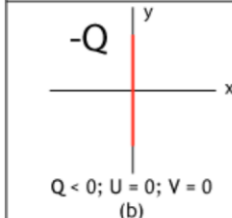
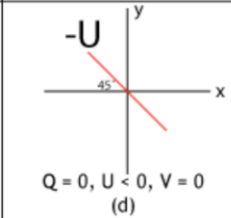
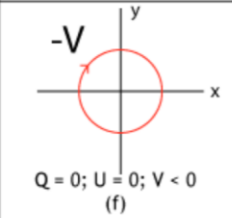


$$\phi = \theta_x - \theta_y$$

$\phi = 0 \Rightarrow$ linear polarization

$\phi = \pi/2, a_x = a_y \Rightarrow$ circular polarization

$$\begin{aligned} I &\equiv \langle a_x^2 \rangle + \langle a_y^2 \rangle \\ Q &\equiv \langle a_x^2 \rangle - \langle a_y^2 \rangle \\ U &\equiv \langle 2a_x a_y \cos(\theta_x - \theta_y) \rangle \\ V &\equiv \langle 2a_x a_y \sin(\theta_x - \theta_y) \rangle \end{aligned}$$

| 100% Q | 100% U | 100% V |
|---|---|---|
|  <p>$Q > 0; U = 0; V = 0$ (a)</p> |  <p>$Q = 0; U > 0; V = 0$ (c)</p> |  <p>$Q = 0; U = 0; V > 0$ (e)</p> |
|  <p>$Q < 0; U = 0; V = 0$ (b)</p> |  <p>$Q = 0; U < 0; V = 0$ (d)</p> |  <p>$Q = 0; U = 0; V < 0$ (f)</p> |

$$I^2 \geq Q^2 + U^2 + V^2$$



$$p \equiv \frac{\sqrt{Q^2 + U^2 + V^2}}{I}$$

- Equal holds for a monochromatic wave or, for superposition of many waves, entirely polarized radiation.
- P = degree of polarization
- In the following, we'll drop V (not produced in standard cosmology model.)

I Q U maps

- Healpix convention: Q and U defined in spherical coordinate system (e_θ, e_ϕ) where e_θ is tangent to the local meridian and directed from North to South, and e_ϕ is tangent to the local parallel, and directed from West to East.

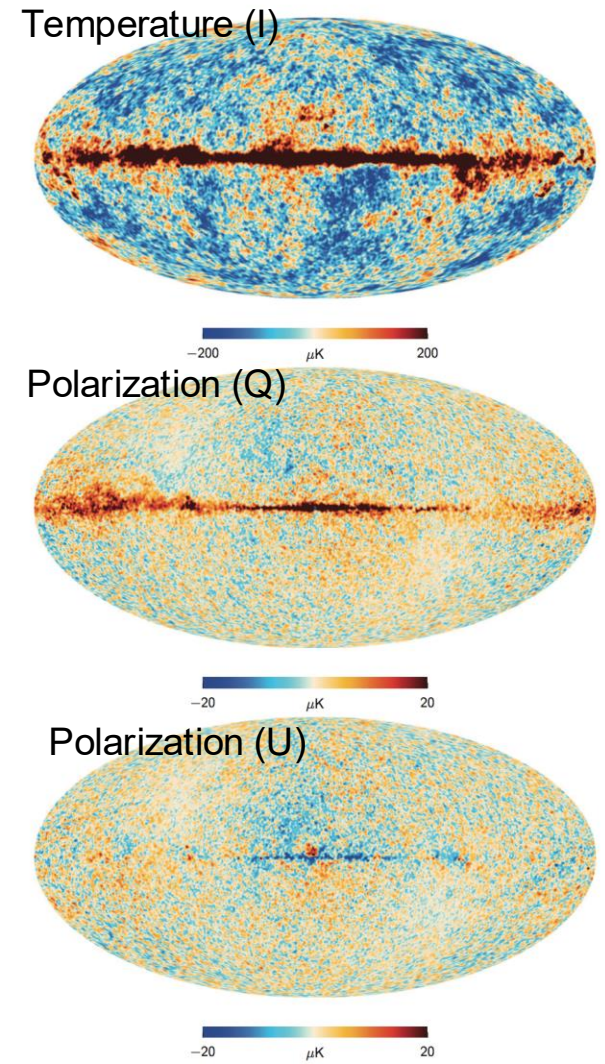


Fig. 12. LFI 70 GHz channel maps. From top to bottom: temperature, Q , and U polarization.

Stokes parameters

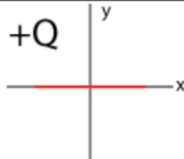
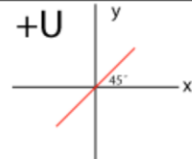
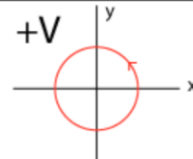
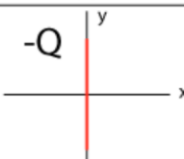
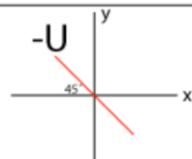
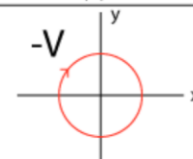
- Polarization is a headless vector, equal to itself after a 180deg rotation=>Q and U spin 2 fields.
- Problem: Q and U depend on reference system.

$$Q' = \cos(2\theta)Q + \sin(2\theta)U$$

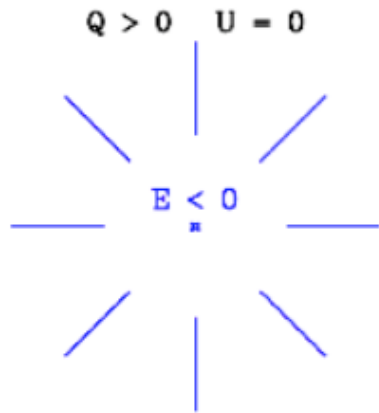
$$U' = -\sin(2\theta)Q + \cos(2\theta)U$$



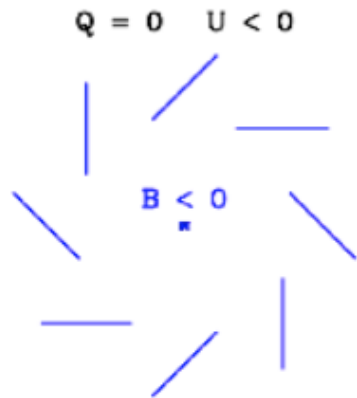
$$Q' \pm iU' = e^{\mp 2i\theta} [Q \pm iU]$$

| 100% Q | 100% U | 100% V |
|--|--|--|
|  <p>$Q > 0; U = 0; V = 0$ (a)</p> |  <p>$Q = 0; U > 0; V = 0$ (c)</p> |  <p>$Q = 0; U = 0; V > 0$ (e)</p> |
|  <p>$Q < 0; U = 0; V = 0$ (b)</p> |  <p>$Q = 0; U < 0; V = 0$ (d)</p> |  <p>$Q = 0; U = 0; V < 0$ (f)</p> |

E and B modes



Q_r and U_r around each point defined in radial coordinate system (e_r, e_t) where e_t is orthogonal to the radius and e_r is parallel



Q and U depend on the reference system \Rightarrow not good to characterize the underlying physics!

A solution is to characterize polarization not by the characteristics in a point, but a **non-local** average 'pattern' around a point the sky.

From Q and U to E and B in a flat patch.

- Take the simplified case of a flat, small patch in the sky. We can define Q and U parameters in the radial coordinates around a given point θ (Zaldarriaga 2001) (\mathbf{e}_r in the radial direction, \mathbf{e}_θ in the orthogonal one).

- E and B are then defined as the C weighted averaged on the full sky**

$$\begin{aligned}
 E(\theta) &= \int d^2\tilde{\theta} \omega(\tilde{\theta}) Q_r(\theta + \tilde{\theta}) \\
 &= \int d^2\tilde{\theta} \omega(\tilde{\theta}) [Q(\theta + \tilde{\theta}) \cos(2\tilde{\psi}) - U(\theta + \tilde{\theta}) \sin(2\tilde{\psi})] \\
 B(\theta) &= \int d^2\tilde{\theta} \omega(\tilde{\theta}) U_r(\theta + \tilde{\theta}) \\
 &= \int d^2\tilde{\theta} \omega(\tilde{\theta}) [Q(\theta + \tilde{\theta}) \sin(2\tilde{\psi}) + U(\theta + \tilde{\theta}) \cos(2\tilde{\psi})]
 \end{aligned}$$

- All the points on this circle contributes with the **same weight** to the construction of E in θ , as they are all at the same distance.

- The value of Q and U depends on the angle ψ , but Q_r is constant on a circle, while $U_r = 0$, so averaging on all the circles, $E \neq 0$, $B = 0$

$$\begin{aligned}
 Q &= I_x - I_y = I_x \neq 0, \\
 U &= I_{45} - I_{-45} = 0 \\
 Q_r &< 0, U_r = 0
 \end{aligned}$$

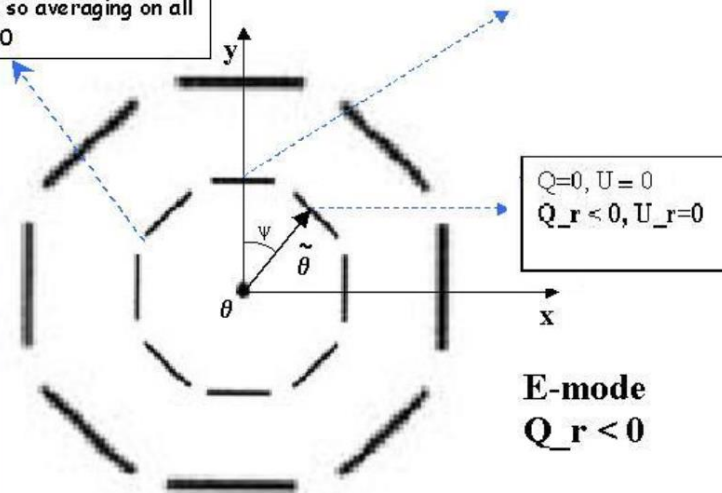
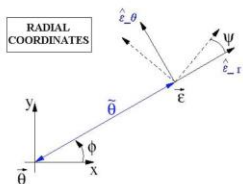
$$\begin{aligned}
 Q &= 0, U = 0 \\
 Q_r &< 0, U_r = 0
 \end{aligned}$$

E-mode
 $Q_r < 0$

In 2D Fourier space, the integrals become multiplications, so that in a flat patch of the sky:

$$\begin{aligned}
 Q(l) &= [E(l) \cos(2\phi_l) - B(l) \sin(2\phi_l)] \\
 U(l) &= [E(l) \sin(2\phi_l) + B(l) \cos(2\phi_l)]
 \end{aligned}$$

Zaldarriaga astro-ph/0106174



E and B on a sphere

- Extend this idea on a sphere (see Zaldarriaga & Seljak 1997, arXiv:astro-ph/9609170).
- 1) Similarly to the (scalar) T maps, we can decompose the Q and U maps (in healpix maps) on the eigenfunctions of the Laplace operator for spin fields, which are an orthonormal, complete basis.

$$(Q \pm iU)(\hat{n}) = \sum_{lm} a_{\pm 2, lm} Y_{lm}(\hat{n})$$

- 2) Obtain scalar quantities by using spin rising/lowering operators (that are rotationally invariant).

$$\delta^2(Q - iU)(\hat{n}) = \sum_{lm} \left[\frac{(l+2)!}{(l-2)!} \right]^{1/2} a_{-2, lm} Y_{lm}(\hat{n})$$

$$a_{E, lm} = -(a_{2, lm} + a_{-2, lm})/2,$$

$$a_{B, lm} = i(a_{2, lm} - a_{-2, lm})/2.$$

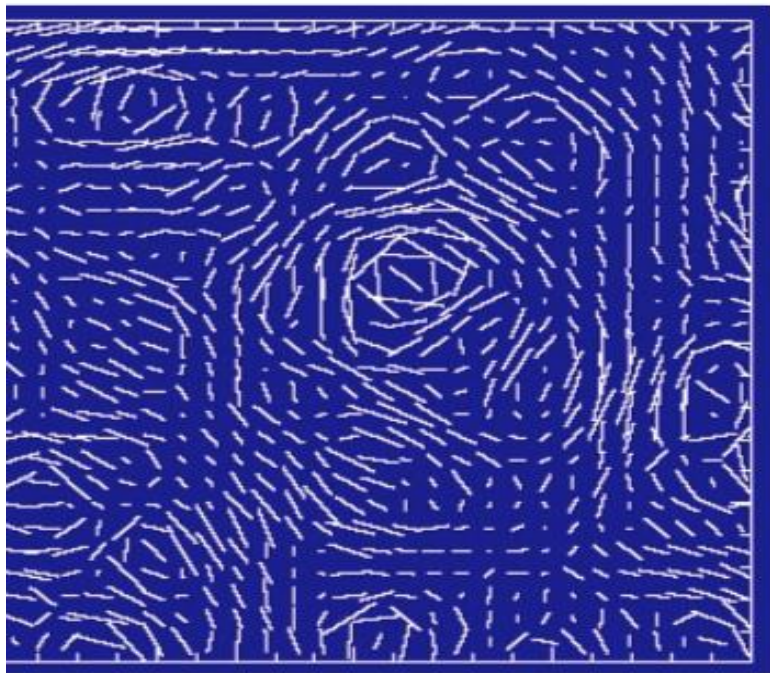
- E and B are rotationally invariant.
- E and B are NON-LOCAL QUANTITIES!
- E remains unchanged under parity transformations (scalar), B changes the sign (pseudo-scalar!)
- In order to obtain (pseudo) scalar fields, need to multiply a_{Elm} by additional geometric factor!

$$\tilde{E}(\hat{n}) \equiv -\frac{1}{2} [\bar{\delta}^2(Q + iU) + \delta^2(Q - iU)]$$

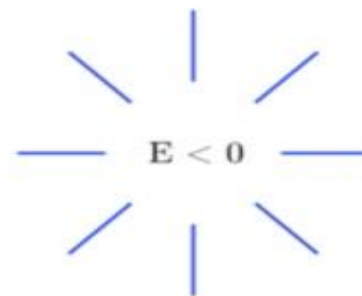
$$= \sum_{lm} \left[\frac{(l+2)!}{(l-2)!} \right]^{1/2} a_{E, lm} Y_{lm}(\hat{n}),$$

$$\tilde{B}(\hat{n}) \equiv \frac{i}{2} [\bar{\delta}^2(Q + iU) - \delta^2(Q - iU)]$$

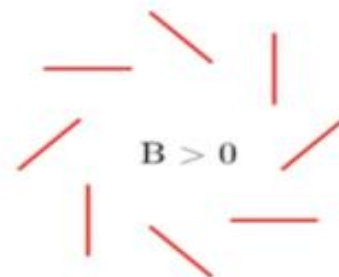
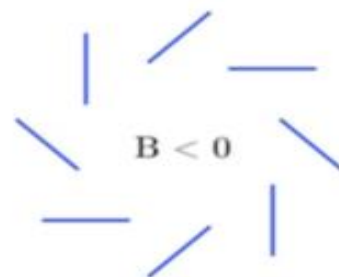
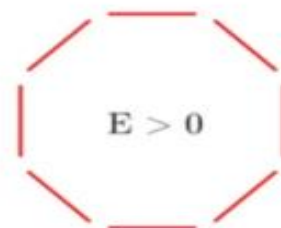
$$= \sum_{lm} \left[\frac{(l+2)!}{(l-2)!} \right]^{1/2} a_{B, lm} Y_{lm}(\hat{n}).$$



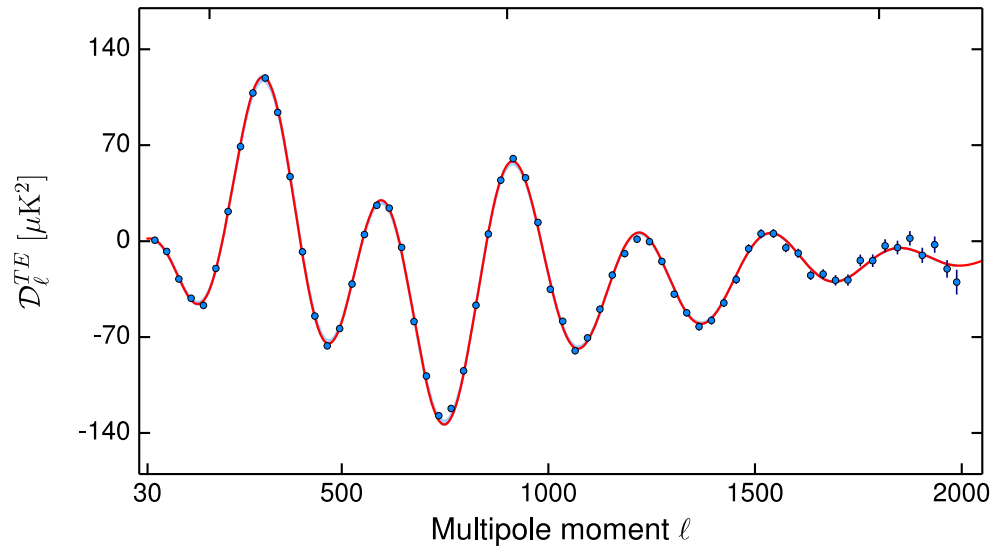
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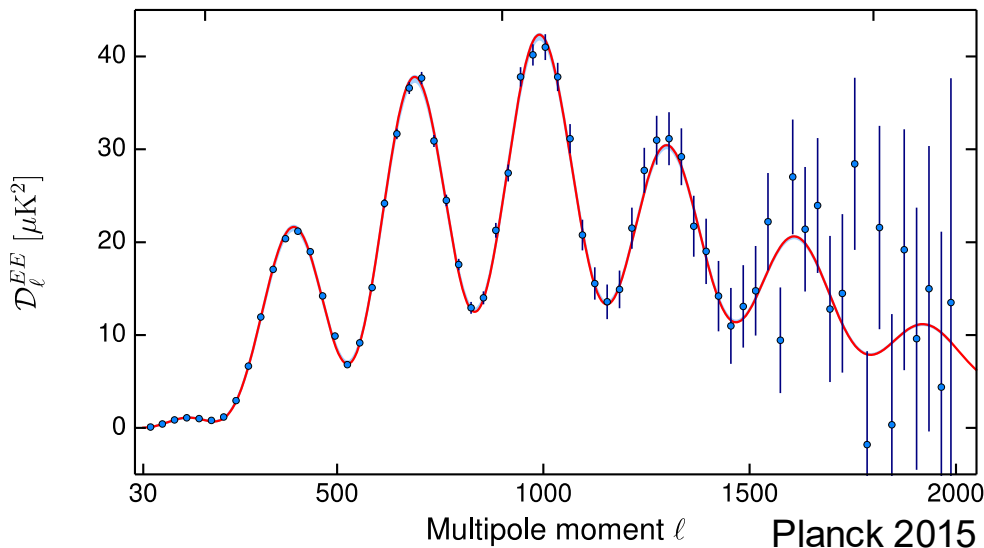


EE and TE power spectra



$$C_{Tl} = \frac{1}{2l+1} \sum_m \langle a_{T,lm}^* a_{T,lm} \rangle$$

$$C_{El} = \frac{1}{2l+1} \sum_m \langle a_{E,lm}^* a_{E,lm} \rangle$$



$$C_{Bl} = \frac{1}{2l+1} \sum_m \langle a_{B,lm}^* a_{B,lm} \rangle$$

$$C_{Cl} = \frac{1}{2l+1} \sum_m \langle a_{C,lm}^* a_{C,lm} \rangle$$

$$D_l = C_l l(l+1)/2\pi$$

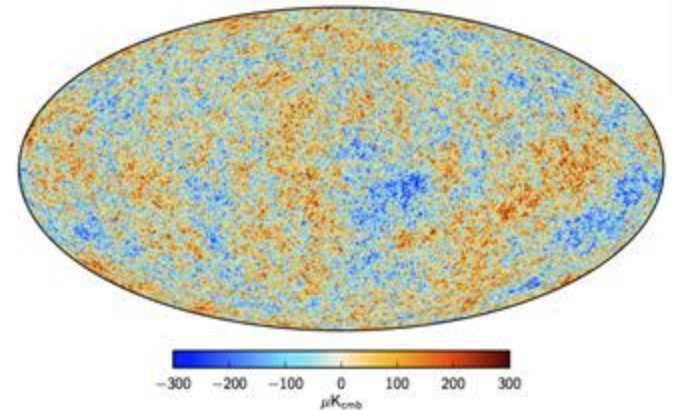
Planck 2015



Miscellanea

The angular power spectrum

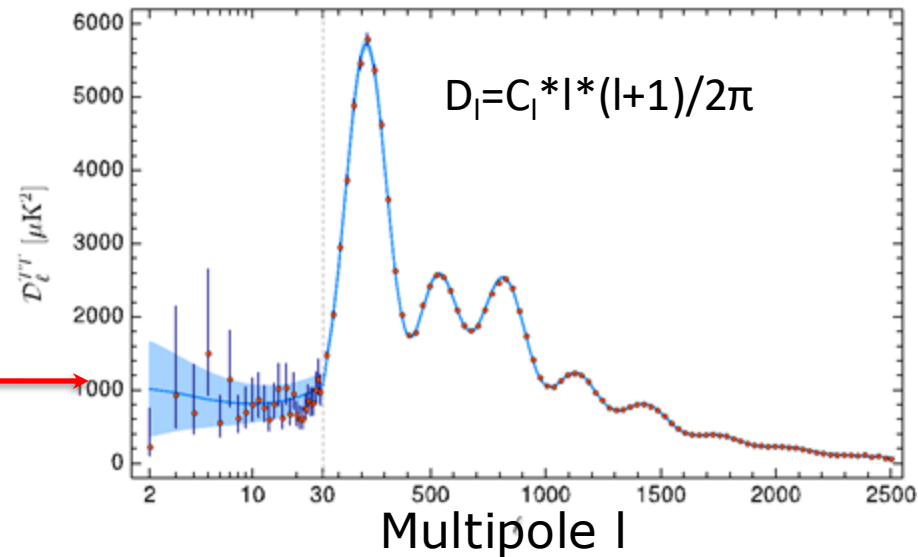
- The anisotropies are distributed as a random **gaussian** field.
- All the information is contained in the two point correlation function, or equivalently in the **angular power spectrum** in harmonic space.



$$\Theta(\vec{x}, \hat{p}, \eta) = \sum_{l=1}^{\infty} \sum_{m=-l}^l a_{lm}(\vec{x}, \eta) Y_{lm}(\hat{p})$$

$$\langle a_{lm} a_{l'm'}^* \rangle = \delta_{ll'} \delta_{mm'} C_l$$

Angular power spectrum



Large scales

Small scales

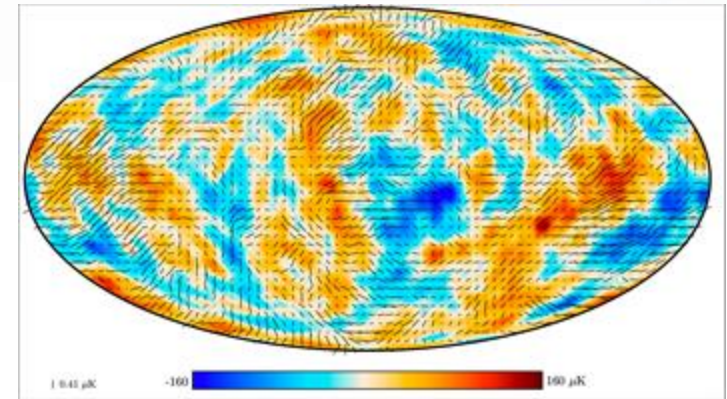
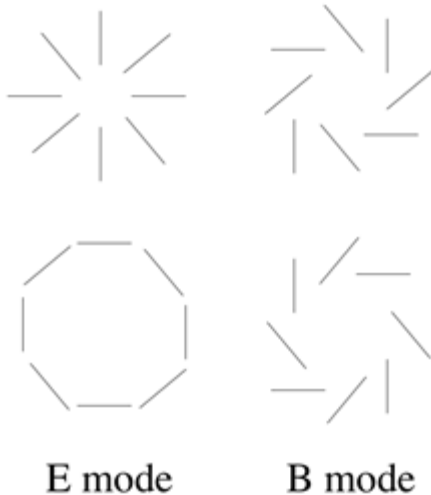
CMB polarization

Polarization generated Thomson scattering in the presence of a temperature quadrupole.

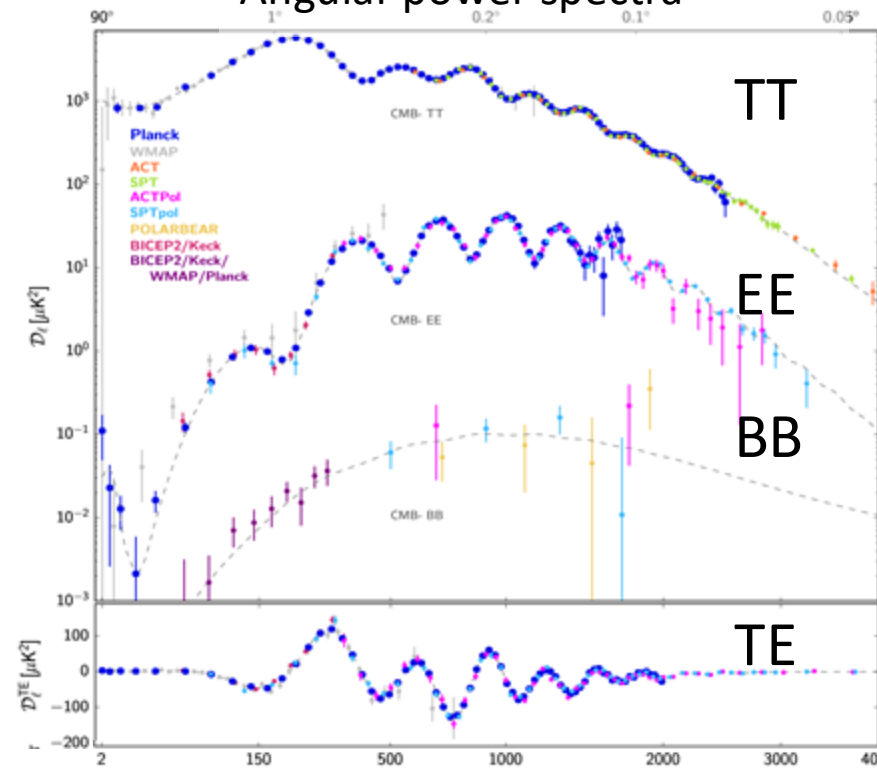
Different sources of quadrupole produce different patterns:

- Scalar (density perturbations): E-mode
- Tensor (e.g. gravitational waves): E-mode and B-mode

Polarization patterns

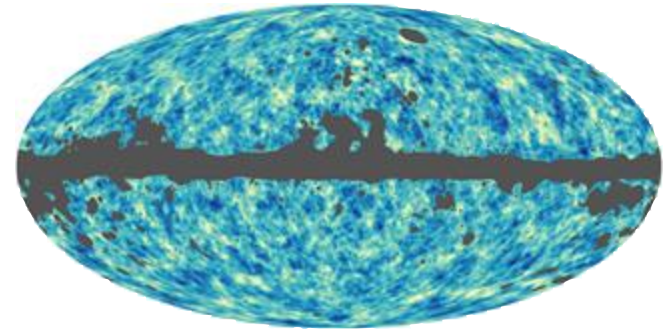
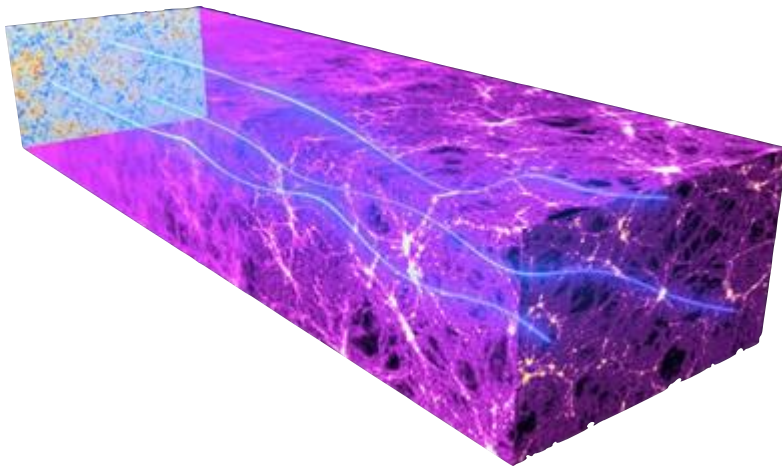


Angular power spectra

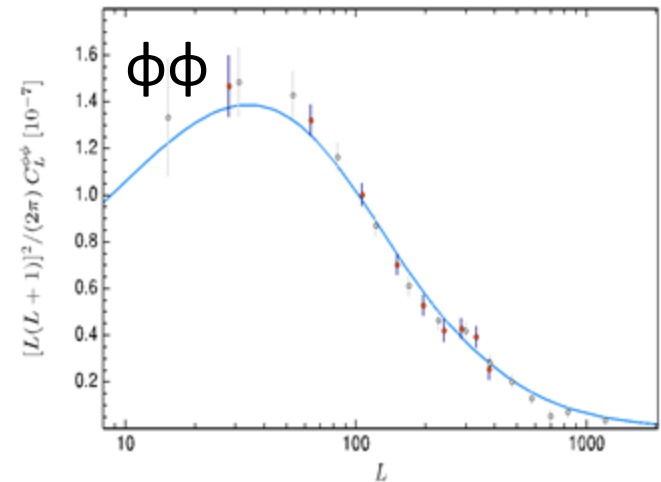


CMB lensing

- CMB lensing breaks isotropy of the CMB.
- Lensing potential map can be extracted from the non-gaussian 4-point correlation function.
- Lensing also impacts the primary power spectra, as well as distorts E-modes into B-modes



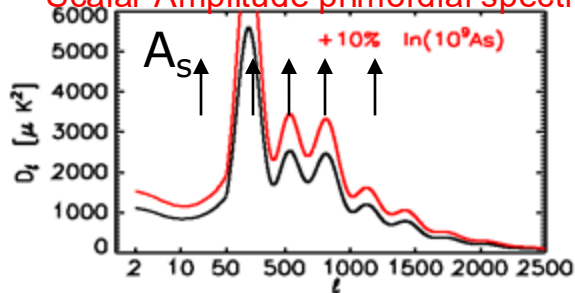
Lensing potential power spectrum



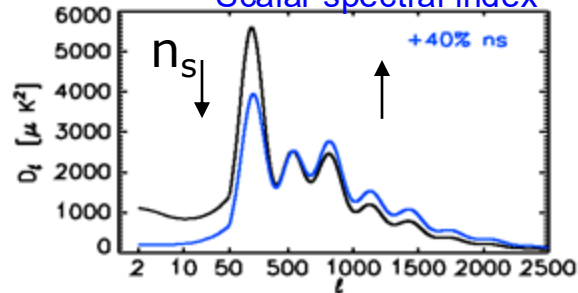
The Λ CDM model

Standard model of cosmology: General relativity to describe gravity, standard model of particles for particle interactions, cosmological constant for dark energy and cold dark matter.

Scalar Amplitude primordial spectrum



Scalar spectral index



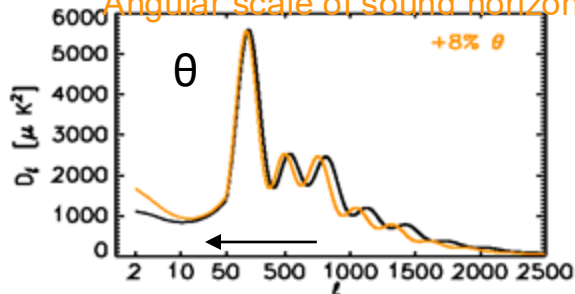
6 parameters:

- Initial conditions A_s , n_s
- Acoustic scale of sound horizon θ
- Reionization τ
- Dark Matter density $\Omega_c h^2$
- Baryon density $\Omega_b h^2$

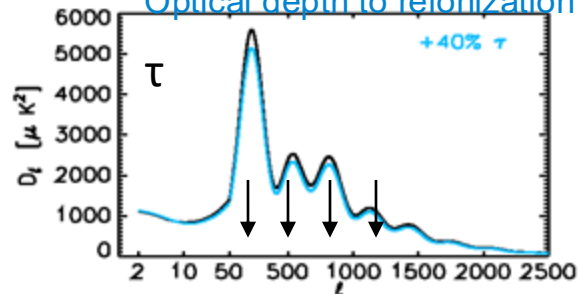
Assumptions:

- Adiabatic initial conditions
- $N_{\text{eff}} = 3.046$
- 1 massive neutrino 0.06eV.
- Tanh reionization ($\Delta z = 0.5$)

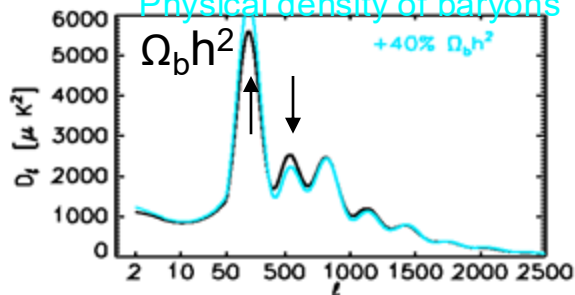
Angular scale of sound horizon



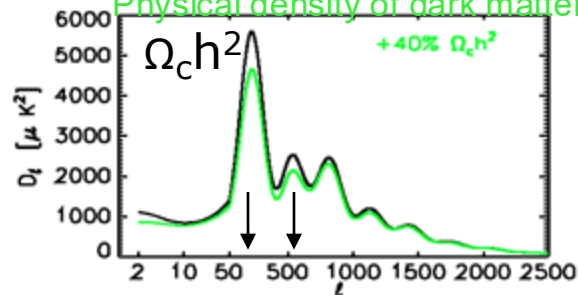
Optical depth to reionization



Physical density of baryons



Physical density of dark matter



The CMB is a laboratory to constrain cosmology and fundamental physics

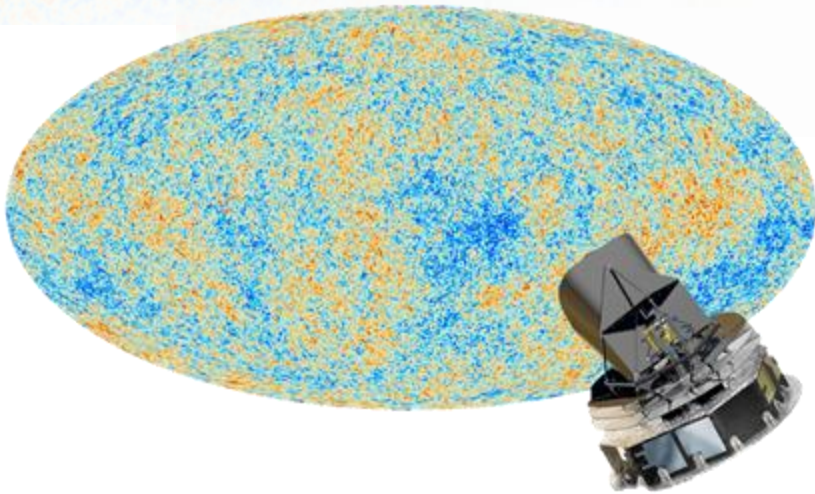


CMB

Planck

SPT-3G

The Planck satellite



3rd generation full sky satellites (COBE, WMAP)
Launched in 2009, operated till 2013.

2 Instruments, 9 frequencies to **disentangle CMB from foregrounds.**

LFI:

- 22 radiometers at **30, 44, 70 Ghz.**

HFI:

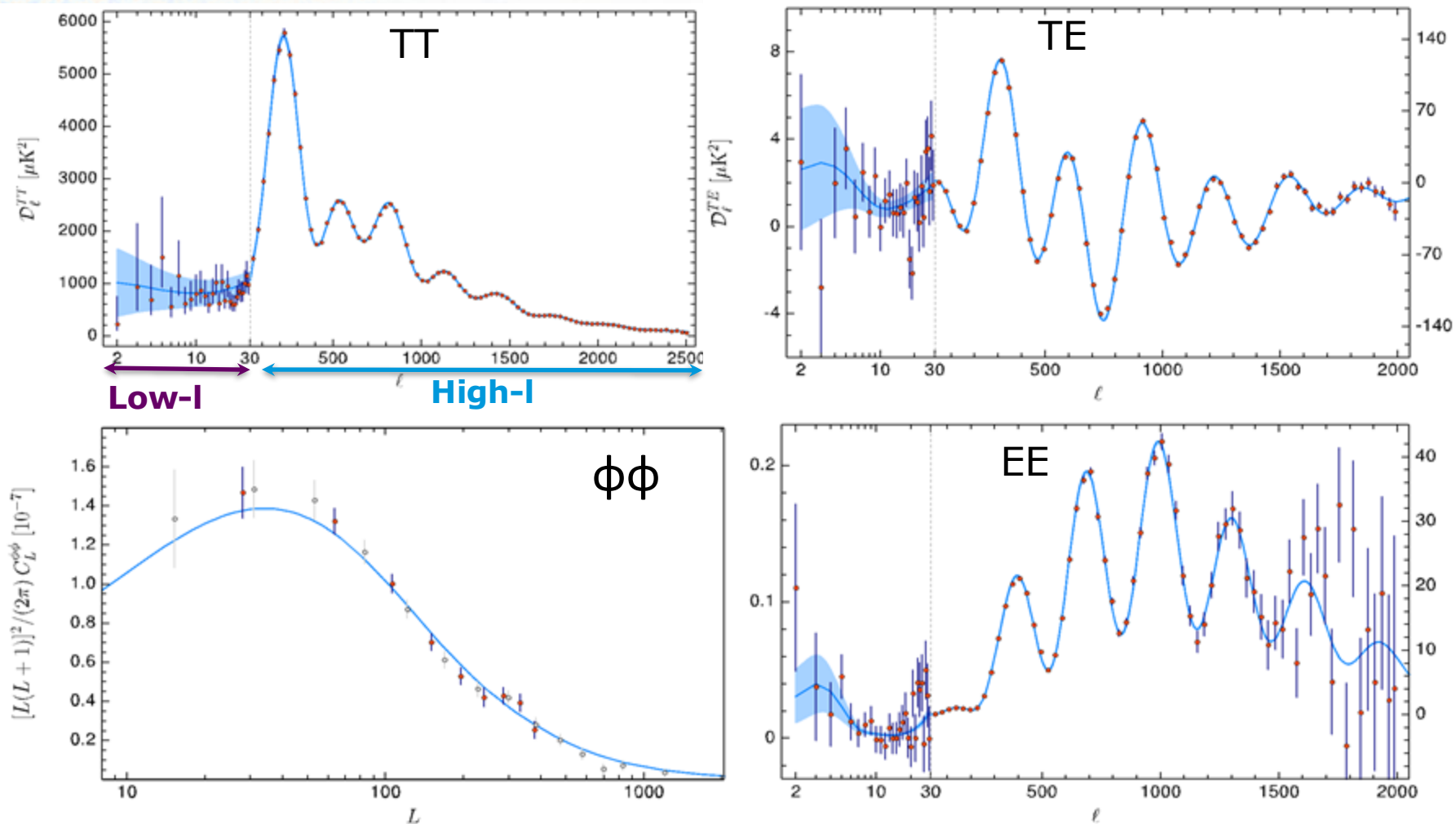
- 50 bolometers (32 polarized) at **100, 143, 217, 353, 545, 857 Ghz.**
- **30-353 Ghz polarized.**

- **1st release 2013: Nominal mission**, 15.5 months, Temperature only (large scale polarization from WMAP).
- **2nd release 2015: Full mission**, 29 months for HFI, 48 months for LFI, Temperature + Polarization, large scale pol. from LFI.
Intermediate results 2016: low- l polarization from HFI
- **3rd release 2018 (PR3): Full mission, improved polarization, low/high- l from HFI.**
Better control of systematics specially in pol., still systematics limited.

Post PR3: new maps PR4 (LFI+HFI map-making, improved low- ell in polarization); new likelihoods and parameters from PR4; new maps Sroll2 (better low- ell polarization); several new estimation of opt. depth to reionization.

No substantially new result compared to PR3.

Planck 2018 power spectra



Planck collaboration 2018 VI.

Λ CDM is an excellent fit to the data. No evidence of preference for classical extensions of Λ CDM from Planck (and non-classical ones: dm annihilation, variation of fund. constants, primordial magnetic fields, variations in recombination, isocurvature, sterile neutrinos, dark energy, modified gravity).....

Baseline Λ CDM results 2018

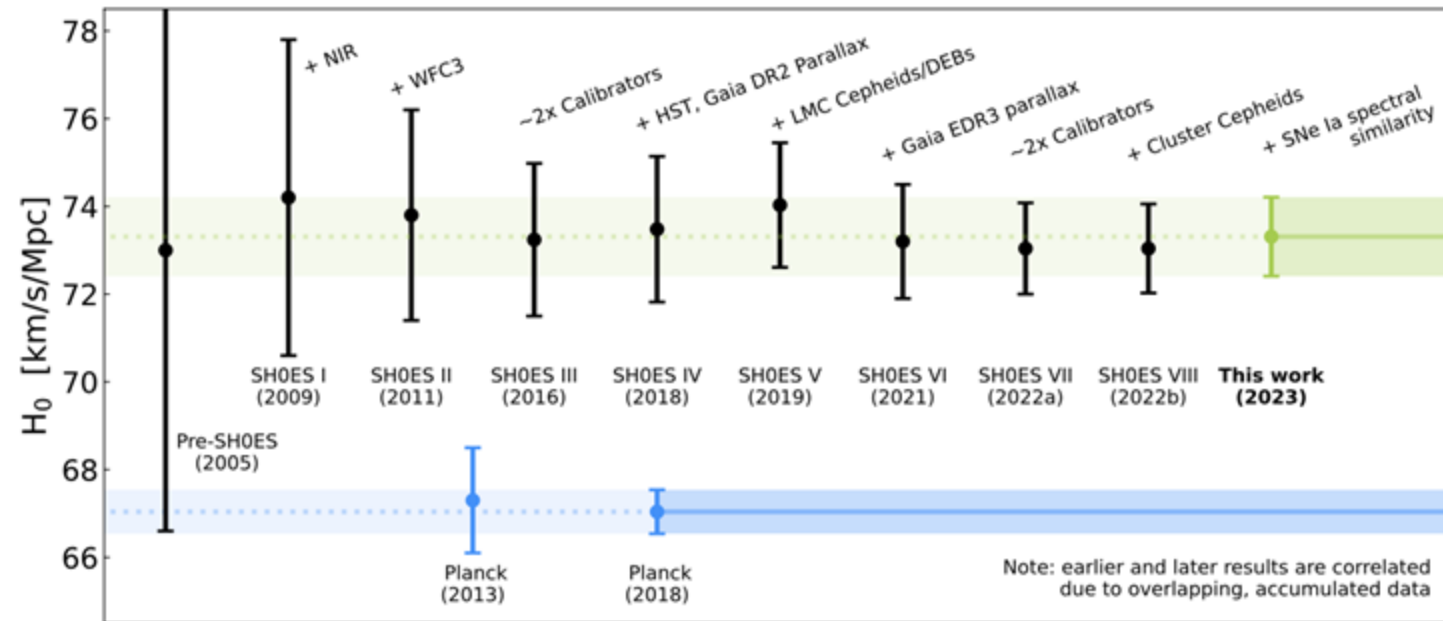
(Temperature+polarization+CMB lensing)

| | Mean | σ | [%] |
|---|---------|----------|------|
| $\Omega_b h^2$ Baryon density | 0.02237 | 0.00015 | 0.7 |
| $\Omega_c h^2$ DM density | 0.1200 | 0.0012 | 1 |
| 100θ Acoustic scale | 1.04092 | 0.00031 | 0.03 |
| τ Reion. Optical depth | 0.0544 | 0.0073 | 13 |
| $\ln(A_s 10^{10})$ Power Spectrum amplitude | 3.044 | 0.014 | 0.7 |
| n_s Scalar spectral index | 0.9649 | 0.0042 | 0.4 |
| H_0 Hubble | 67.36 | 0.54 | 0.8 |
| Ω_m Matter density | 0.3153 | 0.0073 | 2.3 |
| σ_8 Matter perturbation amplitude | 0.8111 | 0.0060 | 0.7 |

- Most of parameters determined at (sub-) percent level!
- **Best** determined parameter is the angular scale of sound horizon θ to **0.03%**.
- τ low and tight, reionization at $z \sim 7$.
- n_s is 8σ away from scale invariance (even in extended models, always $> 3\sigma$)
- **Best (indirect) 0.8% determination of the Hubble** constant to date.

The Hubble tension

The Hubble tension is the difference in the expansion rate of the universe today as measured from Supernovae IA calibrated with Cepheids and from the CMB and other early universe probes. Since the early universe measurements depend on a cosmological model, it could indicate the need of a change in the model, and thus **the discovery of new physics**.



Murakami+ 2023
Breuval+ 2024

5.7 σ tension

Supernovae IA
+ cepheid
SH0ES
 $H_0 = 73.29 \pm 0.90$
km/s/Mpc.

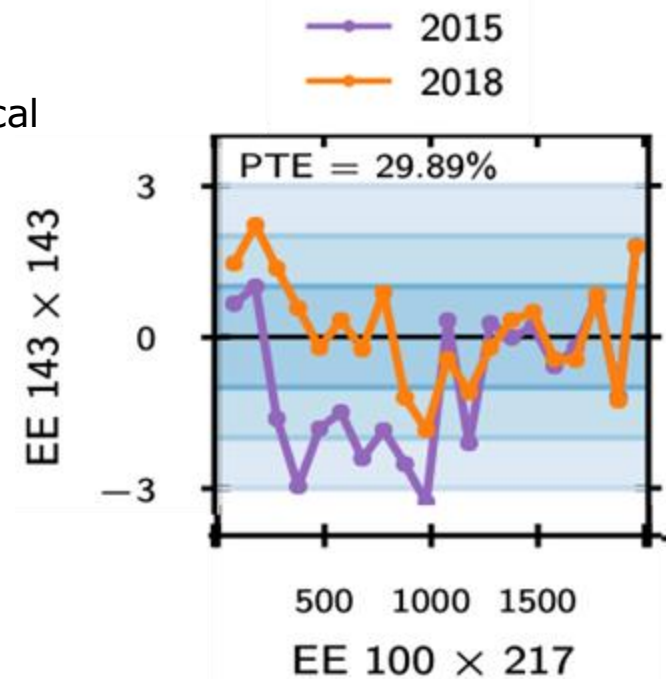
CMB (also BAO)
Planck 2018

$H_0 = 67.36 \pm 0.54$
km/s/Mpc.

The importance of robustness of results

- A **large impact** on the field (tens of thousands of citations)
- **Responsibility to provide the community robust results.**
- Many Tests:
 - **Redundancy** of the data is key in order to be able to do consistency tests at power spectrum or cosmological parameters level from **subsets of the data**, i.e.
 - from different frequency channels, which also corresponds to different detectors.
 - from different map cuts (half mission, versus detector sets)
 - from TT, TE or EE (model dependent).
 - Tested the consistency between a large number of **different analysis choices** on cosmological **parameters** (model dependent).
 - Compared different **analysis pipelines**, which was essential to improve the robustness of the final product.
 - **End-to-end simulations** also allowed us to validate the pipeline.

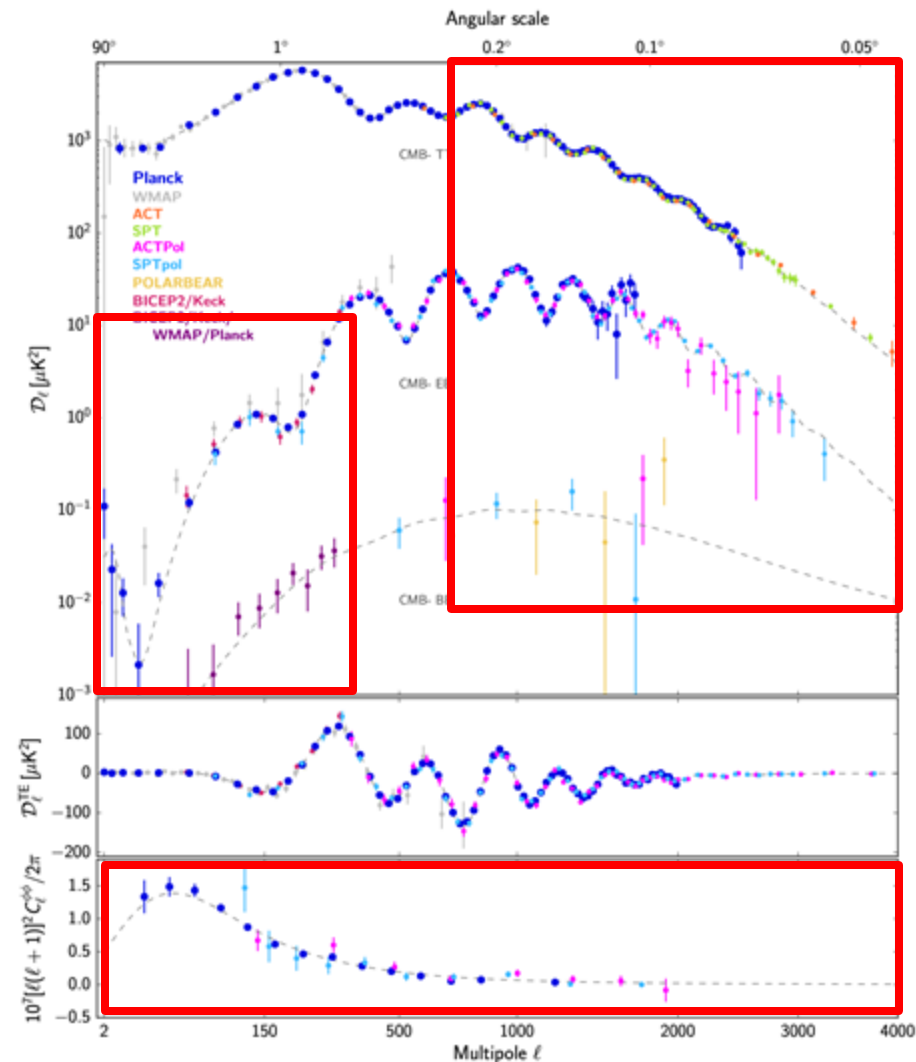
Difference between CMB-only, frequency power spectra in units of error bars in 2015 and 2018



In the second Planck 2015 release, test failed for polarization data. culprit was uncorrected systematics.

From Planck to SPT-3G

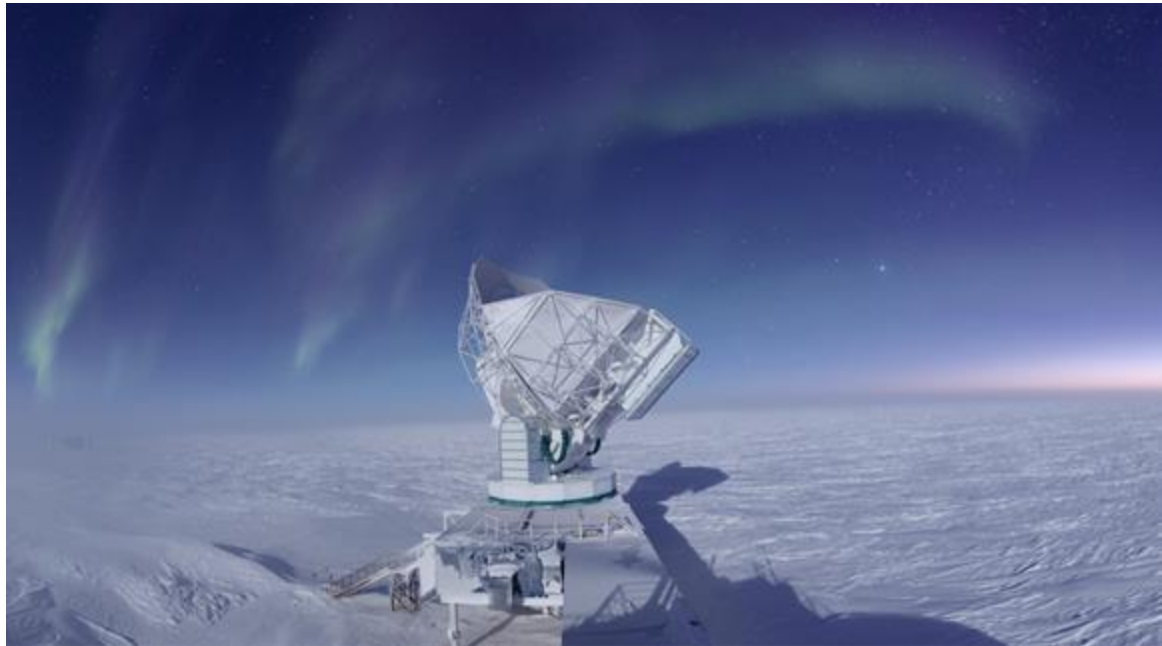
- Planck results had a tremendous impact on cosmology. They confirmed the Λ CDM model. They were checked with many consistency tests.
- Planck opened new mysteries, such as the Hubble tension, that we will explore with upcoming and future experiments.
- There is still a very large amount of information in CMB to be uncovered. Upcoming experiments have two main goals:
 - at large multipoles in polarization, to detect primordial gravitational waves and measure reionization.
 - At small angular scales in polarization, to test cosmological models and the properties of the energy content of the universe.



CMB

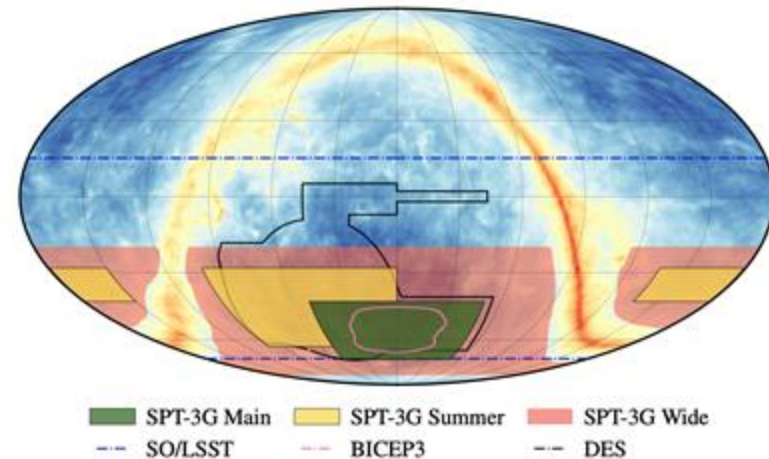
Planck

SPT-3G



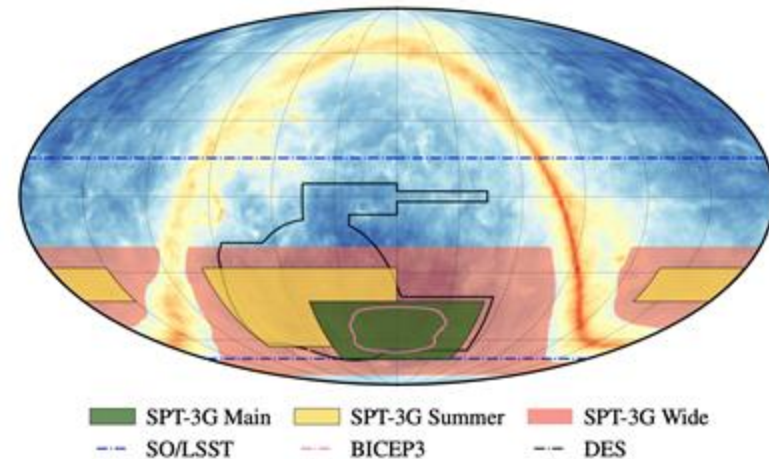
The South Pole Telescope

- SPT is a 10m telescope at the South Pole. It is observing with its third generation camera, SPT-3G, with $\sim 16,000$ detectors at 95, 150, and 220 GHz at high resolution (~ 1 arcmin).
- It has observed **25%** of the sky:
 - **Main** Winter Field **1500 deg²** (5yrs done, 2yr TBD)
Coadded noise **$\sim 17X$** better than Planck.
 - **Summer** fields **2650 deg²** (4yrs done),
coadded noise **4X** better than Planck.
 - **Wide** field **6000 deg²** (1yr)
coadded noise **3X** better than Planck.
- Many scientific goals:
 - **Cosmological constraints from CMB primary anisotropies and CMB lensing**
 - Delensing of the BICEP/Keck field to improve constraints on tensor to scalar ratio r .
 - High-ell TT foregrounds (including kSZ), Cross-correlations with other surveys, High- z galaxies, Clusters of galaxies, Transients etc...



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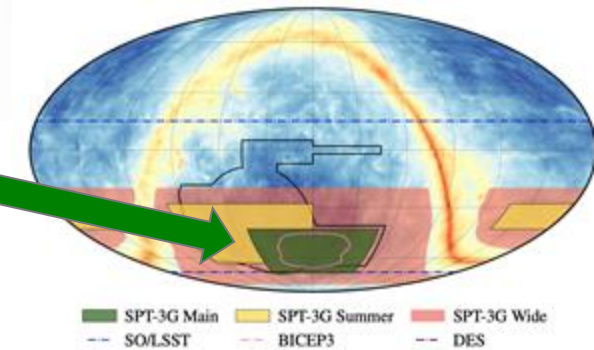
Forecasts show we will measure parameters better than Planck by a factor of 2, and in combination with it by a factor of 3.

SPT-3G have the potential to take the lead on our understanding of the universe in the next years.

Ensuring the robustness of the results is even more critical.

Published power spectrum SPT-3G results from Main field

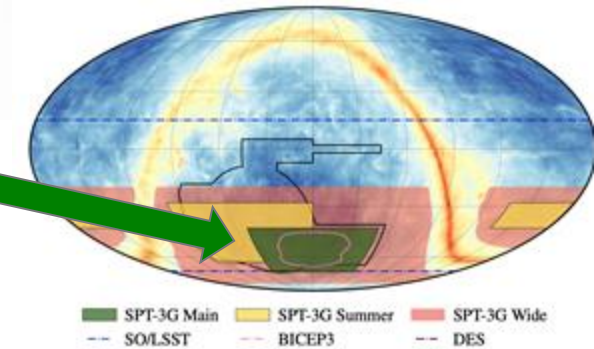
Results based on the observations of the small but **deep Main field** observed during the Austral winter.



- **4 months of observations "SPT-3G 2018"**: Used only half of the focal plane [published in 2021-2023]:
 - Analysis of **TEEE** (Dutcher,...SG+ 2021, Balkenhol,...SG+ 2021)
 - Analysis of **TTTEE** (Balkenhol,...SG+ 2023)
- **2 years of observations "SPT-3G 2 year Main field"**: Based on ~16 months of observations with the full focal plane.
 - Cosmology from CMB lensing and delensed EE power from polarization with **MUSE** (Ge., Millea,...SG+ 2024)
 - Constraints on Inflationary Gravitational Waves from large scale **BB** (Zebrowski,...SG+ 2025)
 - Constraints on duration of reionization from non-gaussianity of KSZ (Raghunathan+ 2024)

Published power spectrum SPT-3G results from Main field

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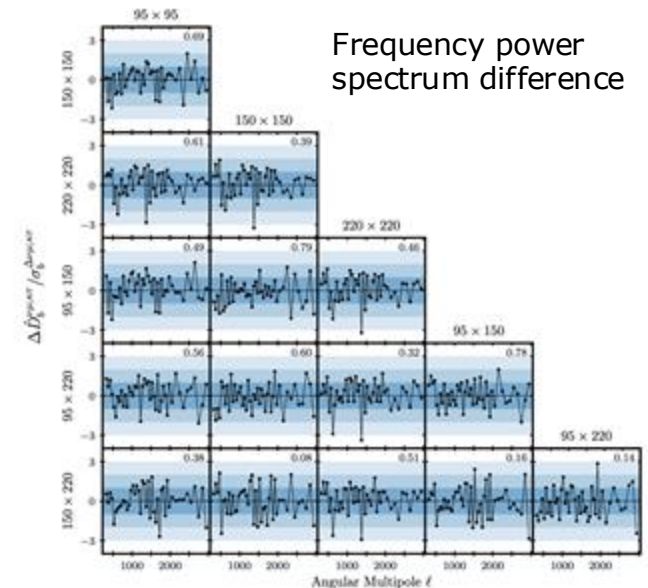
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Early Data Release: SPT-3G TTTEEE 2018

- Data taken from **first 4 months** of data in **2018** from **half of the focal plane**, only from **winter main field**. First analysis only analysed **TEEE** (Dutcher,...SG+ 2021, Balkenhol,...SG et al. 2021).
- We added **TT** in a new analysis in Balkenhol,...SG+ 2023.
 - More **redundancy for tests**.
 - breaks degeneracies, specially for extensions of LCDM model e.g. primordial magnetic fields, EDE (Galli et al. 2022, Smith,... SG et al. 2022).
- The new **TTTEEE** analysis allowed us to introduce many **improvements** to increase the robustness of the results:
 - Developed **consistency** tests inspired on Planck
 - Implemented **blinding**
 - **Speed-up** MCMC chains with ML emulator of Boltzmann code (CosmoPower)
 - **Increased accuracy** of the power spectrum covariance matrix (based on simulations and in flat sky approximation, large off-diagonals terms tricky to model)

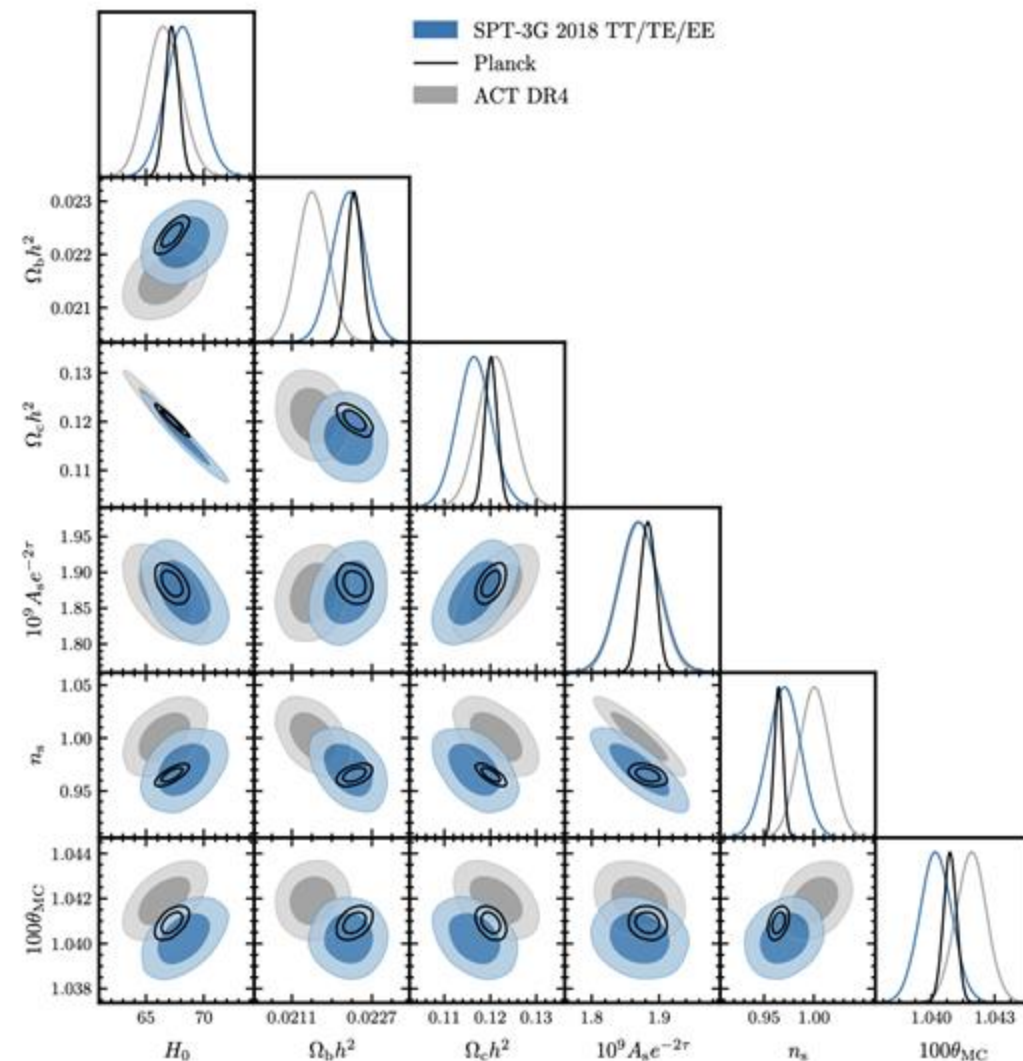


Lennart
Balkenhol



L. Balkenhol, D. Dutcher, A. Spurio Mancini, A. Doussot, K. Benabed, **SG** and the SPT collaboration arXiv:2212.05642

Results from SPT-3G TTTEEE 2018



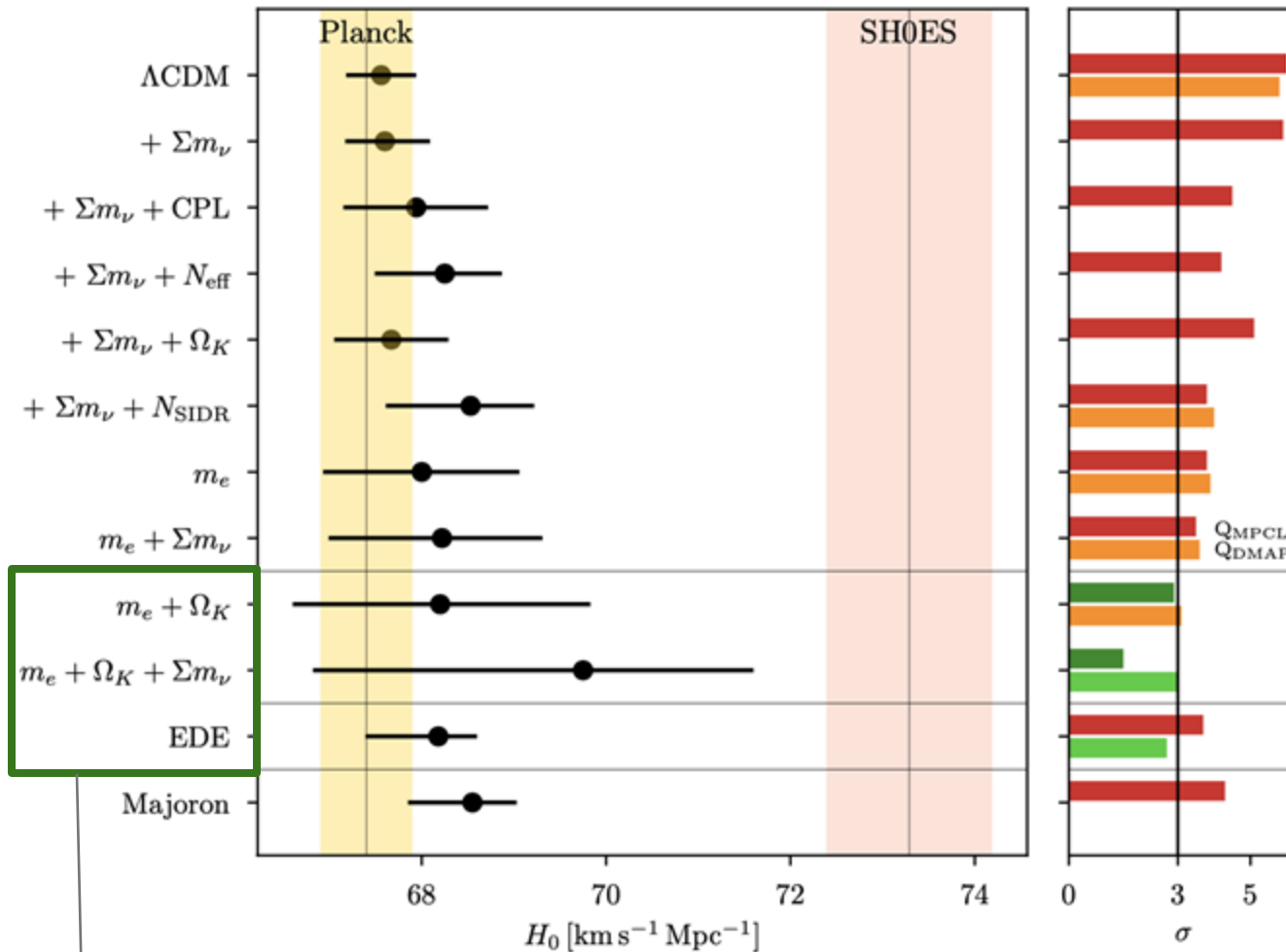
- Errors from SPT3G were three times larger than Planck.
- Λ CDM is a good fit to the data, $\chi^2=763$ for 728 bandpowers (PTE=15%).
- Under Λ CDM, on the 5 parameters, Planck and SPT are consistent with PTE 76%.
- Results are consistent between ACT-DR4 and SPT3G-2018. Constraining power is comparable.

$$H_0 = 68.3 \pm 1.5 \text{ km s}^{-1} \text{ Mpc}^{-1}$$

$$S_8 = 0.797 \pm 0.042$$

No deviations from LCDM

Hubble tension models with SPT-3G 2018 TTTEE



Ali Rida Khalife

Data:

Planck

(TT,TE,EE+lowE+lensing),

SPT-3G 2018 (TT+TE+EE),

BAO (6dFGS+DR7+DR12+DR16),

and **SN IA**

A few still viable models to solve the tension (others in the literature...)

First analysis of polarization from Main 2 years: MUSE

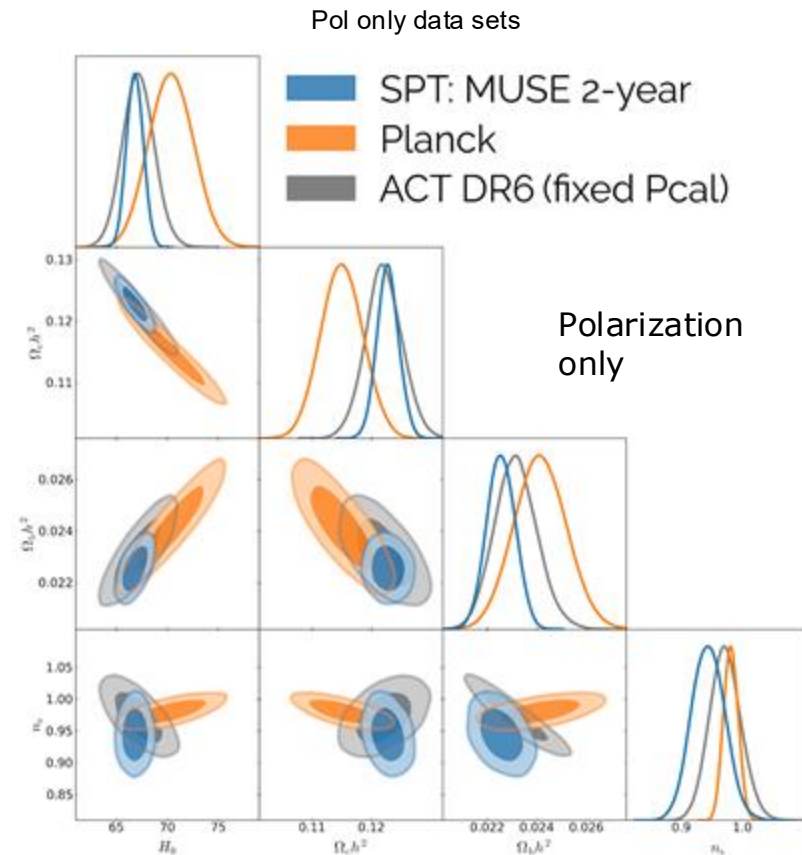
Fei Ge



Marius Millea



- **MUSE**: Algorithm similar to simulation-based inference with semi-analytic compression statistic.
- Used only polarization data to estimate the CMB **unlensed EE power spectrum and lensing reconstruction power spectrum**.
- **Tightest bandpower** measurement of $\varphi\varphi$ at $L>350$ and **EE at $l>2000$**
- Tightest constraints on **Λ CDM parameters from CMB polarization-only inference EE**

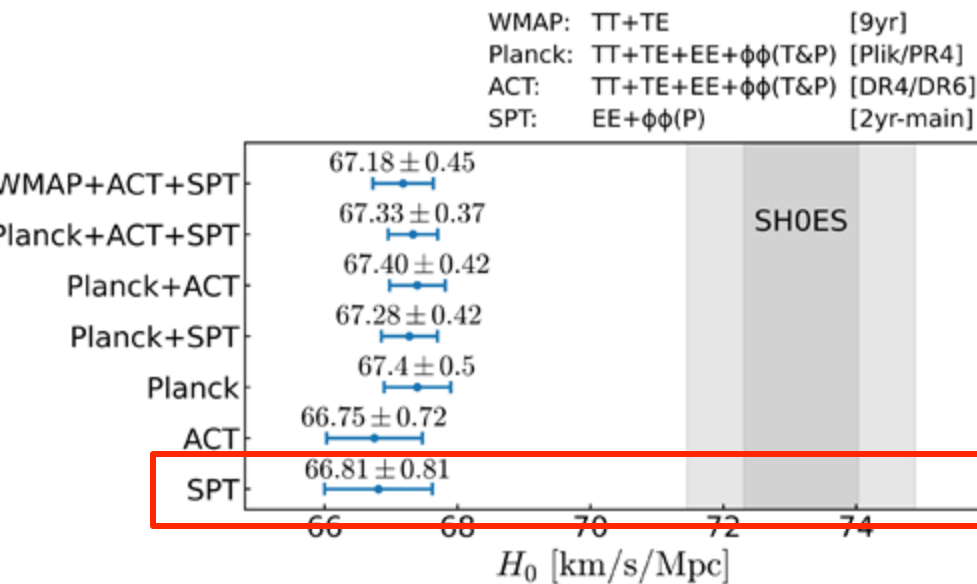


Likelihood public on our website

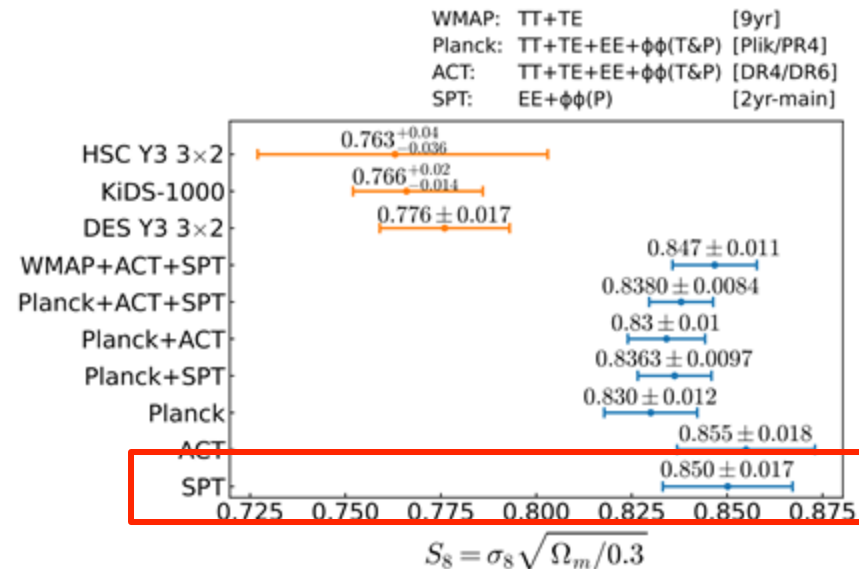
<https://pole.uchicago.edu/public/data/ge25/index.html>

Ge, Millea...SG+ 2025

First analysis of polarization from Main 2 years: MUSE

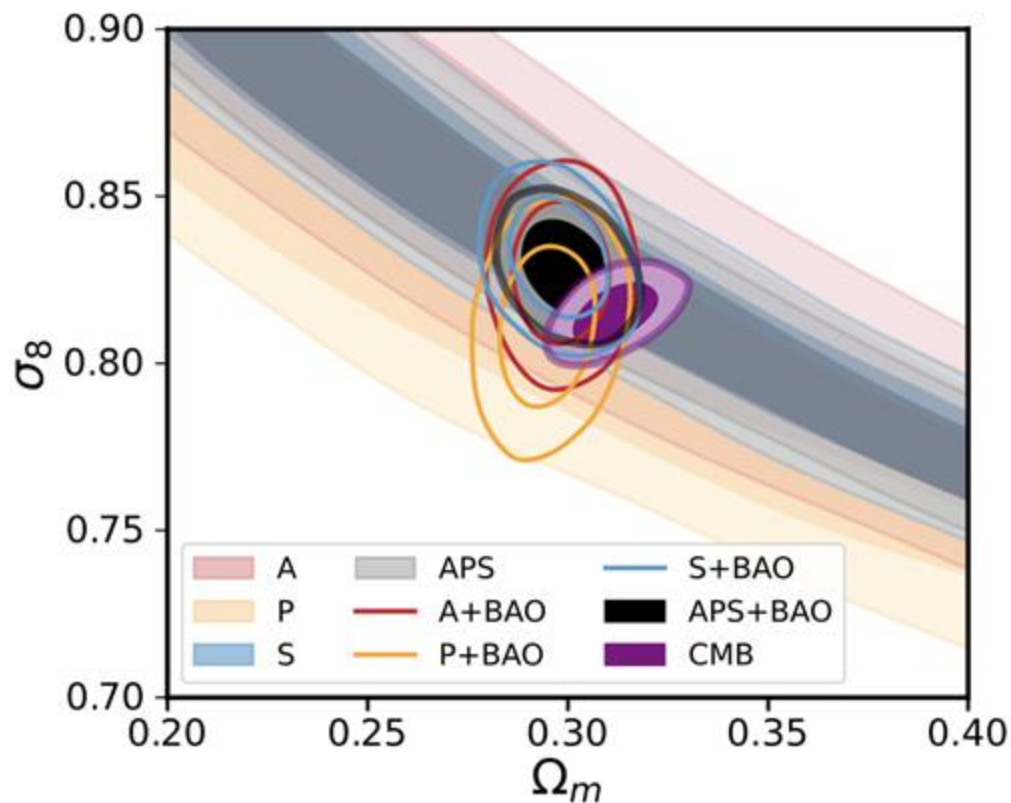


Hubble tension with SH0ES confirmed at 5.4σ from SPT-3G alone.



S_8 results in agreement with other CMB experiments.

ACT Planck SPT lensing combined



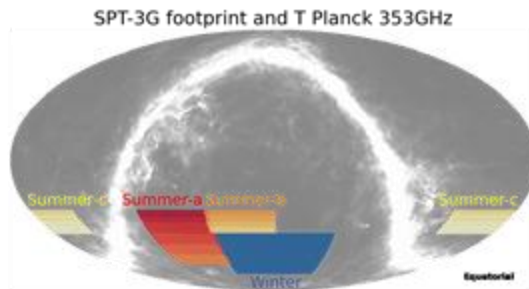
$$S_8^{\text{CMBL}} = 0.825^{+0.015}_{-0.013} \text{ (68\% C.L., APS).}$$

1.6% measurement from CMB lensing alone

Coming next: Summer and Wide fields



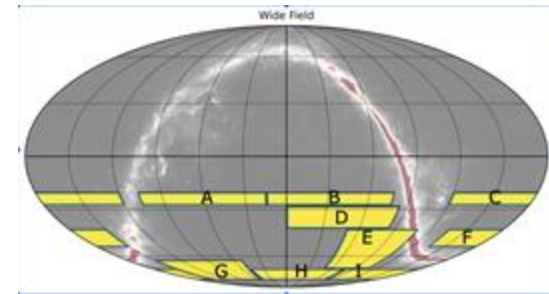
Federica Guidi



- Challenges: Larger atmospheric fluctuations, galactic contamination, lower elevation.
- Mature analysis, to be in a few months (Guidi,...SG et al., in preparation).



Aline Vitrier

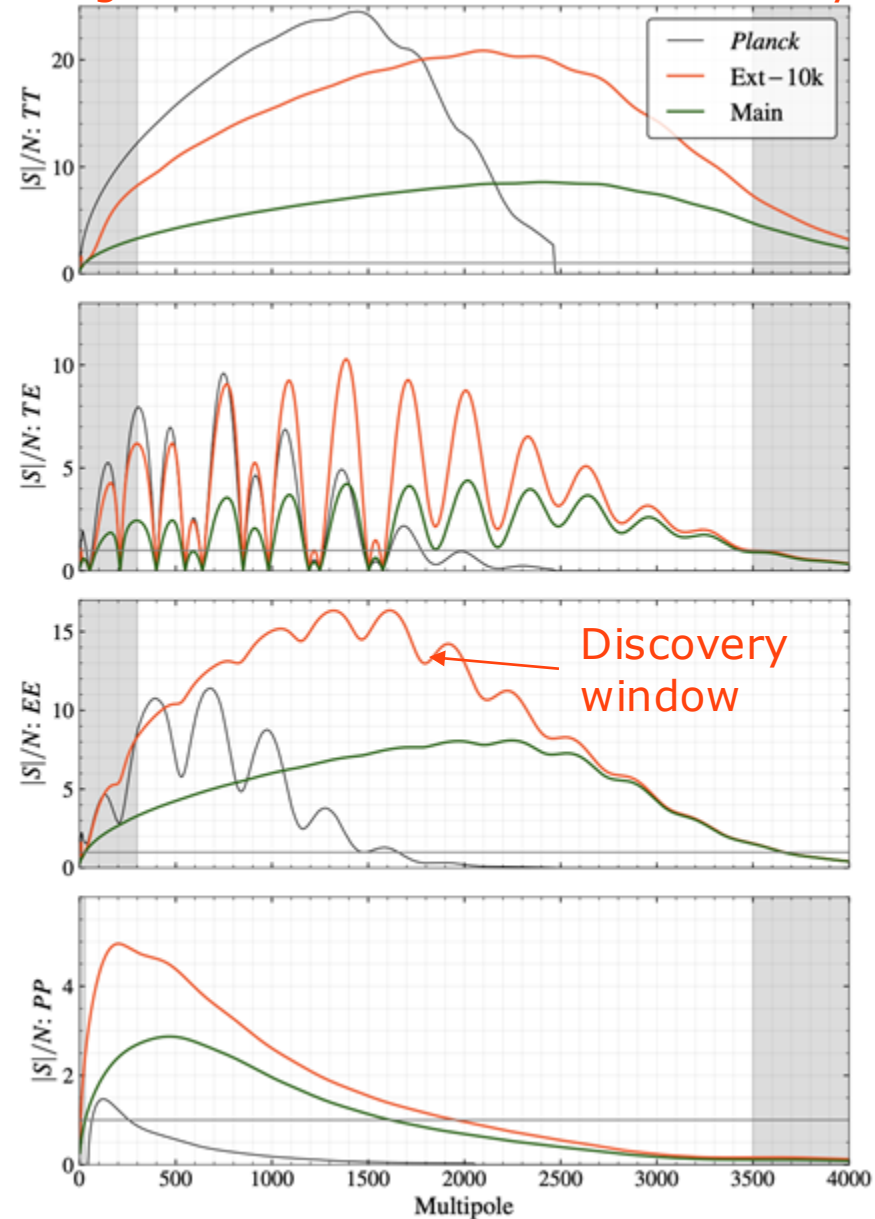


- Challenge of analyzing 9 different subfields, each with different characteristics.
- Early stages of analysis, to be released within a couple of years (Vitrier,...SG et al., in preparation)

Conclusions and prospects

1. CMB is a leading probe of cosmology.
2. SPT-3G will improve over Planck results. We have the potential to constrain proposed solutions to the cosmological tensions. This power will come with the great responsibility of checking for consistency to ensure robustness.
3. Total dataset will be 7+ years of winter field, 4 of summer fields, 1 of wide fields.
4. SPT-3G will be replaced by SPT-3G+ in 2027.
5. The lessons we are learning will pave the way for the next generation of CMB experiments.

Signal-to-Noise of full SPT-3G survey

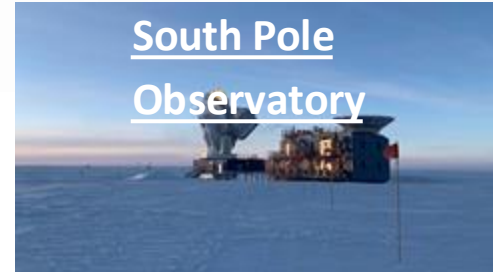


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Current/upcoming

South Pole
Observatory



Future

