

The Early Universe

Les Houches summer school July, 28th & 29th 2025







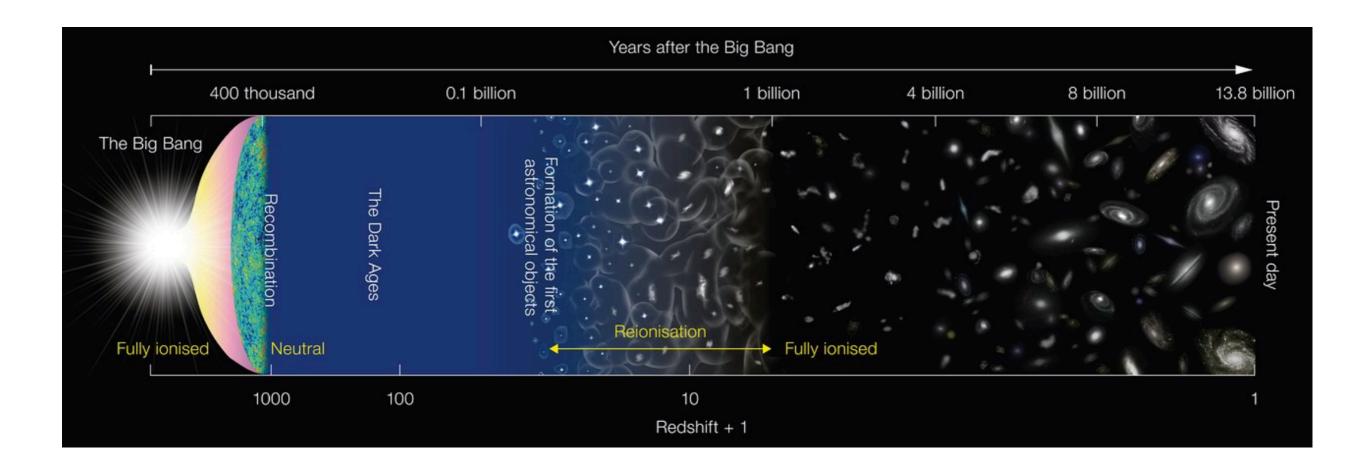


#### Astonishing success of \CDM Cosmology: GR+ Cosmological Principle

$$\omega \equiv \Omega h^2$$
,  $H_0 = 100h \text{ km/s/Mpc}$   $\{H_0, \omega_h, \omega \}$ 

$$\{H_0, \omega_b, \omega_{\text{cdm}}, A_s, n_s, \tau_{\text{reio}}\}$$

$$\Omega_{\Lambda} = 1 - \Omega_m$$



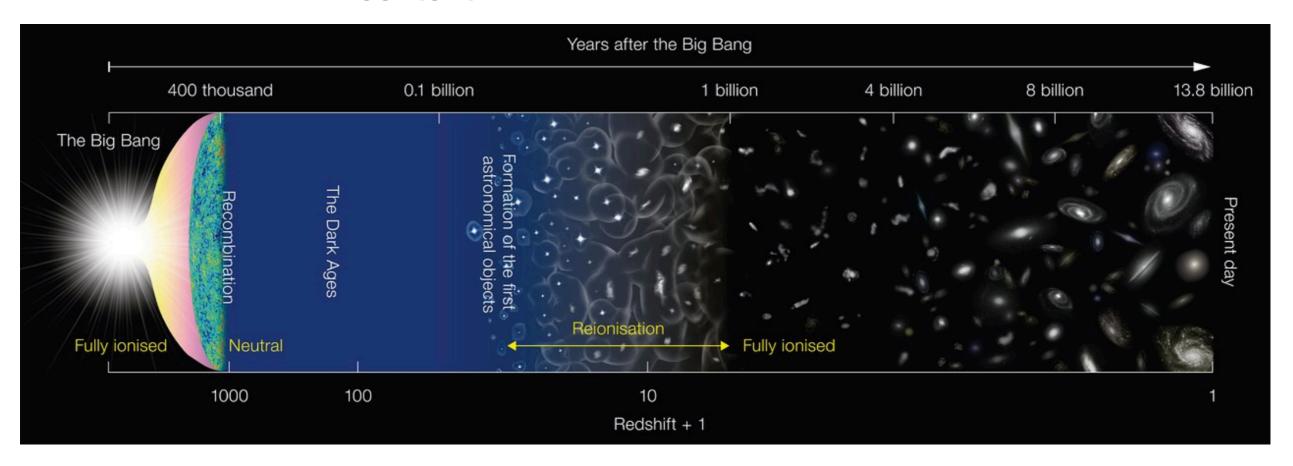
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## **Expansion/Energy** content



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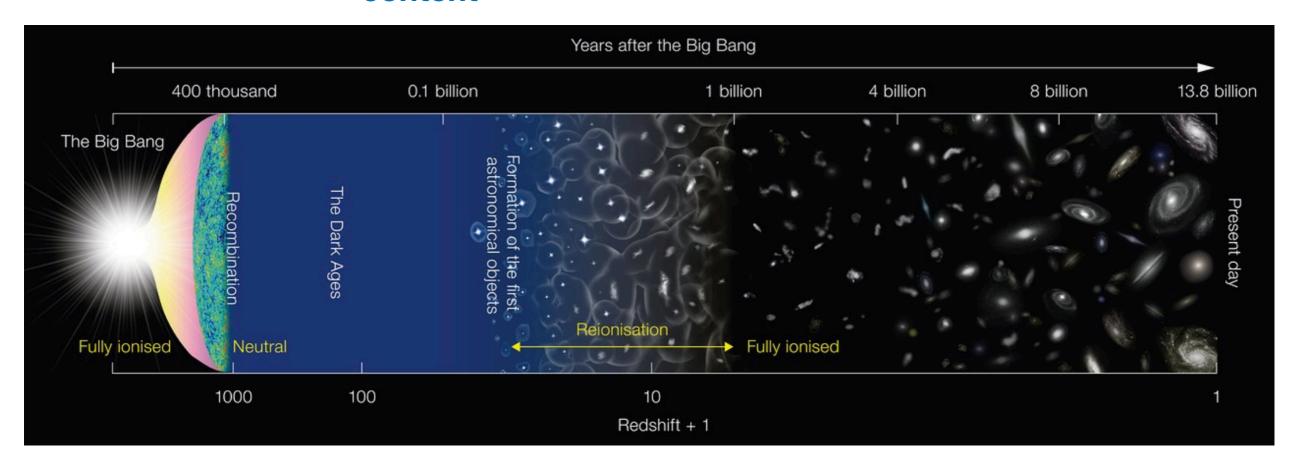
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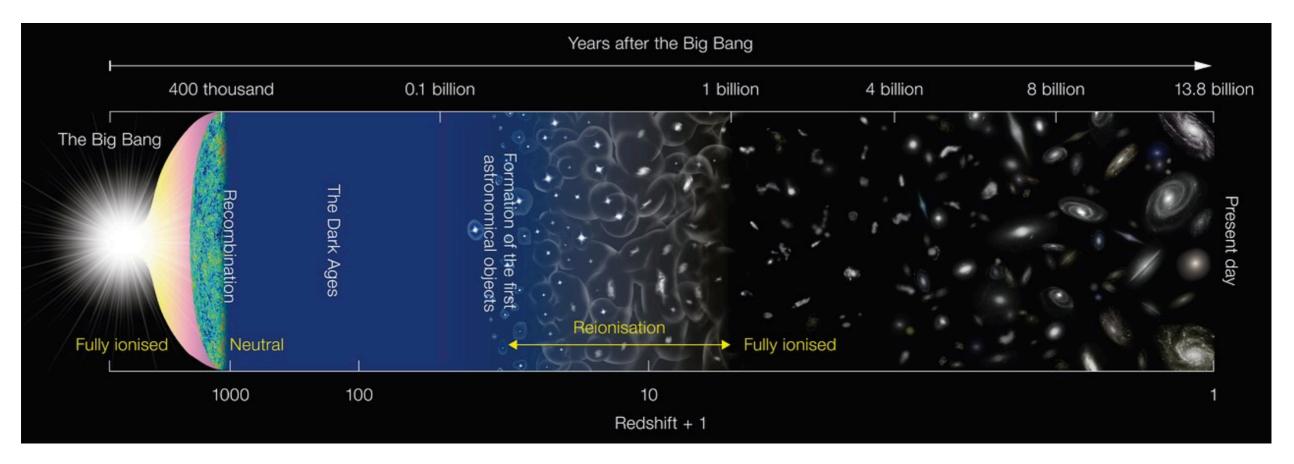


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**Expansion/Energy** content

**Inflation** 

star formation



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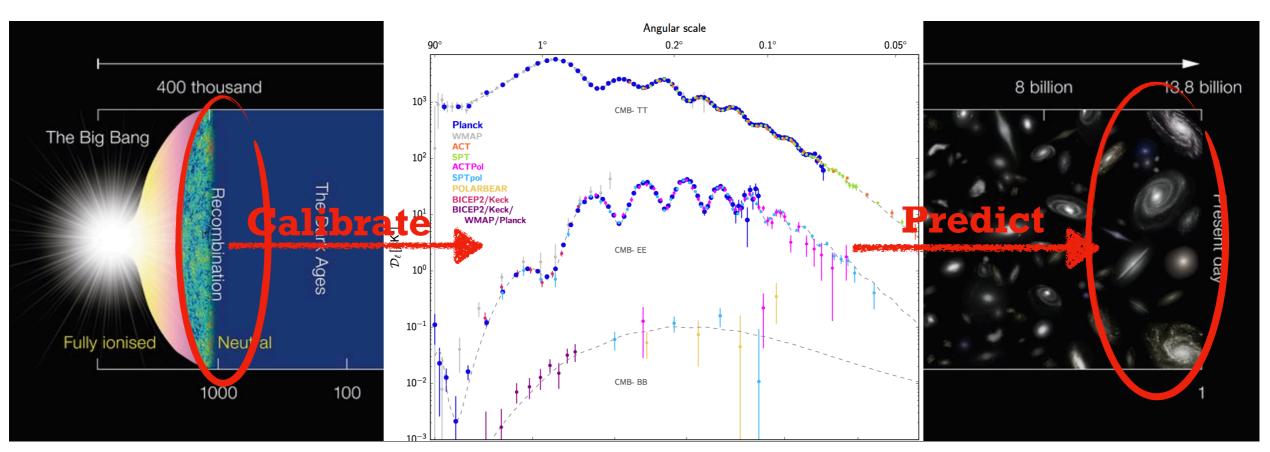


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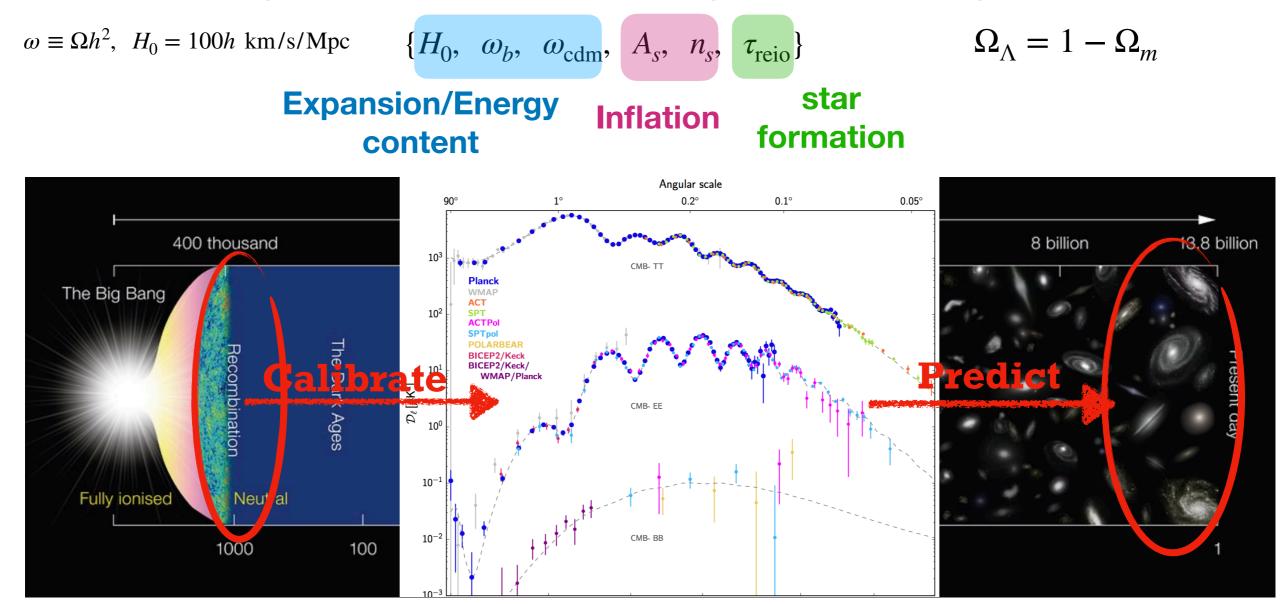
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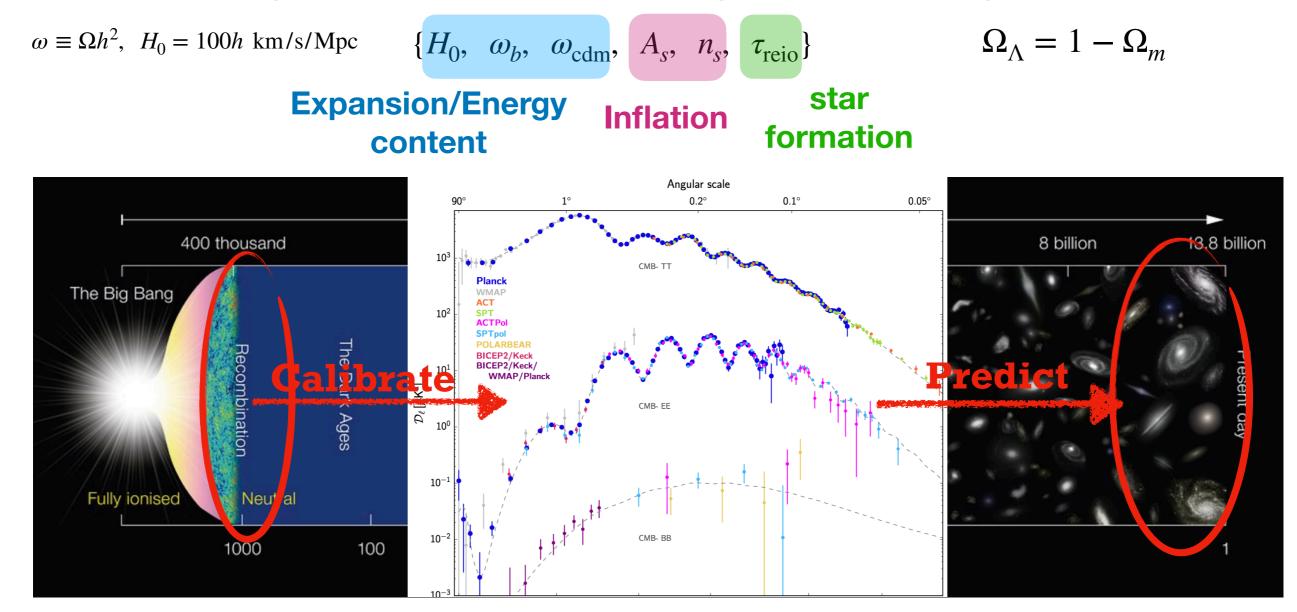


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95% of the energy budget today is unknown! 70% Dark Energy, 25% Dark Matter.

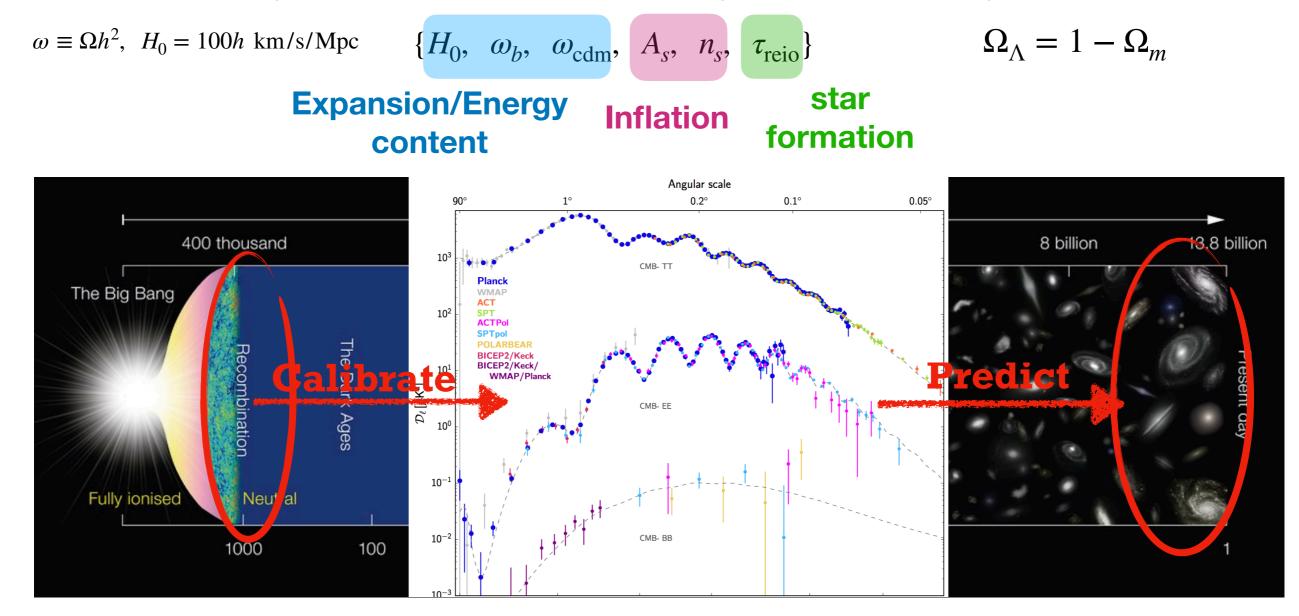
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The mechanism behind initial conditions is unknown.

How star formation happened and re-ionized the universe is unknown.

## Precision Cosmology or Cosmic Discordance?

#### The ∧CDM Cosmology is under extreme pressure

Cosmic dipole anomaly? The universe is not isotropic?

Colin++ 1703.09376, 1808.04597, Secrest++ 2009.14826, Alari++ 2207.05765, Guandalin++ 2212.04925

Cosmic void? The universe is not locally homogeneous?

Wu&Huterer 1706.09723, Kenworthy++ 1901.08681, Cai++ 2012.08292, Camarena++ 2205.05422

Tensions in cosmological parameters?

*Abdalla*++ 2203.06142

• Anomalies in CMB data? Evidence for a curved universe?

*Di Valentino++ 1911.02087, Calderón++ 2302.14300* 

Hints of dynamical dark energy and negative neutrino masses?

Union3 2311.12098, DES 2401.02929, DESI 2404.03002

(Too) High redshift galaxies with JWST?

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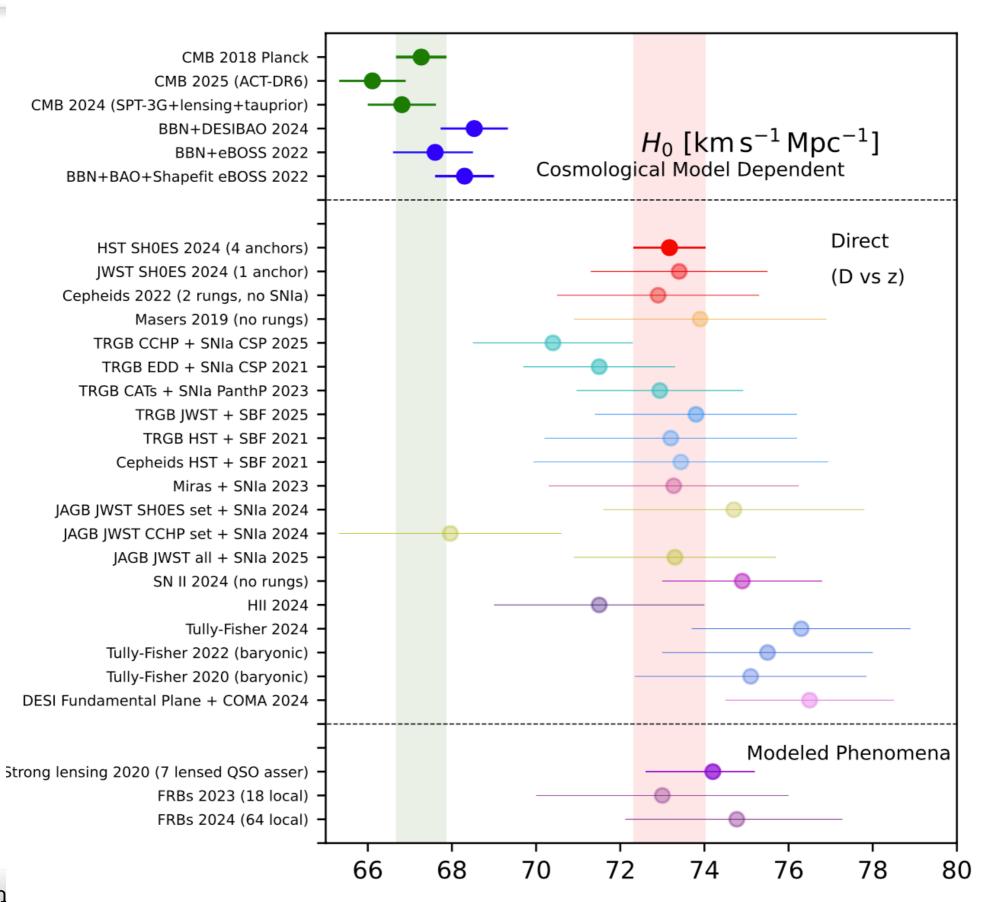
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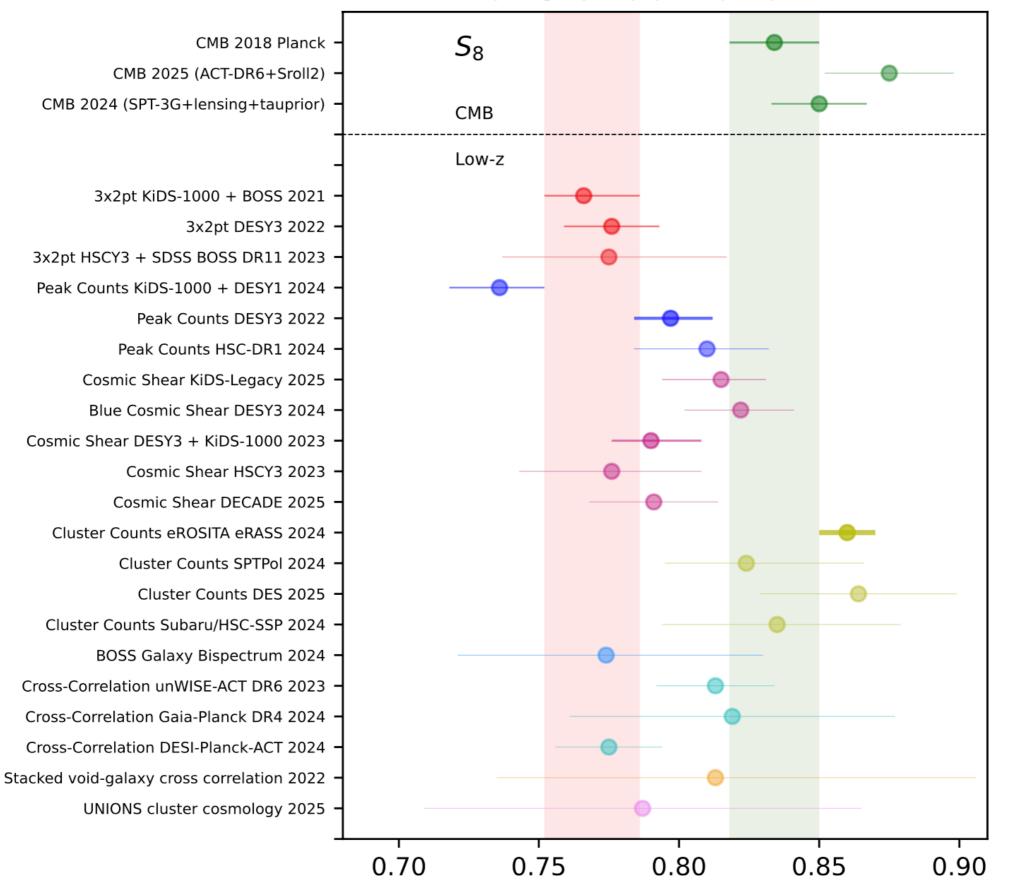
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# Is this a sign of a break down in the cosmological principle or GR? Are these the first signs of the true nature of DM and DE?

### The Hubble tension

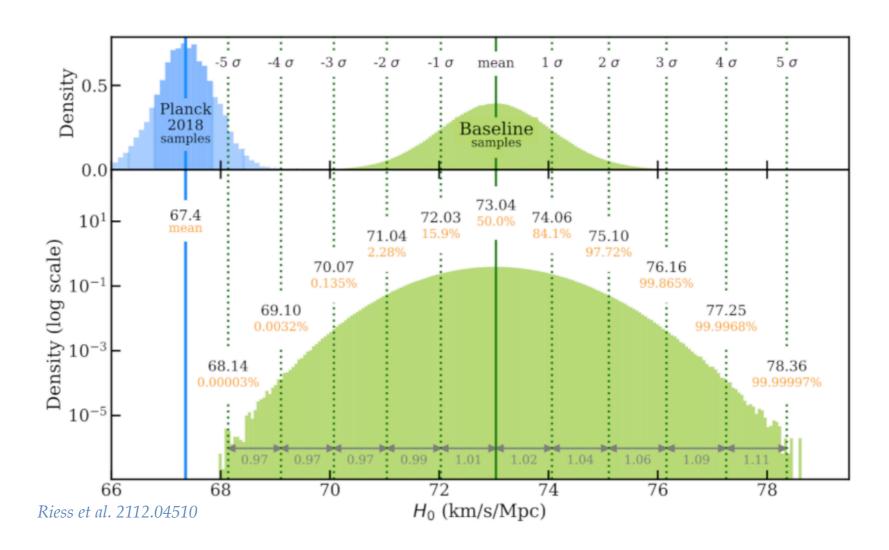


### The S8 tension



### The Hubble tension between SH0ES & Planck

There is a  $5\sigma$  discrepancy between the SH0ES and *Planck* determination of the Hubble parameter

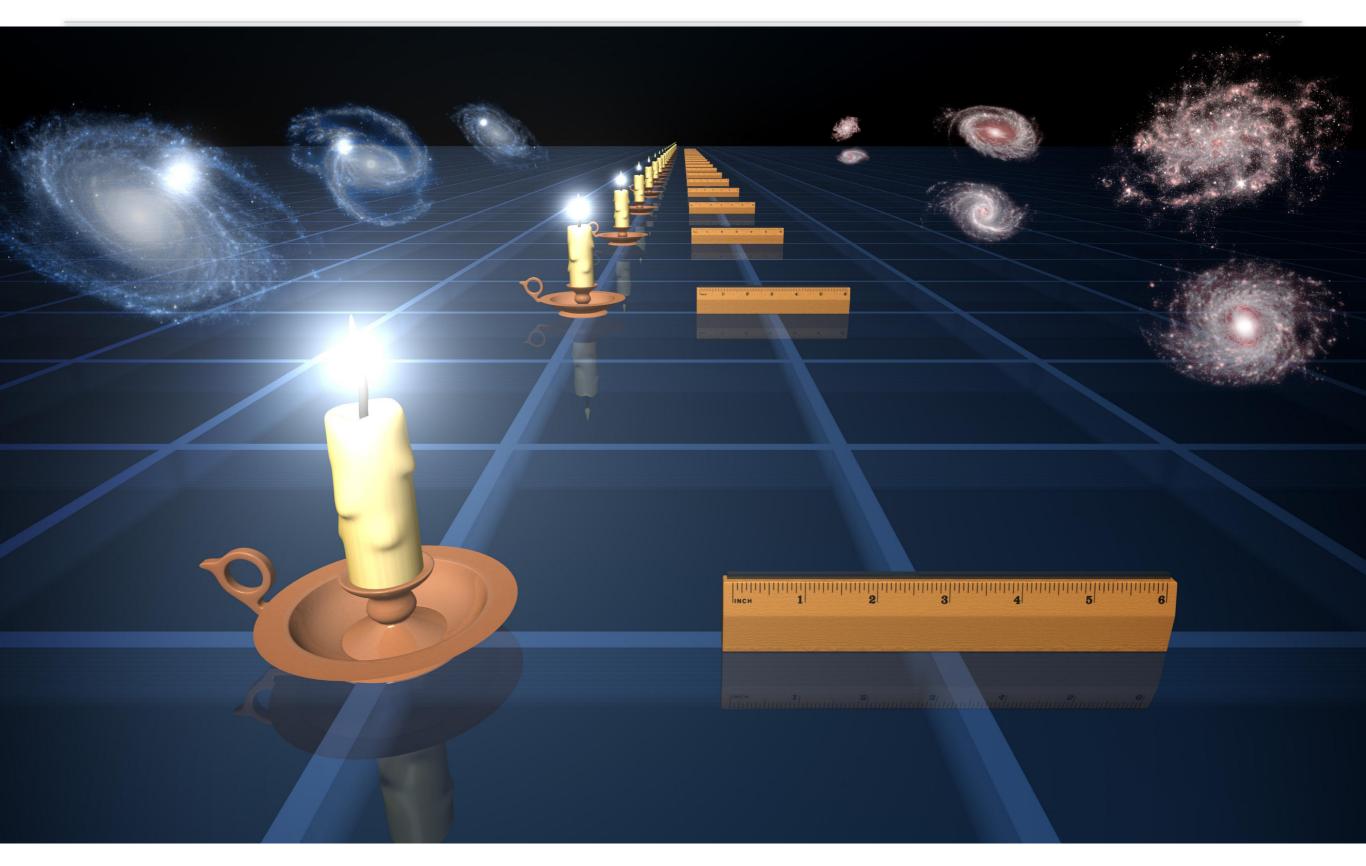


$$H_0(\Lambda \text{CDM/Planck}) = 67.4 \pm 0.5 \text{ km/s/Mpc}$$

$$H_0(SH0ES) = 73.04 \pm 1.04 \text{ km/s/Mpc}$$

Is this first sign of a crack in  $\Lambda$ CDM?

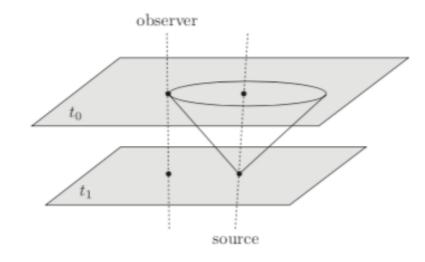
### Measuring distances in Cosmology



The comoving distance  $\chi = \int_{t_e}^{t0} \frac{dt}{a(t)} = \int_{0}^{z} \frac{dz}{H(z)}$  from imposing  $ds^2 = dt^2 - a^2 d\chi^2 = 0$ 

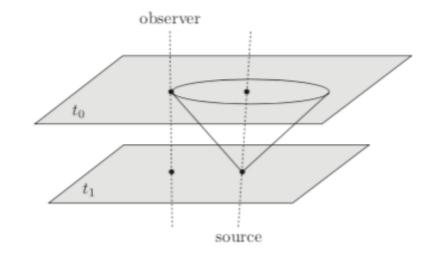
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- The luminosity distance to a source of known flux and luminosity

$$F = \frac{L}{4\pi d_L(a)^2}$$
  $d_L = \frac{\chi(a)}{a} = \chi(z)(1+z)$ 



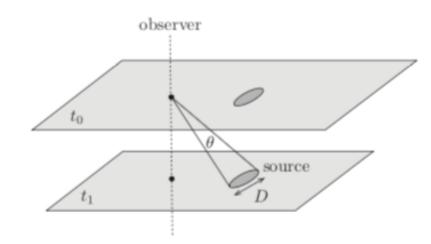
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• The angular diameter distance to a source of known apparent and intrinsic size.

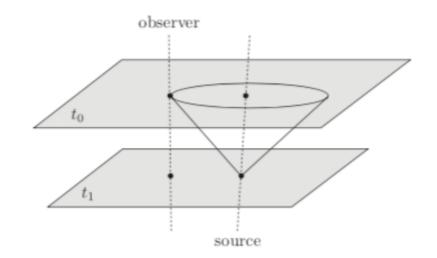
$$\delta\theta = \frac{D}{d_A}$$
  $d_A = a\chi(a) = \frac{\chi(z)}{(1+z)}$ 



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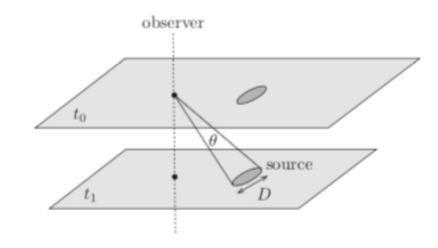
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• In GR the Etherington distance-duality relationship imposes

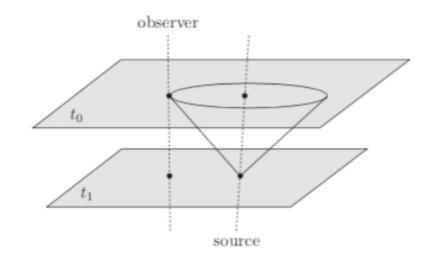
$$d_A = d_L/(1+z)^2$$

$$d_{A} = \frac{1}{1+z} \int_{0}^{z} \frac{dz}{H_{0}\sqrt{\Omega_{m}(1+z)^{3} + \Omega_{r}(1+z)^{4} + \Omega_{\Lambda}}}$$

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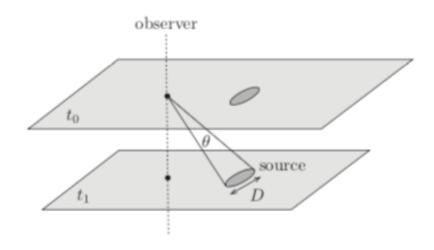
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Measuring  $H_0 ==$  measuring distances Mismatch in  $H_0 ==$  wrong distance!

Standard candle: Object of known luminosity L

$$F = \frac{L}{4\pi D_L^2}$$

$$m \equiv -2.5\log F + \text{const.}$$
  
 $M \equiv -2.5\log F(10 \text{ pc}) + \text{const.}$ 

$$m - M = 5\log D_L[\text{Mpc}] + 25$$
  
\*1 Mpc =  $3.2 \times 10^6 \, 1. \, \text{y}$ .

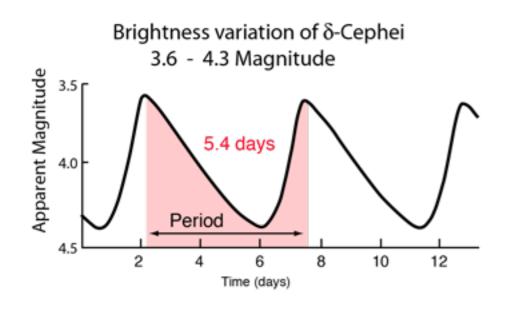
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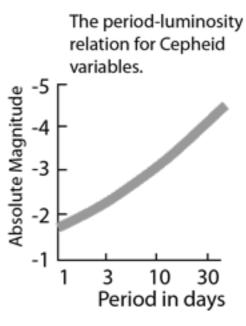
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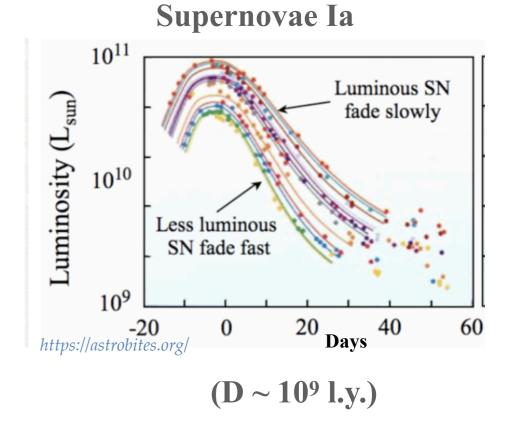
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#### **Cepheid stars**



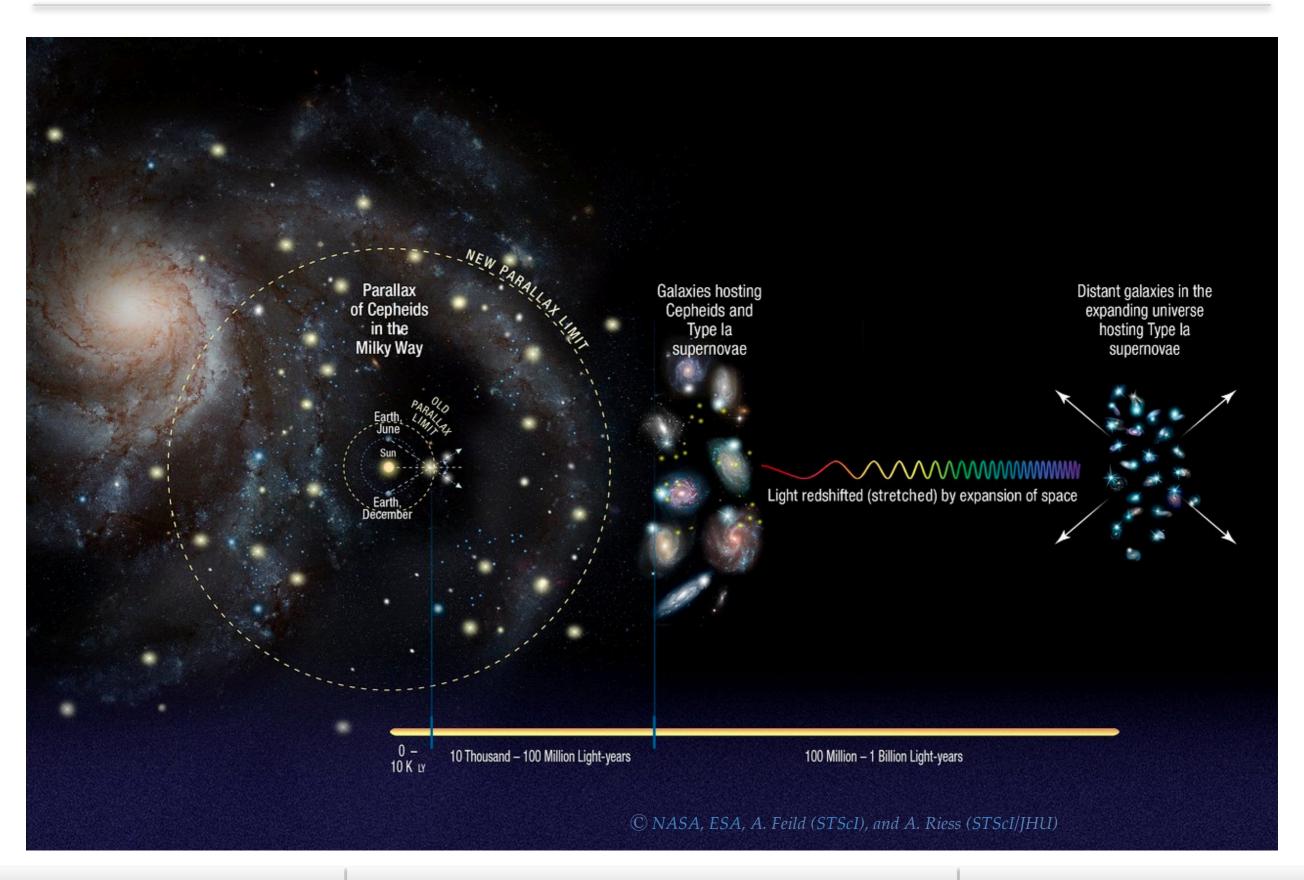




(D ~  $10^3$  -  $10^6$  l.y.)

Pulsating starsPeriod-luminosity relation

Bright "standardizable" objects
 relation between peak & slope of the light curve

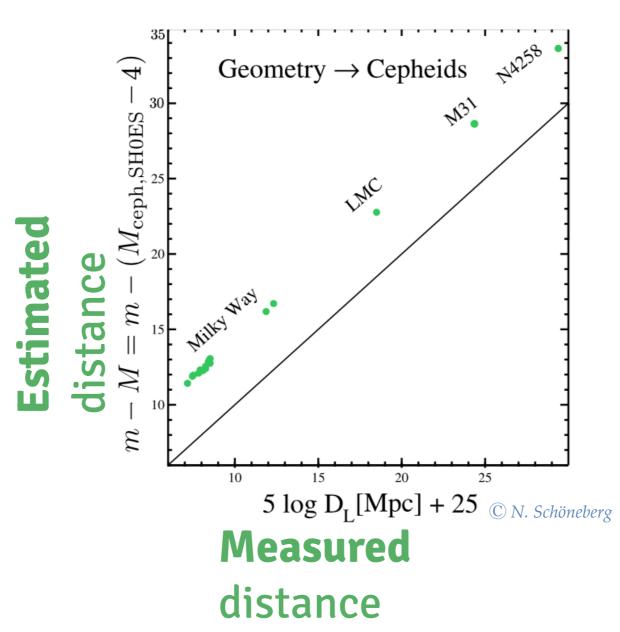


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#### Different geometric calibrators: GAIA parallaxes, masers NGC 4258, DEB in the LMC

$$m - M = 5\log D_L[\text{Mpc}] + 25$$





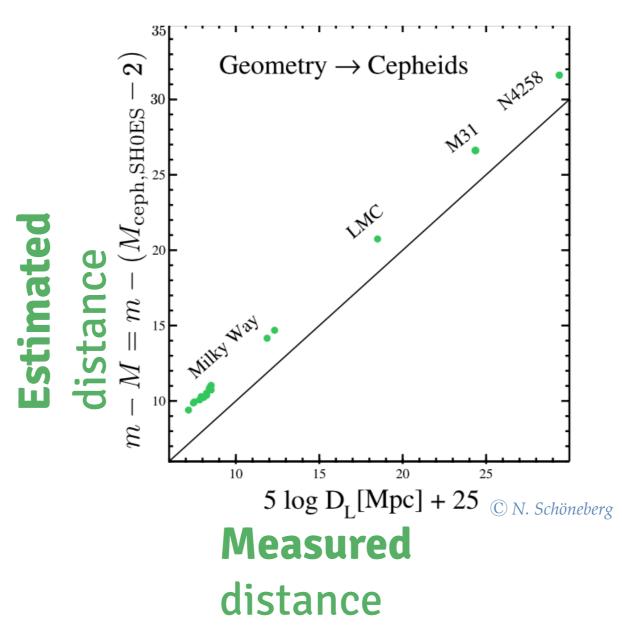
[Riess+1604.01424 (edited)]

11

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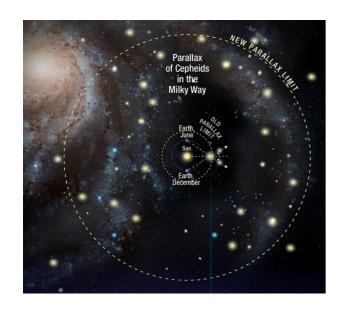


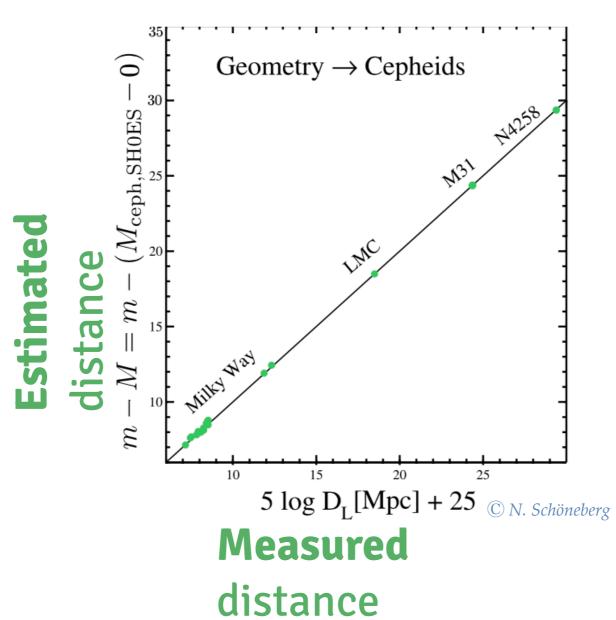
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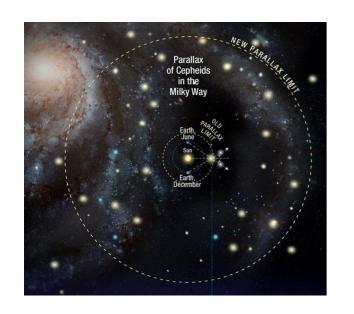


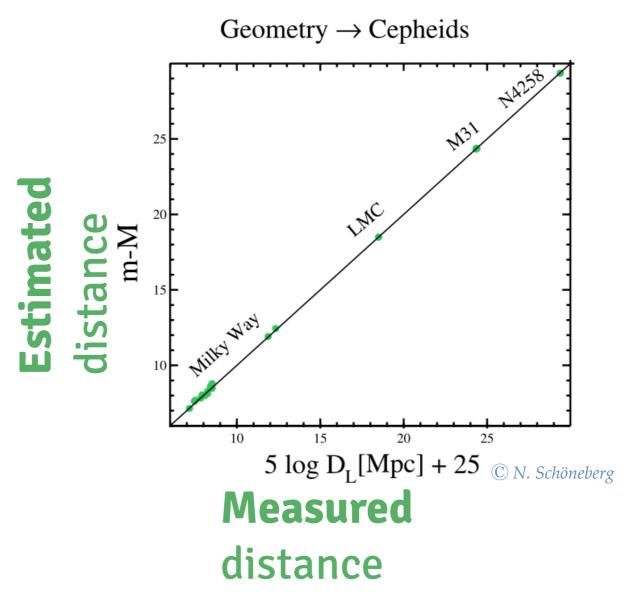
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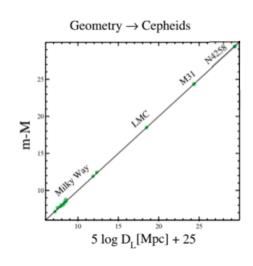
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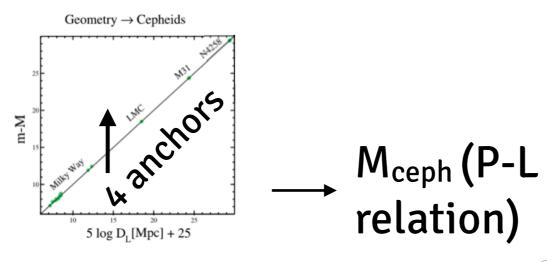
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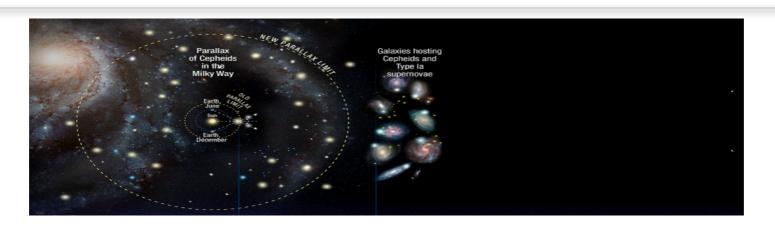


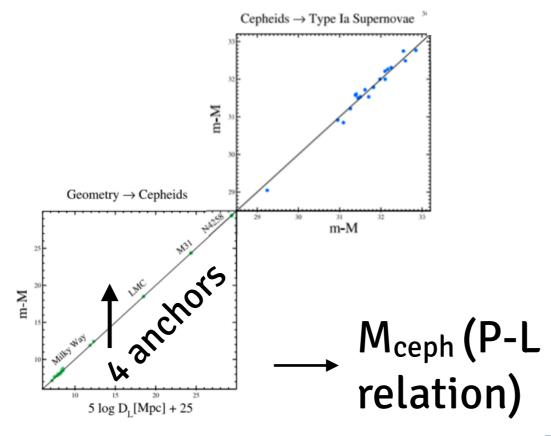
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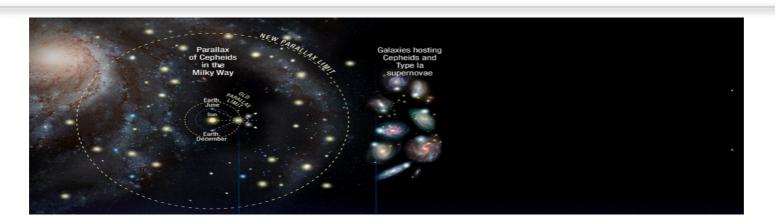


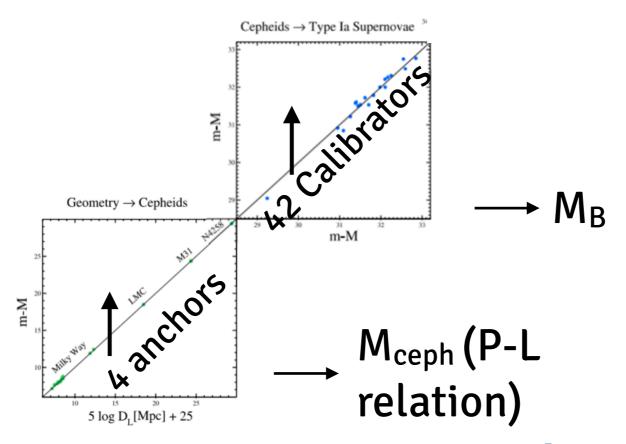
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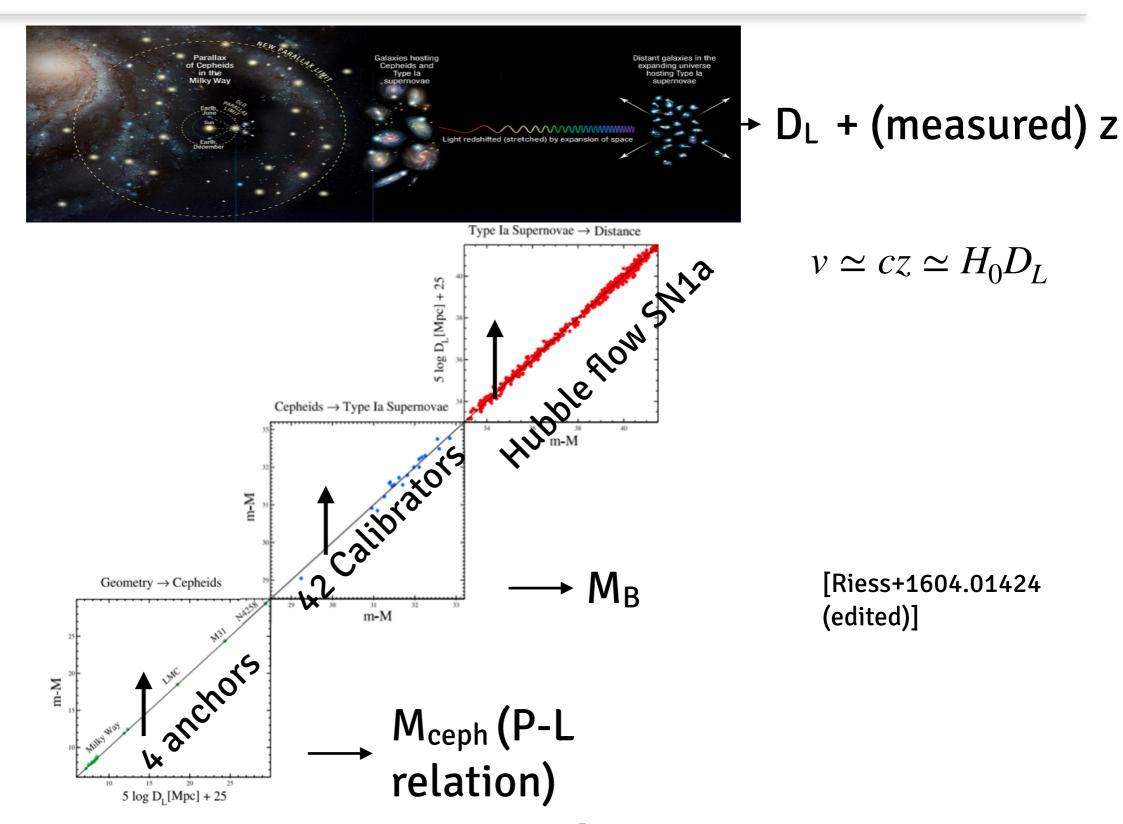


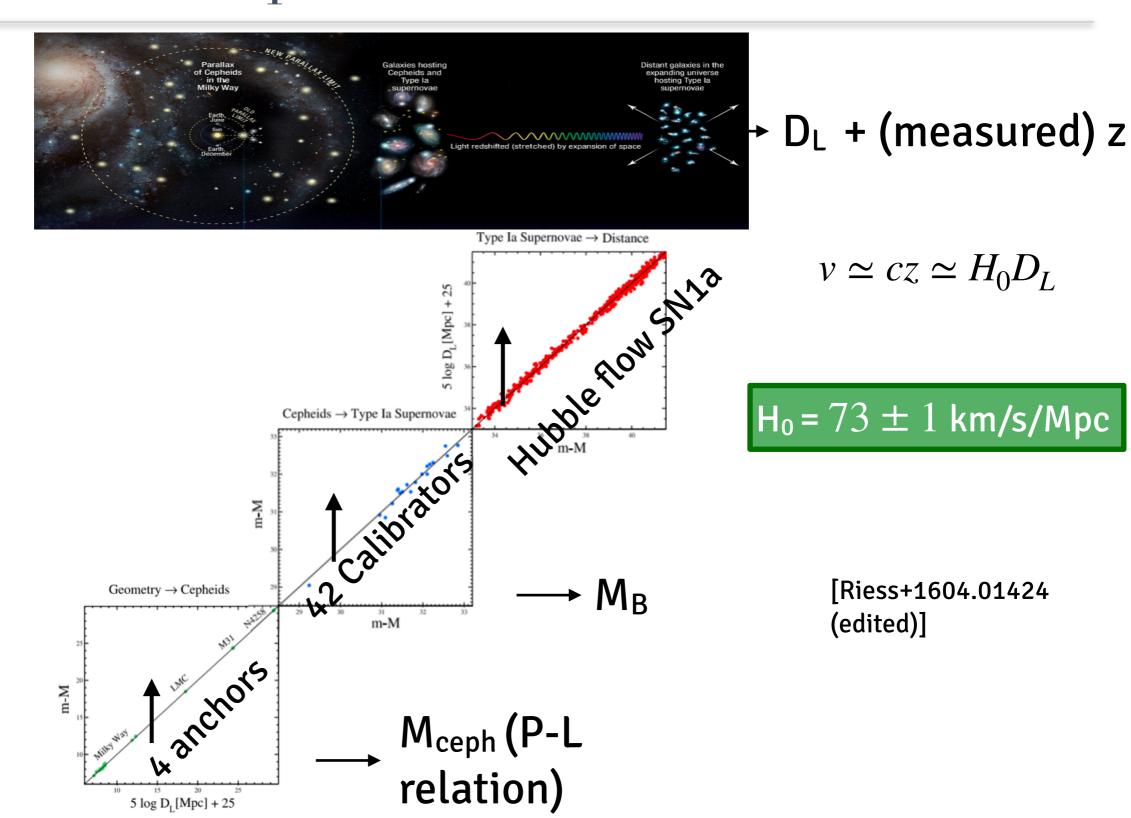
[Riess+1604.01424 (edited)]





[Riess+1604.01424 (edited)]





## Systematics? A non-exhaustive list

See review Di Valentino++ 2103.01183 for all relevant references

- SH0ES builds a 3 steps distance ladder: anchors => cepheids => SN1a
- Are there issues with distance anchor? (GAIA, LMC, NGC4258)

*Efstathiou++ 2007.10716, Soltis++2012.09196* 

- Are there issues with cepheids?
  - Cepheids vs TRGB: disagreement?

Freedman++ 2106.15656, Anand++ 2108.00007

Effect of Dust?

Mortsell++ 2105.11461

Cepheid crowding?

Riess++ 2401.04773

Is the metallicity correction correct?

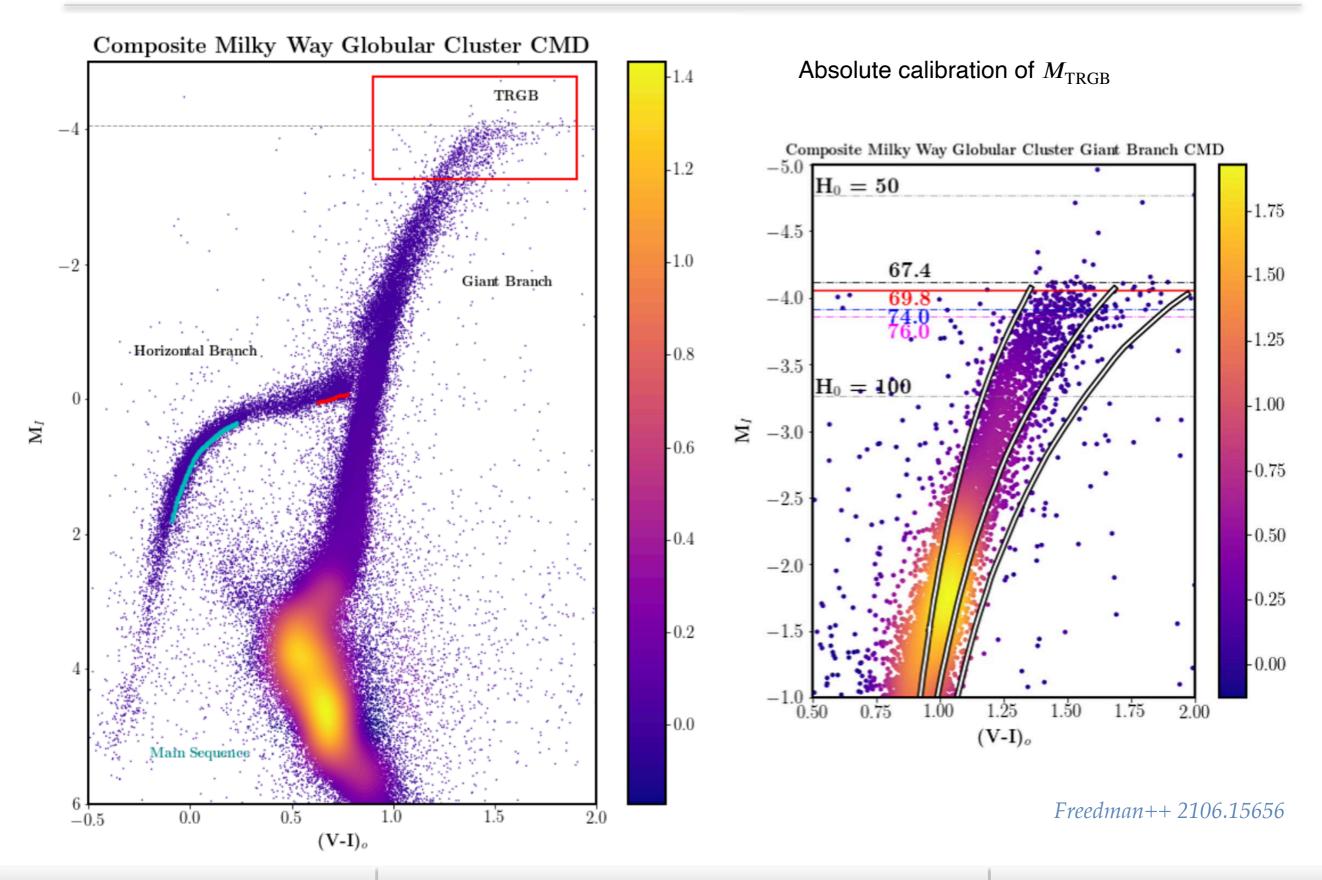
*Efstathiou++* 2007.10716

• Are there issues with SN1a? different populations of SN1a between "cepheid-SN1a calibrator" and Hubble flow SN1a?

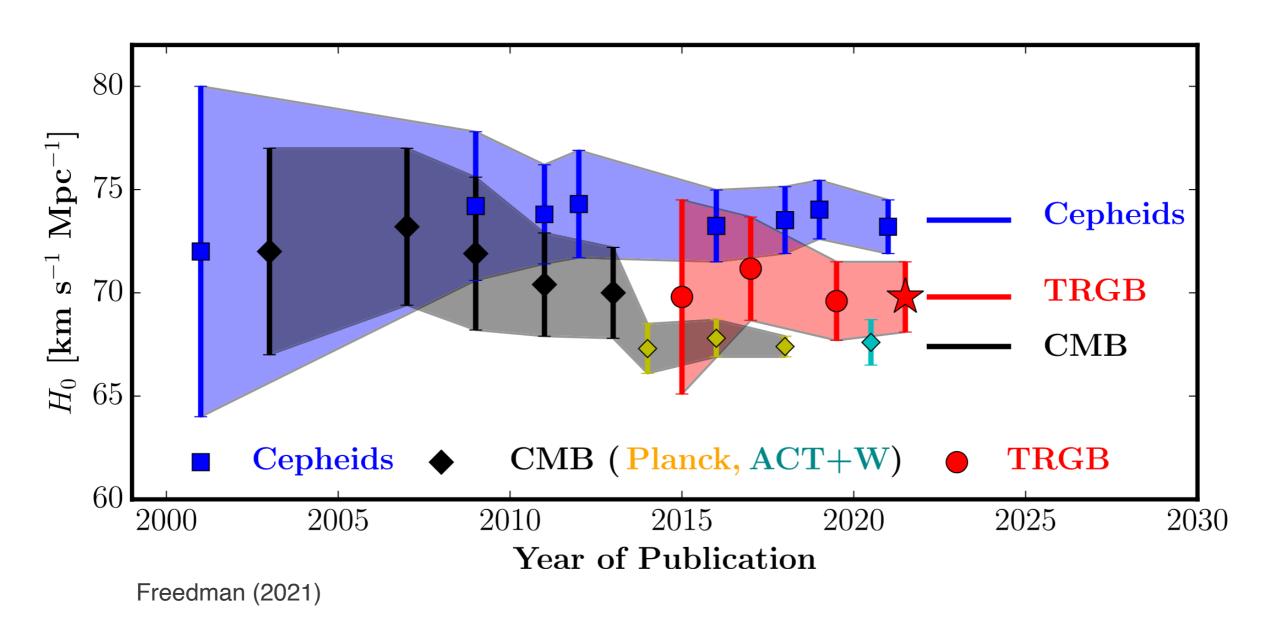
Rigault++ 1412.6501, Jones++1805.05911, Brout&Scolnic 2004.10206

• Are there issues with the CMB? Di Valentino++ 1911.02087, Calderón++ 2302.14300

## Tip of the Red Giant branch



## The Chicago Carnegie Hubble Program

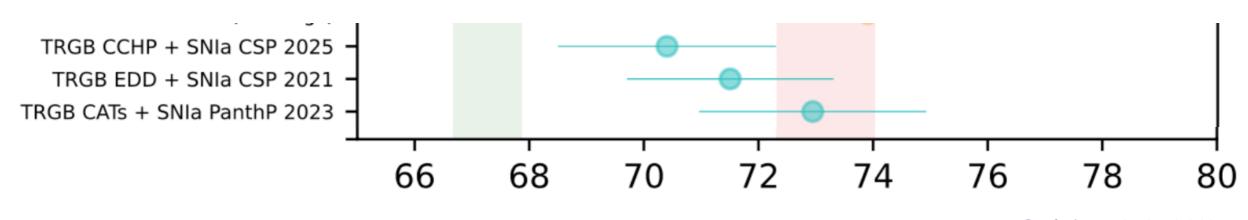


*Freedman++ 2106.15656* 

## Tip of the Red Giant branch revisited

**Table 5.** Sources of Differences in  $H_0$  Between TRGB analysis by CATs (Here), CCHP, EDD (in  $H_0$ )

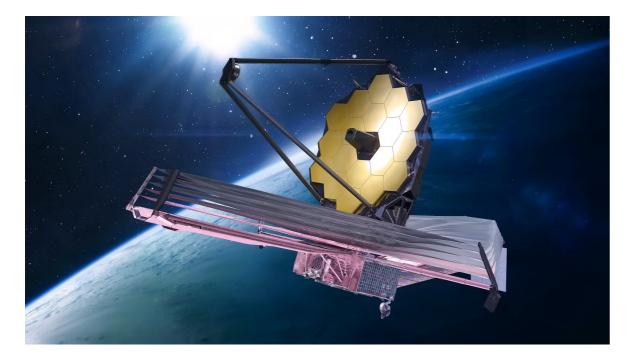
Term	$\Delta \text{CCHP}$	$\Delta  ext{EDD}$
	(km/s/Mpc)	(km/s/Mpc)
SN Related		
1. Include SN 2021pit,2021rhu,2007on	0.6	1.3
2. No TRGB detected in N5584,N3021,N1309,N3370	0.0	0.0
3. Peculiar Flows (Pantheon+)	0.4	0.0
4. Hubble Flow Surveys (Pantheon+)	1.1	0.0
SN subtotal	2.0	1.3
TRGB Related		
5. Fiducial TRGB Calibration/Tip-Contrast Relation	1.4	-0.3
Total	3.4	1.0

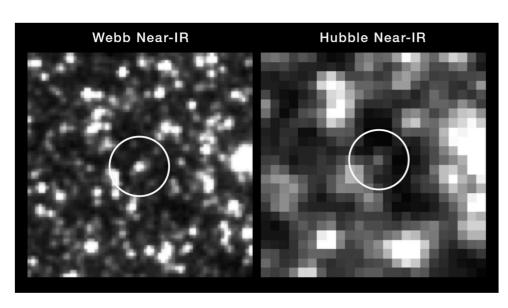


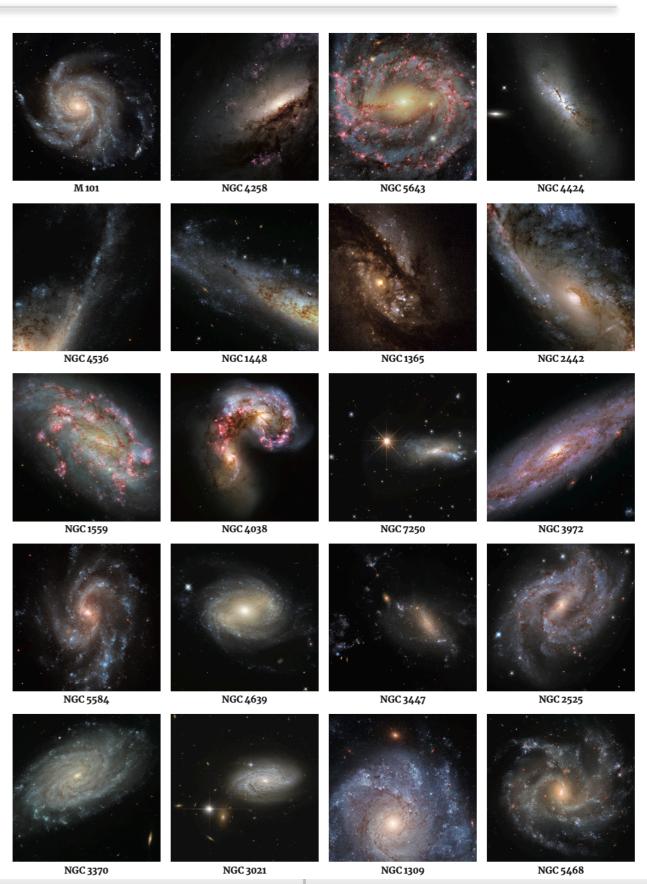
Scolnic++ 2304.06693

### JWST and the Hubble tension

- First analyses with JWST from CCHP and SH0ES.
- Re-observations of 20 key galaxies to check HST results + develop new calibration method.

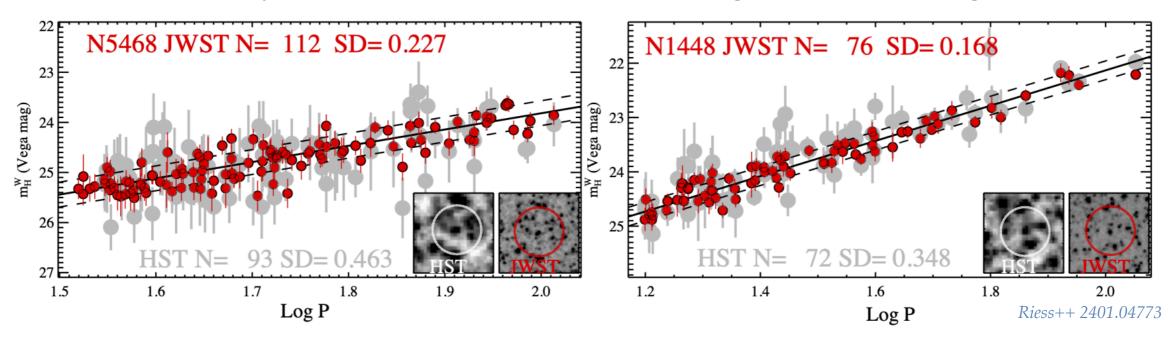




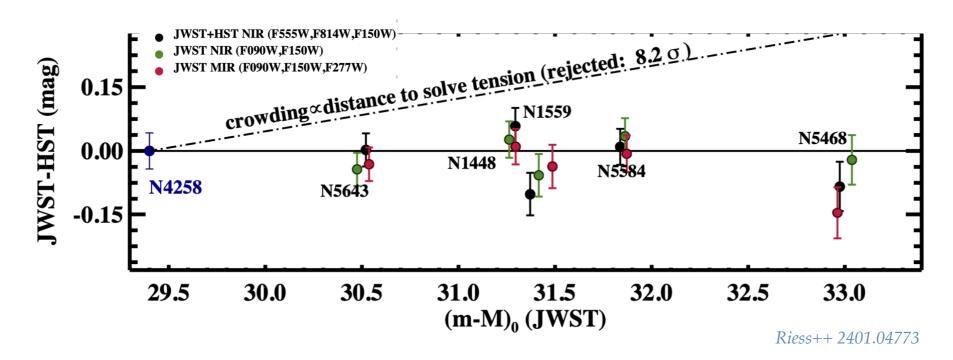


# SHOES finds excellent agreement with HST

• From SH0ES: reanalysis of the PL relations in 8 SN1a hosts galaxies, excellent agreement with HST

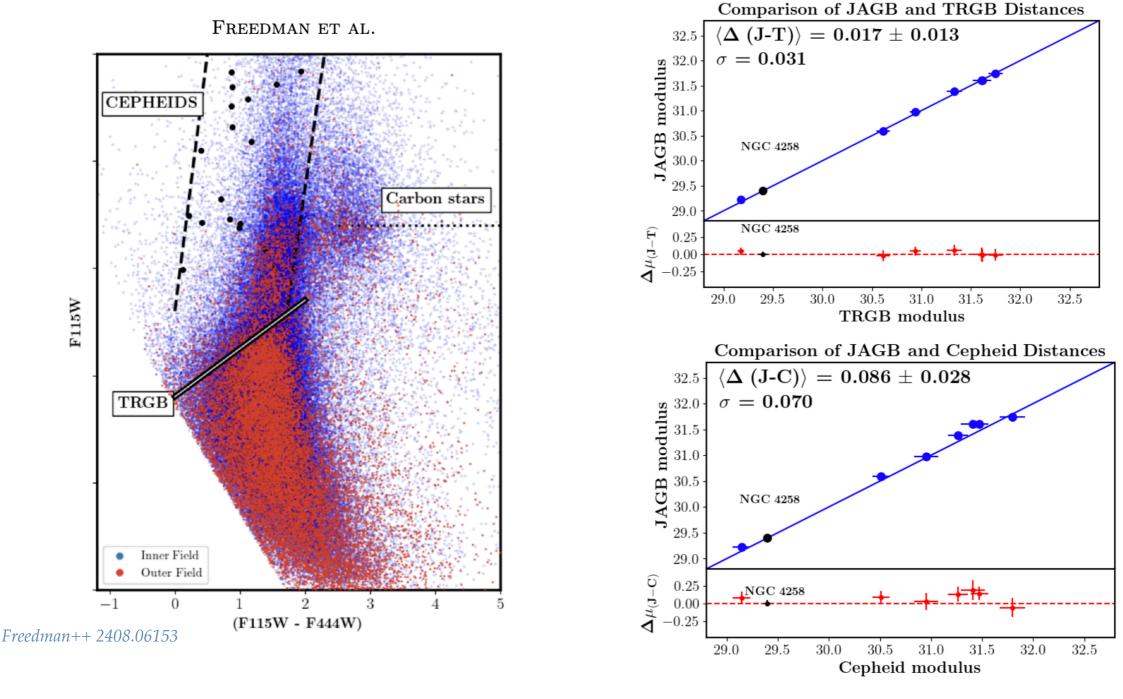


• Rejects "crowding" of cepheids as an explanation for the tension



# CCHP: 3 JWST-only measurements of Ho

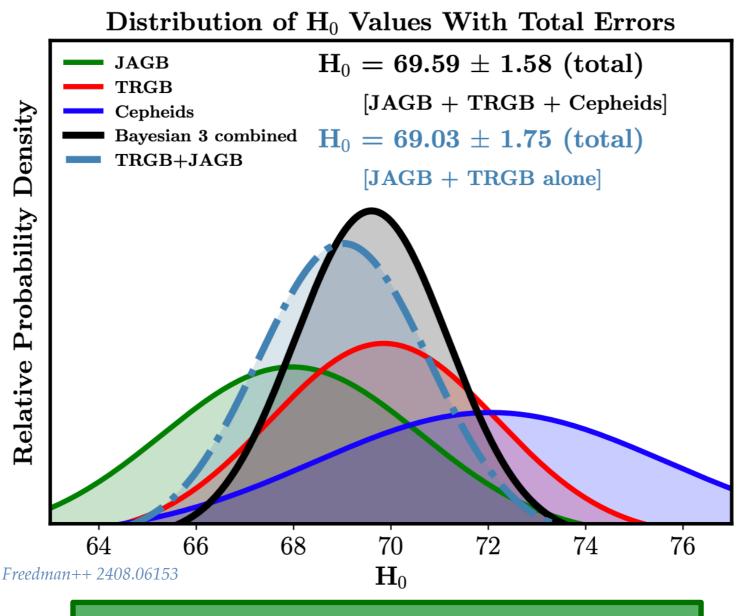
Observations of 10 SN1a hosts + 1 anchor to re-calibrate cepheids, TRGB and a new 'JAGB' method.



• Finds a bias in the cepheids distance while TRGB and JAGB distance are in good agreement

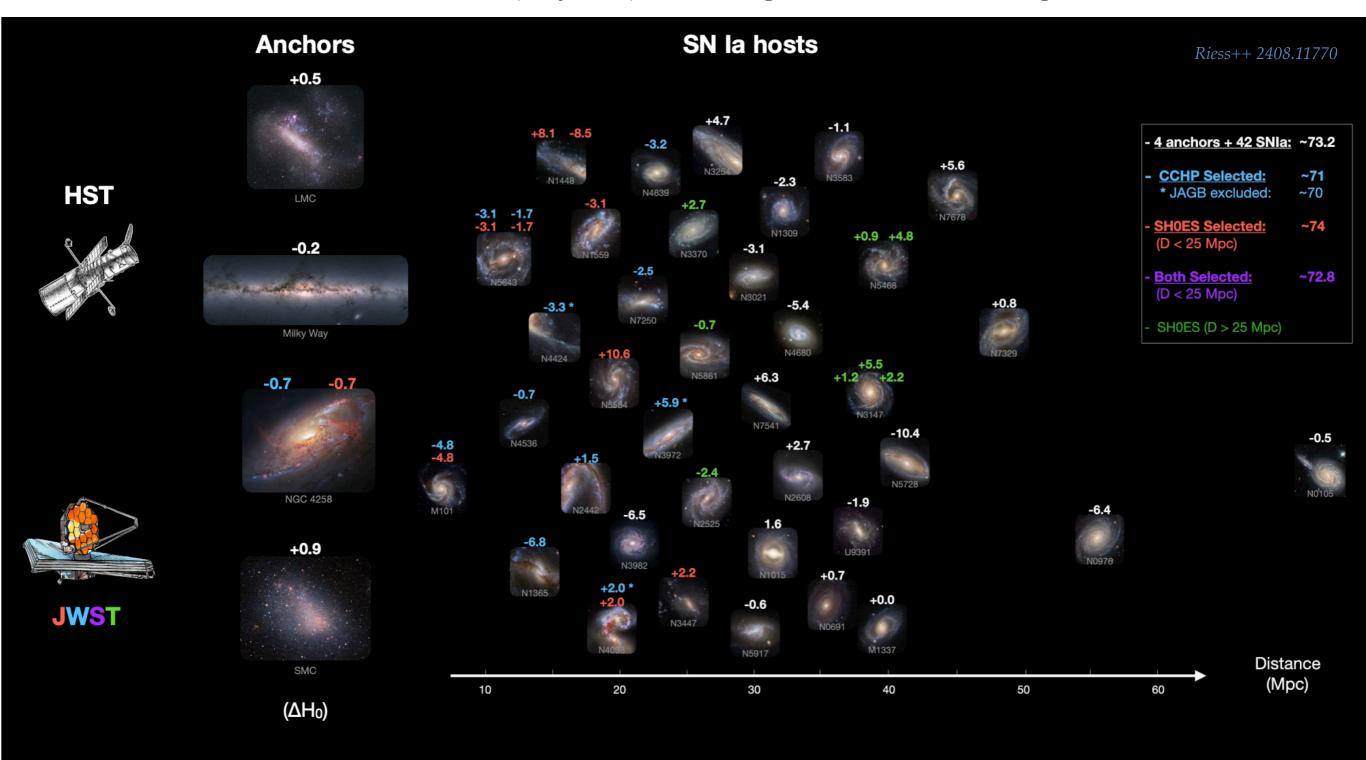
#### CCHP finds no Hubble tension

- JAGB and TRGB value of  $H_0$  in good agreement with  $\Lambda$ CDM, Cepheids are 'biased high'.
- Error bars are large: JWST alone is not (yet) as good as HST, only 10 hosts galaxies and one anchor.



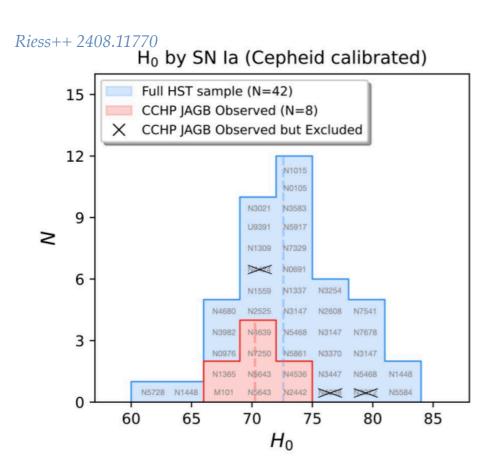
# HST provides a 'complete' picture

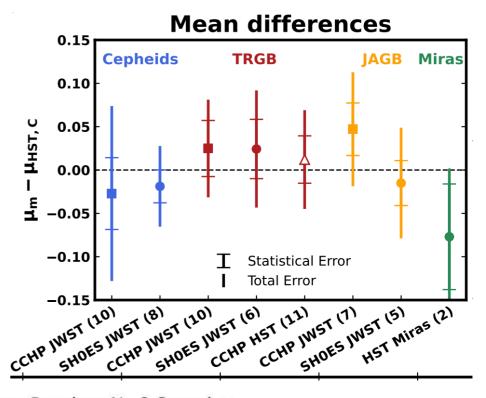
JWST measures (very well) a sub-sample of the full HST sample

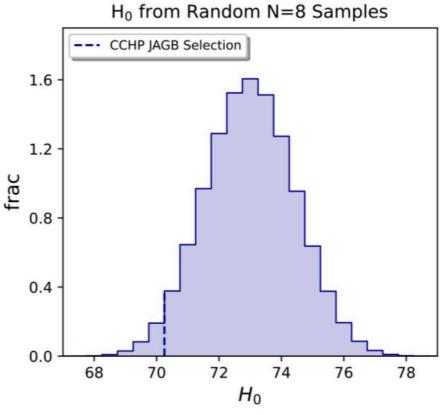


# SHOES suggests a 'bias low' in CCHP samples

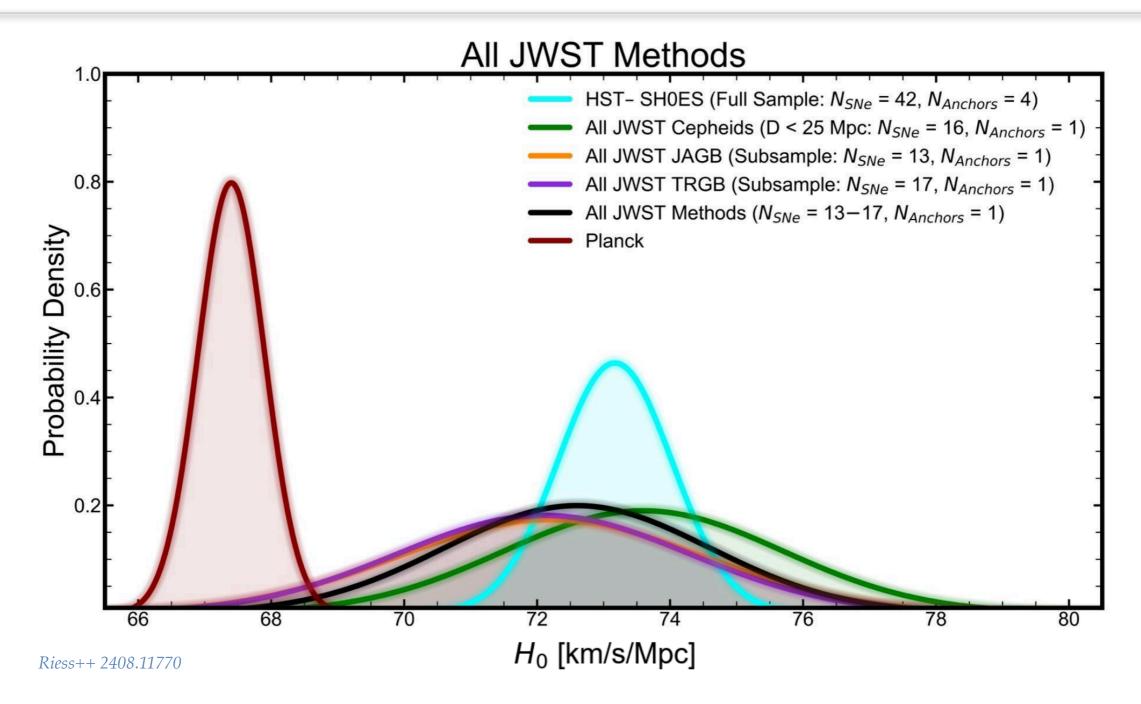
- JWST in very good agreement ( $< 1\sigma$ ) between cepheid distances and all other methods from HST
- Identified a missing source of error in the CCHP cepheid
- JAGB sample of host galaxies is 'biased low' and this is expected!







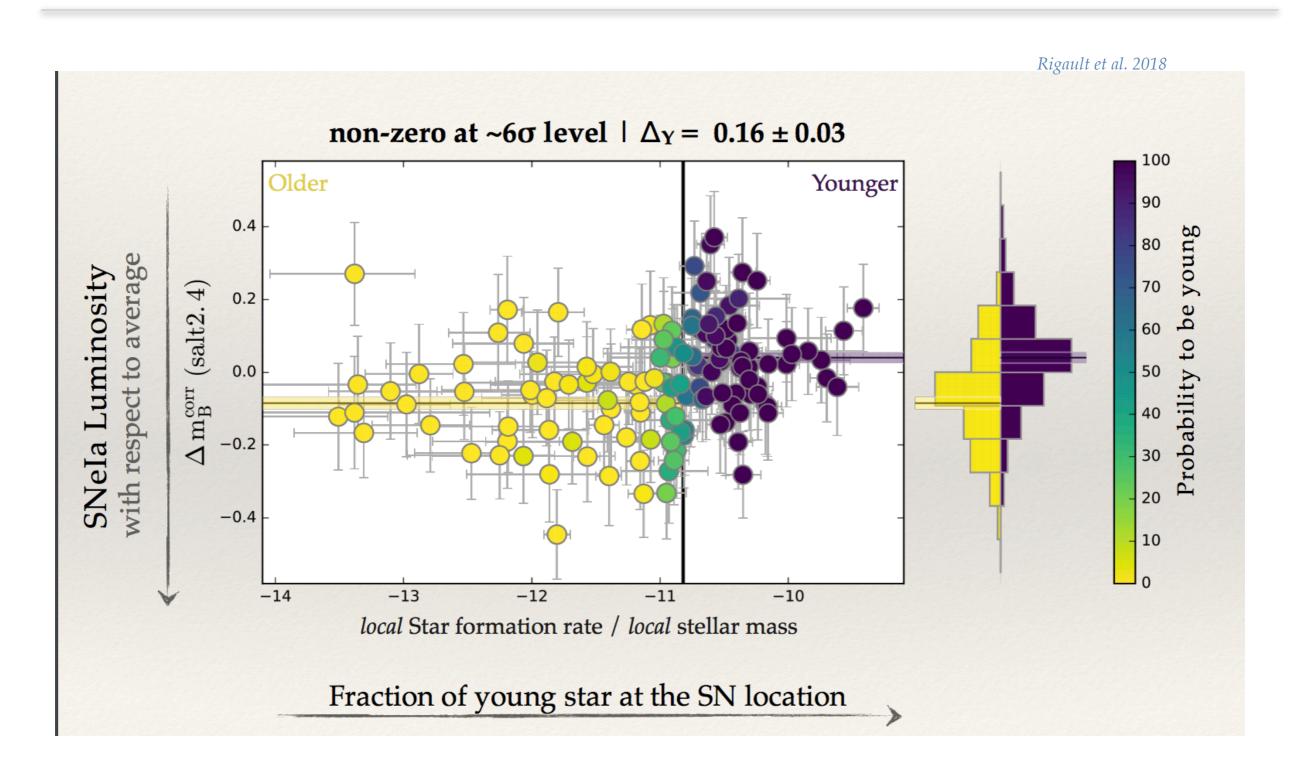
### SHOES confirms the Hubble tension with JWST



• The situation will be settled by (re-)measuring the remaining SN1a host galaxies and anchors.

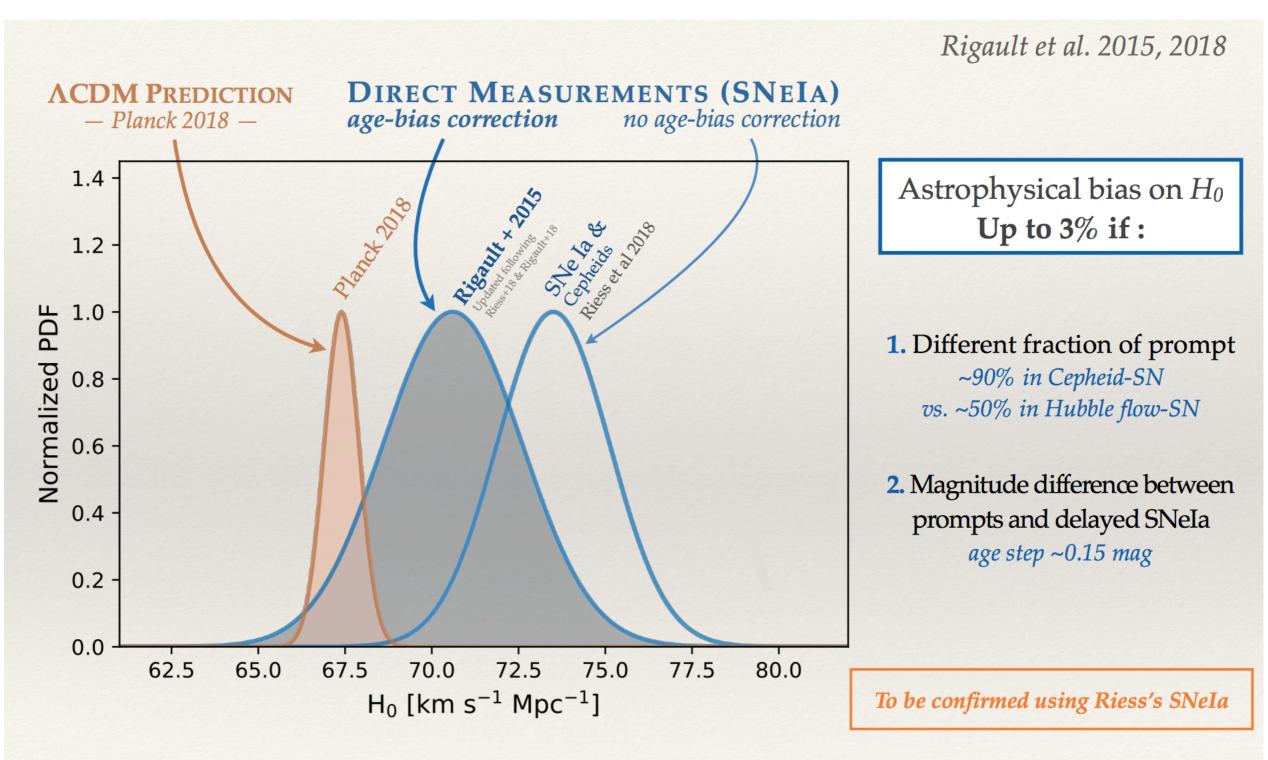
The Hubble tension is alive and well!

# The progenitor bias



Slide by M. Rigault, IAP 2018

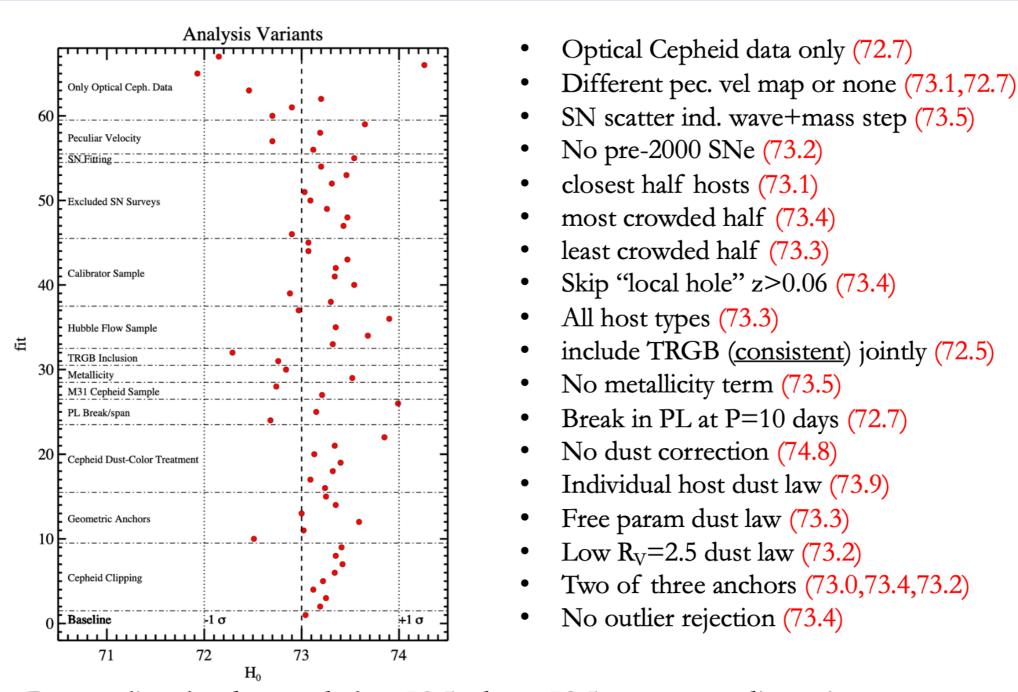
## Impact for $H_0$



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## SHOES estimates of systematic uncertainties

Analysis Variants: 12 categories, 67 variants, bifurcations, extensions, etc



Bottom line: hard to get below 72.5, above 73.5, propagate dispersion as extra systematic

#### A local void?

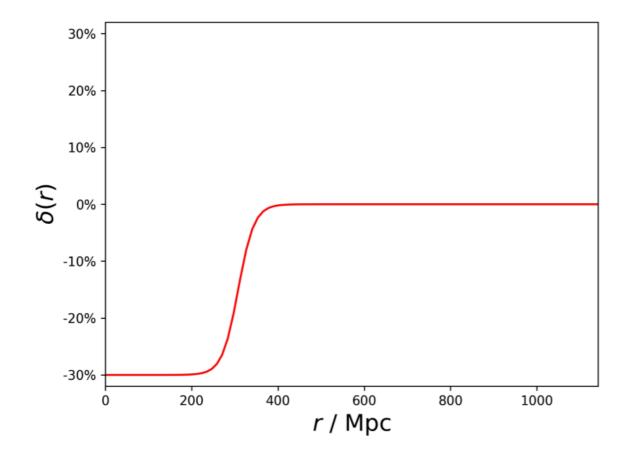
Kenworthy++ 1901.08681

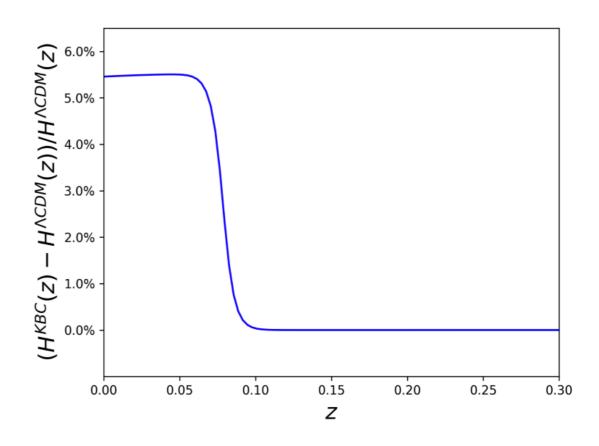
$$\frac{\Delta H}{H} = -\frac{1}{3}\delta \cdot f$$

$$ds^2 = -dt^2 + rac{R'^2(r,t)}{1 + 2r^2k(r)\tilde{M}^2}dr^2 + R^2(r,t)d\Omega$$

$$\delta(r) = \delta_V \frac{1 - \tanh((r - r_V)/2\Delta_r)}{1 + \tanh(r_V/2\Delta_r)}$$

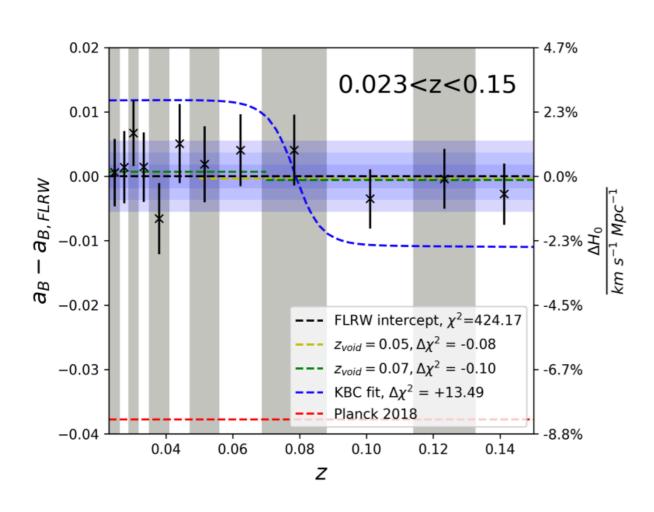
$$\delta(r) = \delta_V \frac{1 - \tanh((r - r_V)/2\Delta_r)}{1 + \tanh(r_V/2\Delta_r)} \qquad \frac{\dot{R}(r, t)^2}{R(r, t)^2} = H_0(r)^2 \cdot (\Omega_M(r) \frac{R_0(r)^3}{R(r, t)^3} + \Omega_k(r) \frac{R_0(r)^2}{R(r, t)^2} + \Omega_\Lambda(r))$$

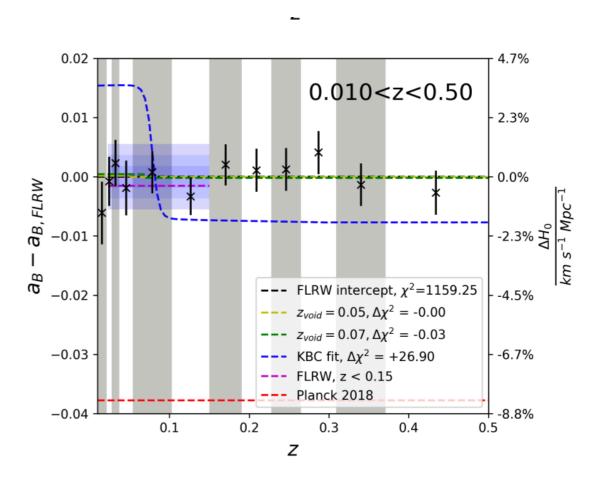




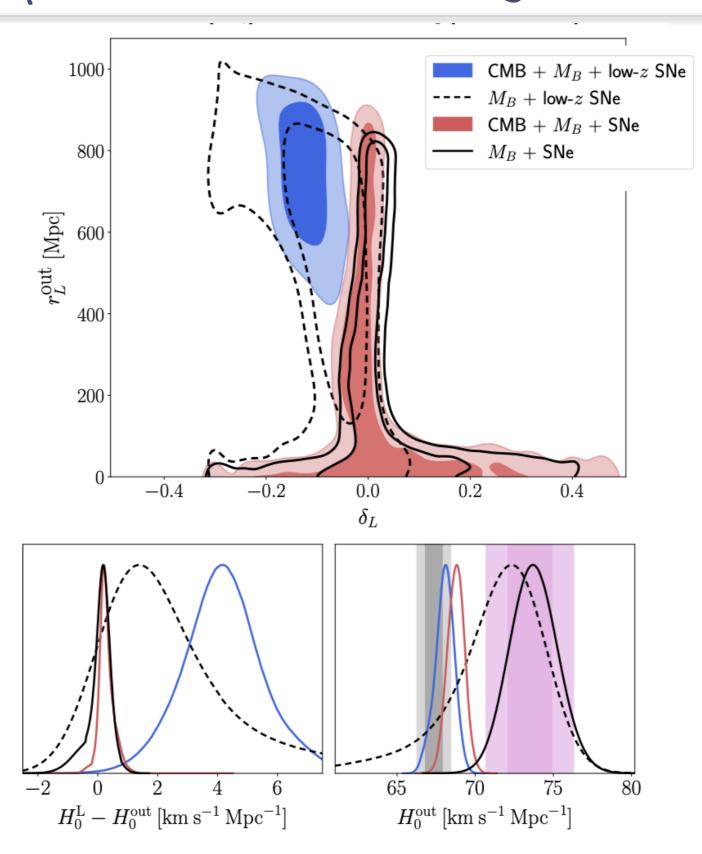
## Incompatible with the SNIa data

Kenworthy++ 1901.08681



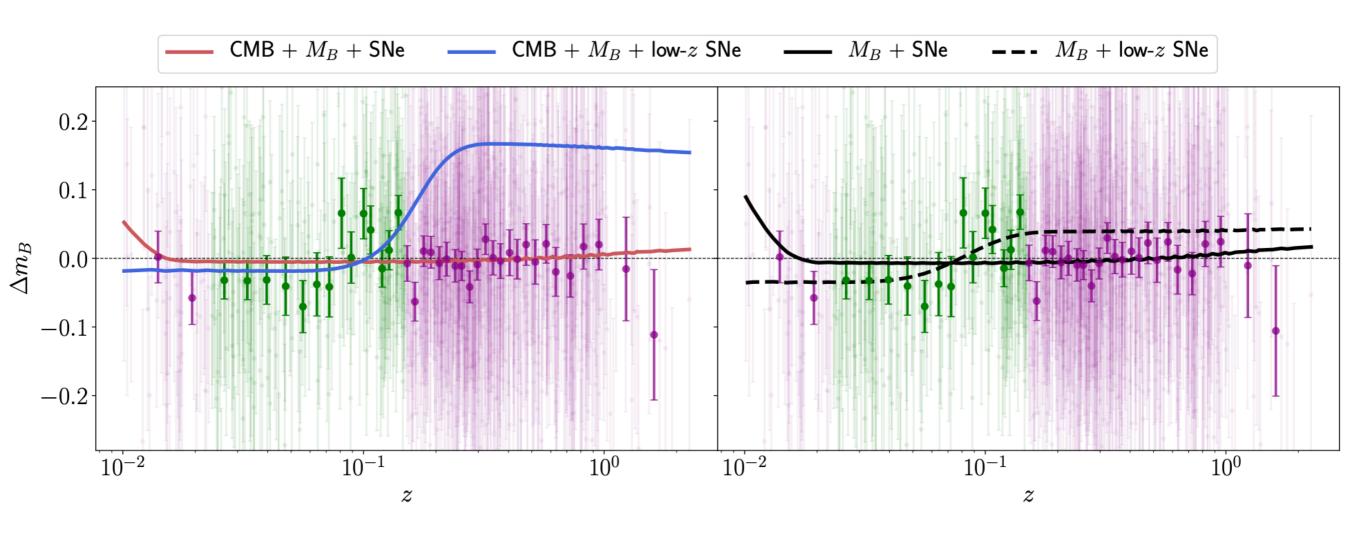


## Impact of local void beyond KBC



*Camarena++ 2205.05422* 

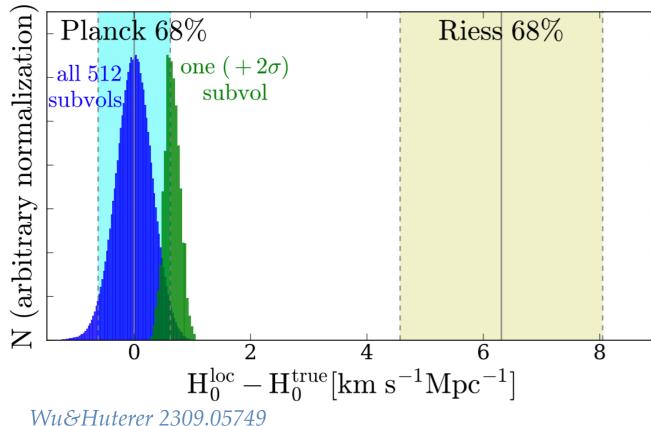
# It cannot explain the Hubble tension



*Camarena++ 2205.05422* 

## Sample variance cannot explain the Hubble tension

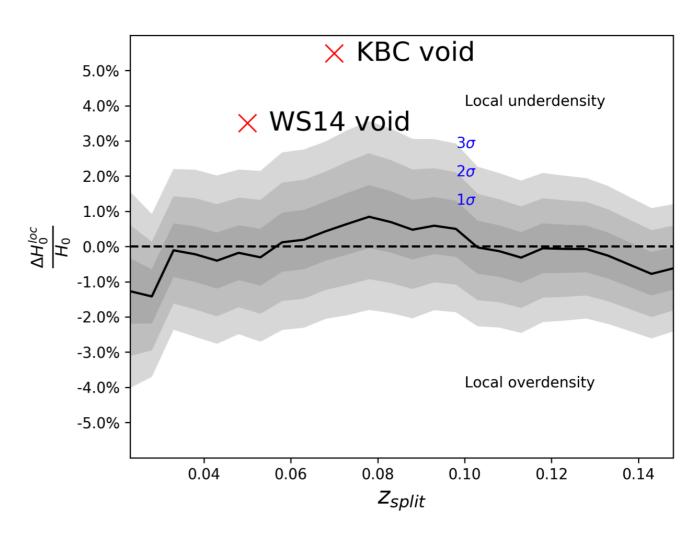
$$\Delta H_0^{\text{loc}} = \frac{1}{N} \sum_{i=1}^{N} \frac{v_{r,i}}{r_i}$$



1011111111111 2309.03749

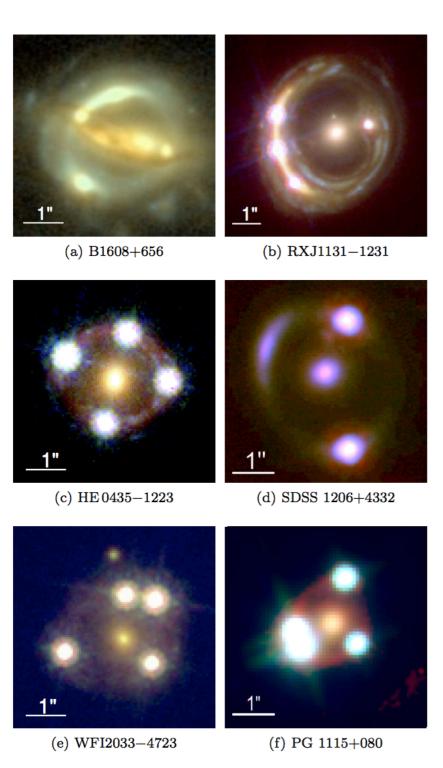
$$\sigma_{\text{sample variance}} = 0.31 \, \text{km s}^{-1} \, \text{Mpc}^{-1}$$

$$|H_0^{\rm Planck} - H_0^{\rm loc}|/\sigma_{\rm sample\,variance} \simeq 20.$$

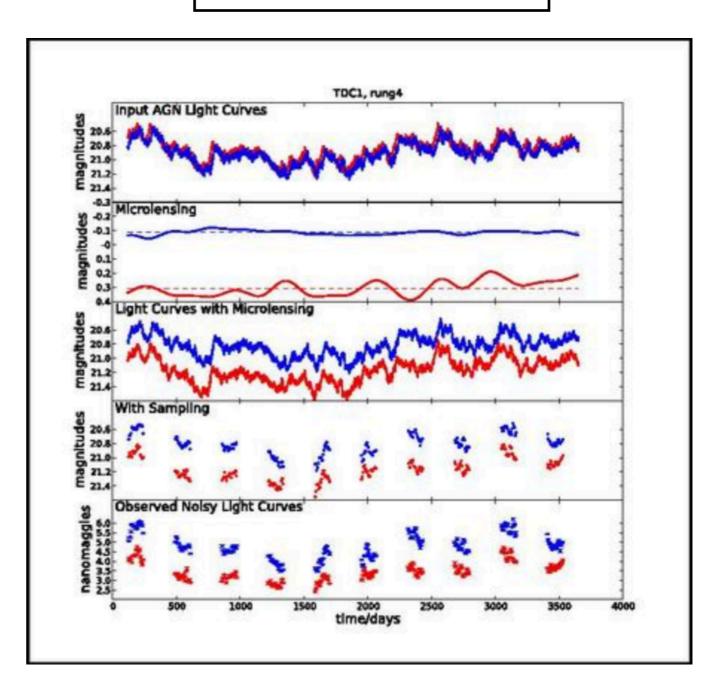


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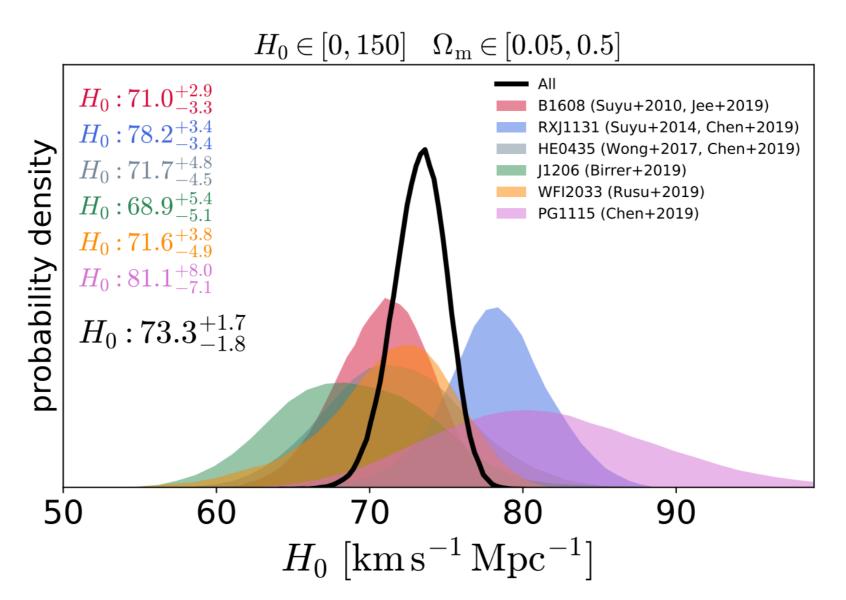
# Strongly lensed Quasars measurement of HO



$$D_{\Delta t} \equiv (1+z_{
m d}) rac{D_{
m d}D_{
m s}}{D_{
m ds}}$$

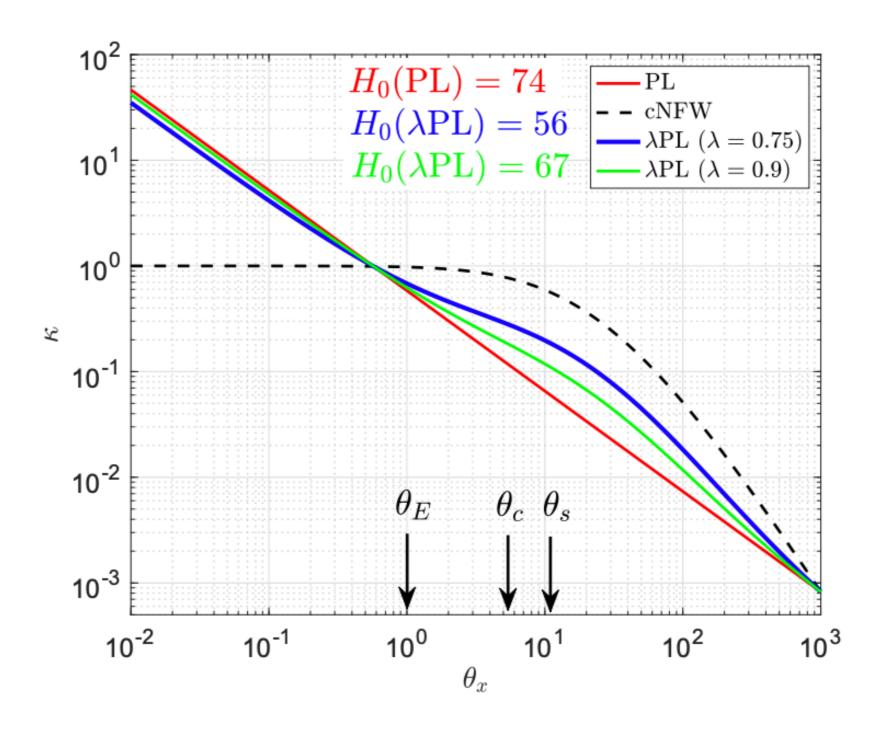


#### H0LiCOW: *H*<sub>0</sub> measurement to few %

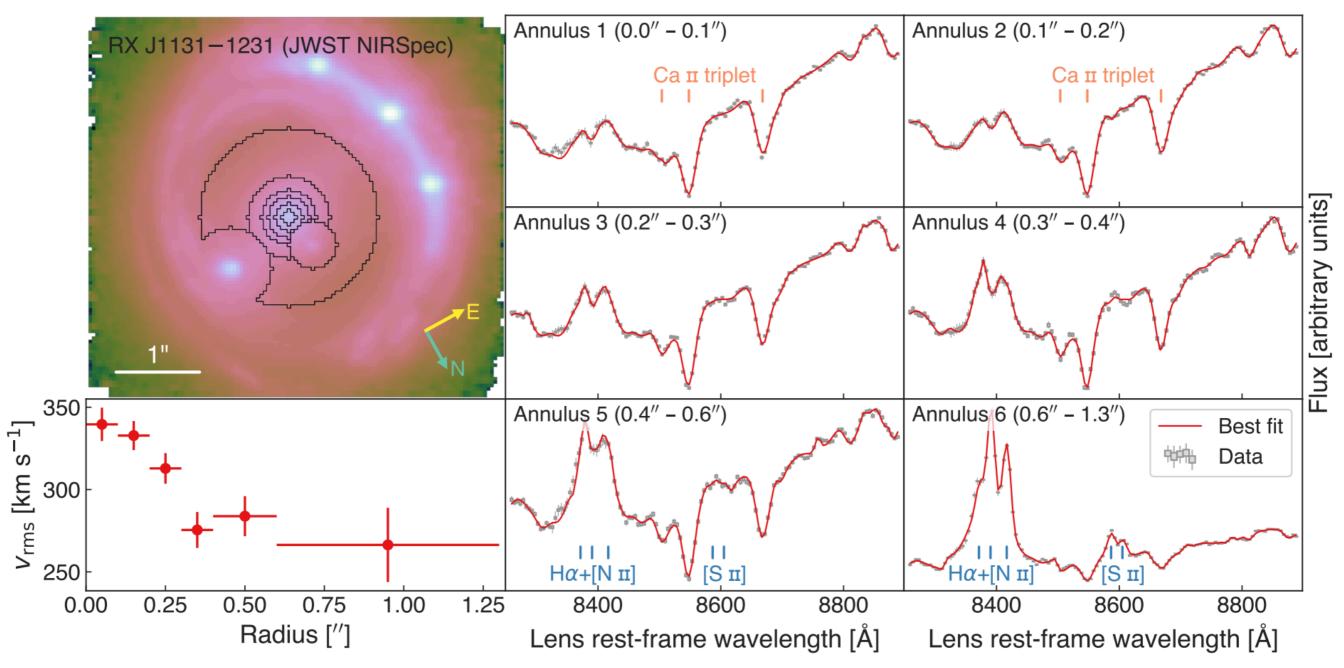


Wong et al. 1907.04869

## The "mass-sheet" degeneracy

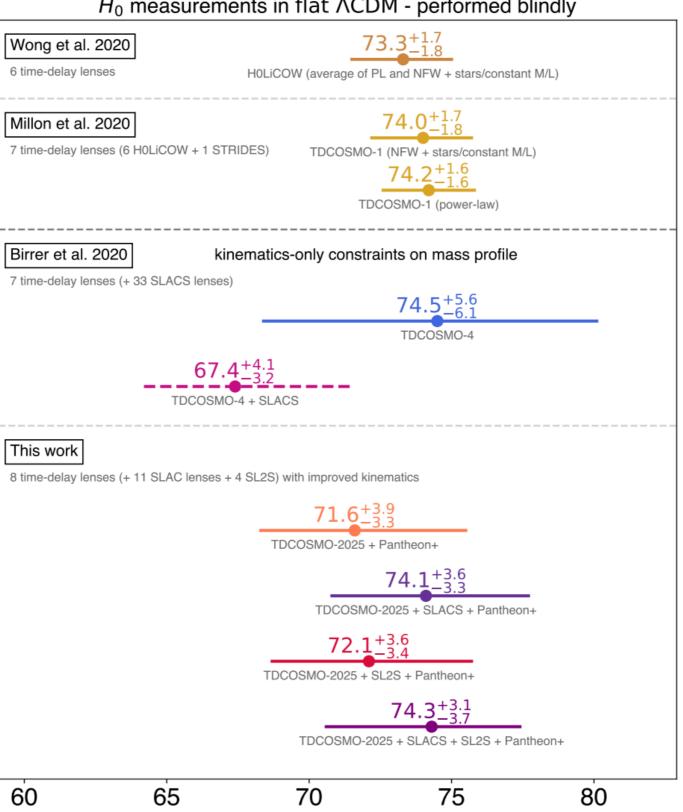


#### TDCosmo: Conservative H0 measurement



## High $H_0$ but large error bars

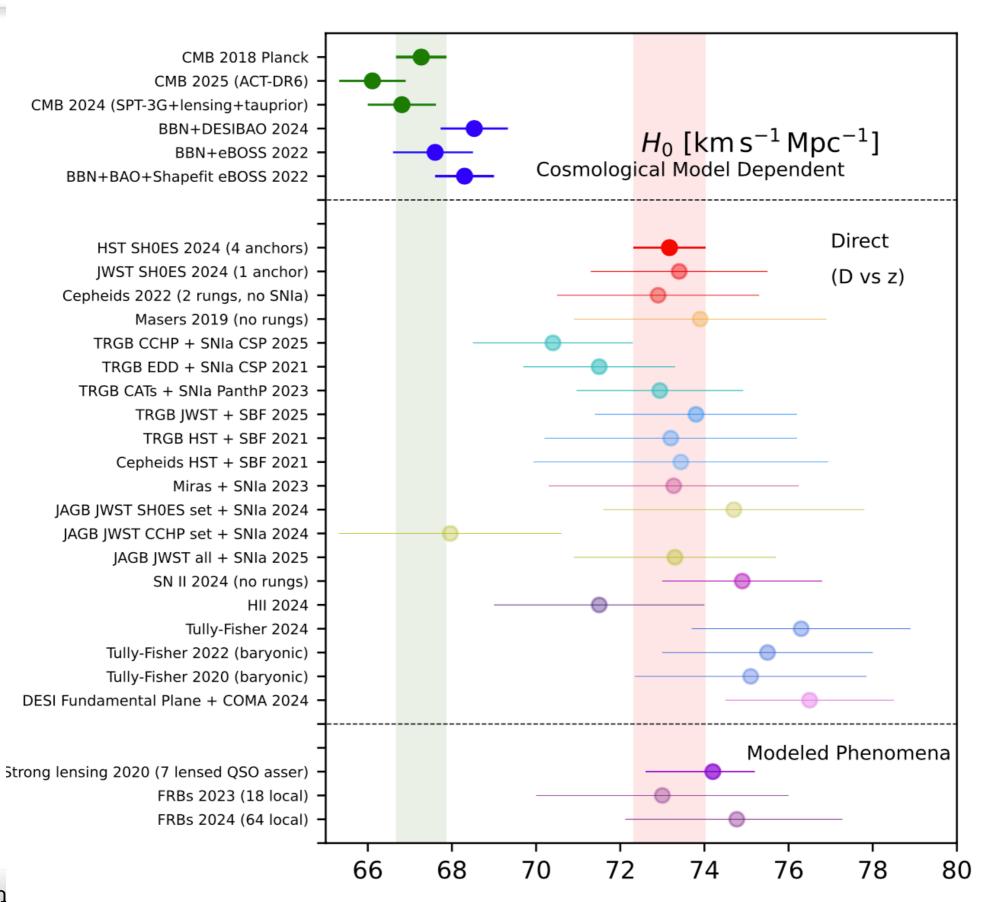
#### $H_0$ measurements in flat $\Lambda$ CDM - performed blindly



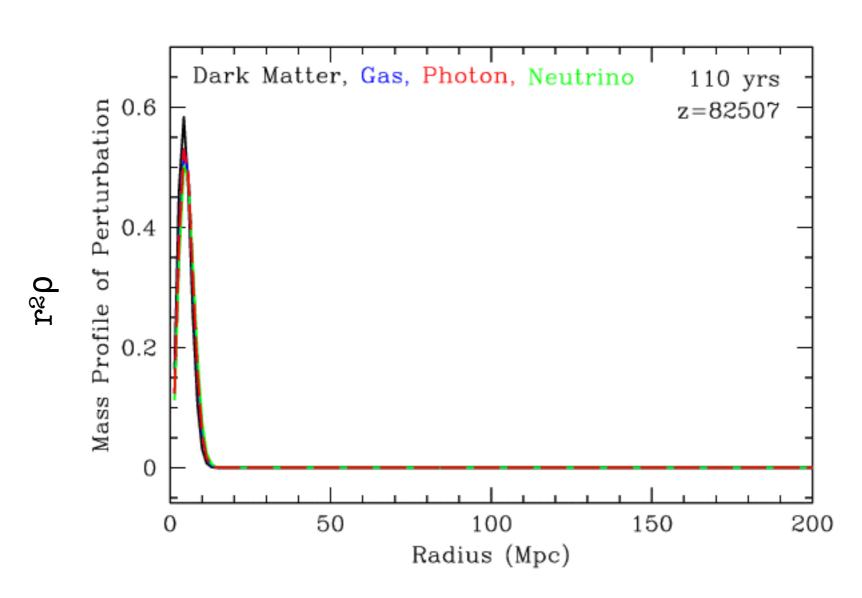
Birrer et al 2506.03023

 $H_0 [\text{km s}^{-1} \text{Mpc}^{-1}]$ 

#### The Hubble tension



- Before recombination (380 000 yr): a tightly coupled photon-baryon plasma with acoustic waves.
- The 'sound horizon' is the distance travelled by sound wave until recombination.

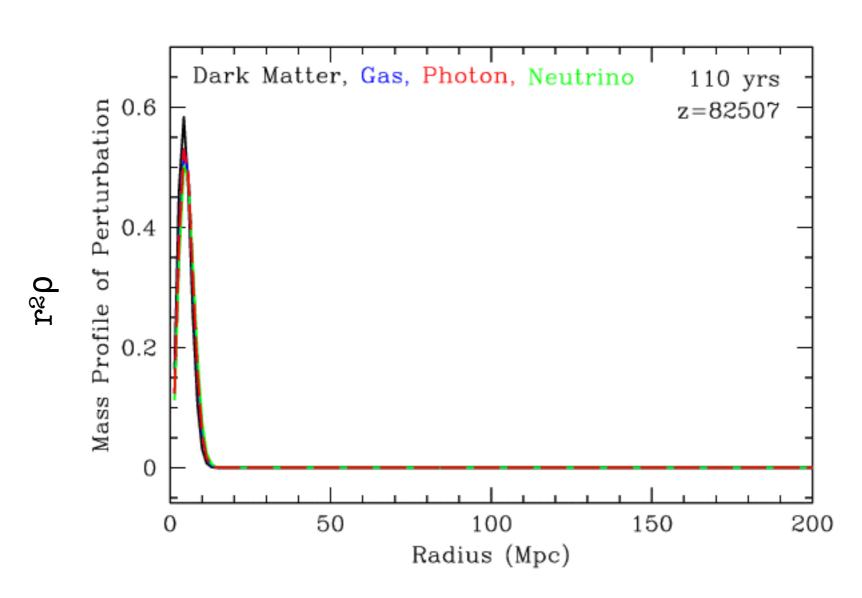


$$r_{s} = \int_{t_{rh}}^{t_{rec}} c_{s}(t)dt$$

credit: Bassett & Hlozek

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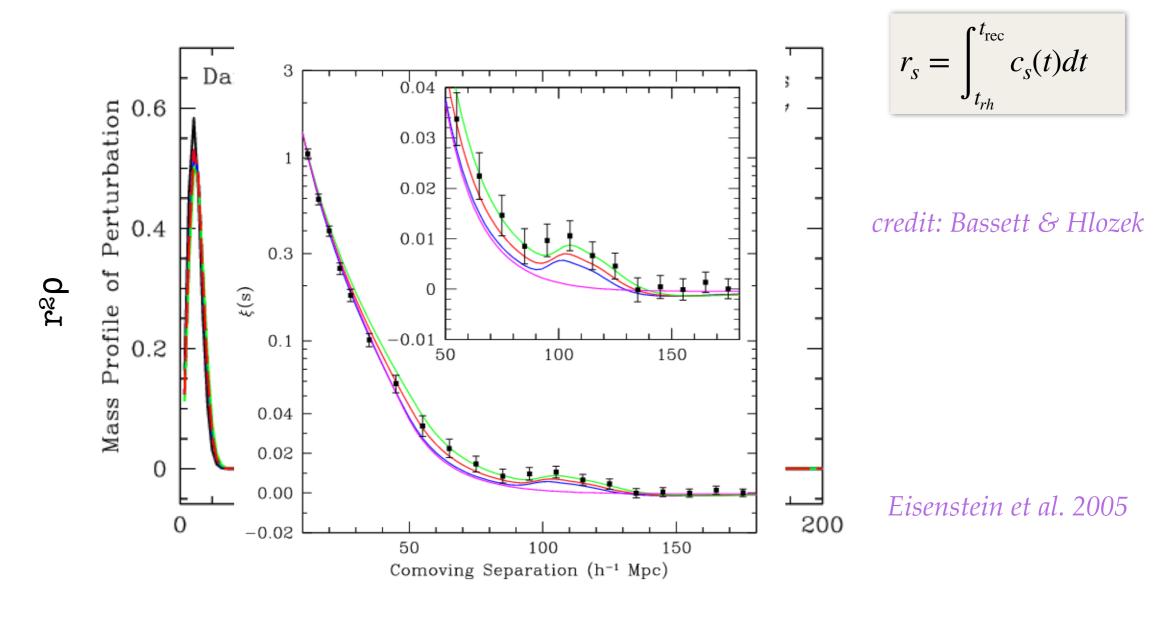


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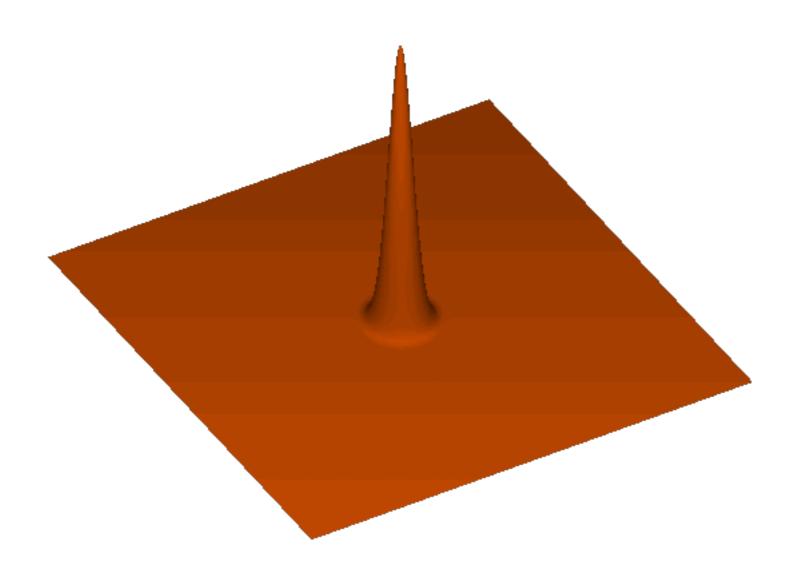
credit: Bassett & Hlozek

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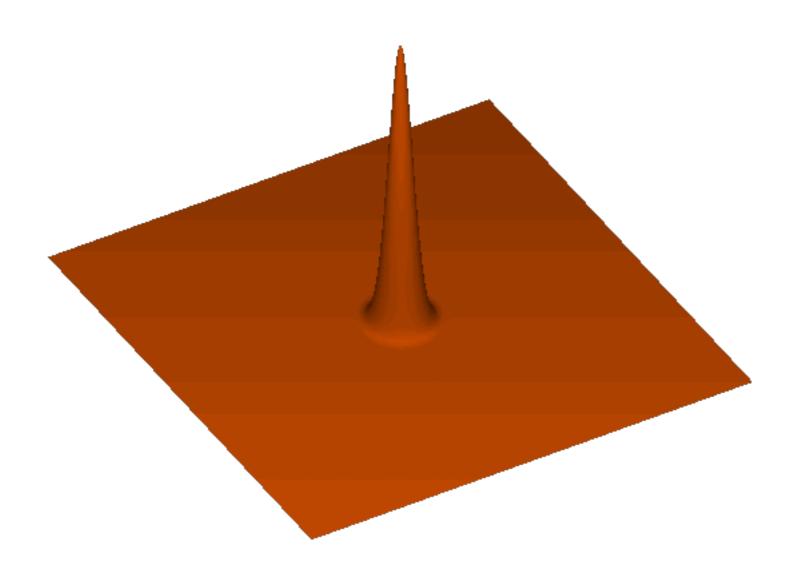
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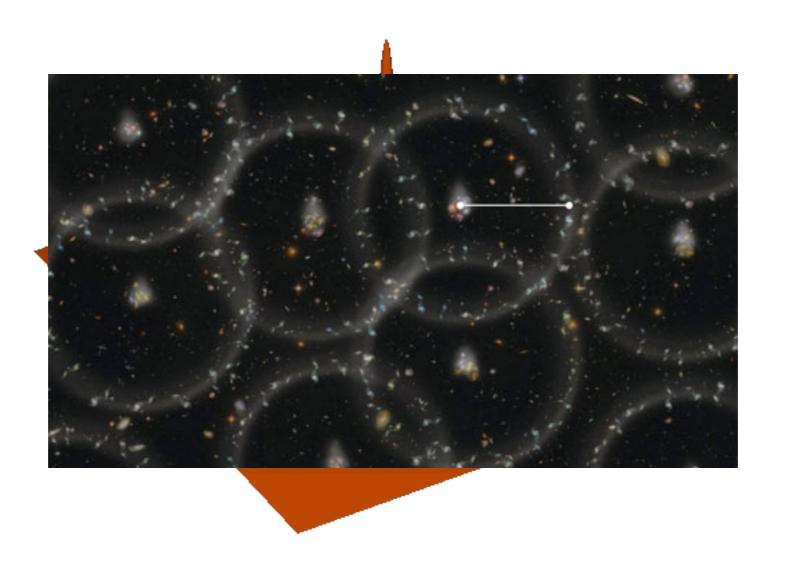
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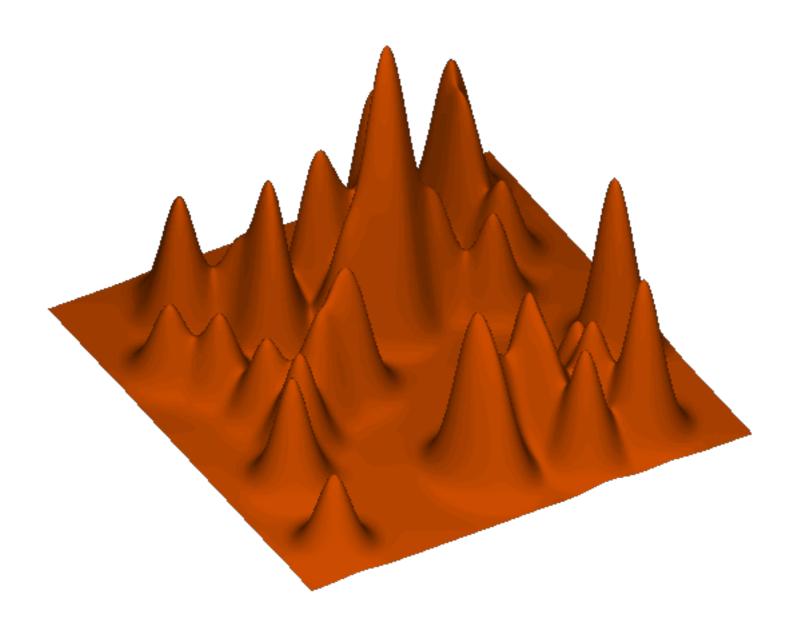
credit: Bassett & Hlozek



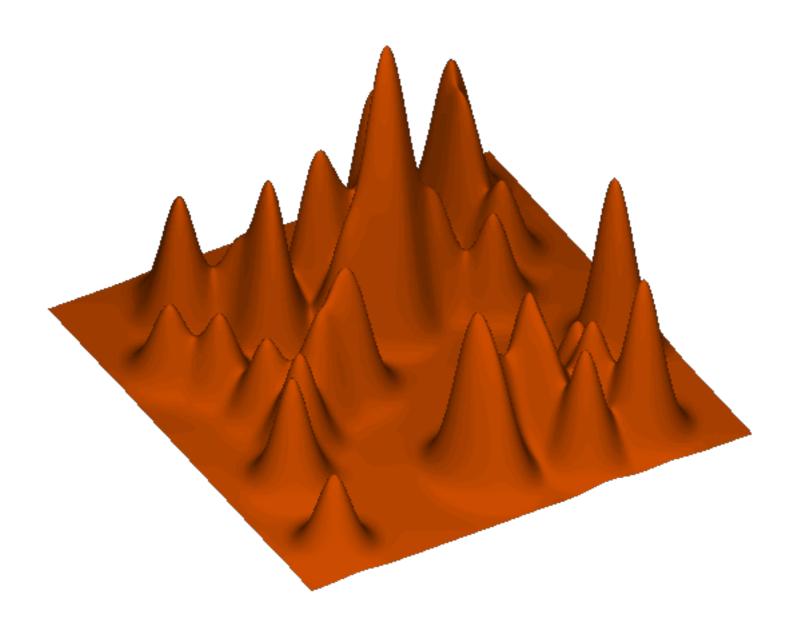
credit: Bassett & Hlozek



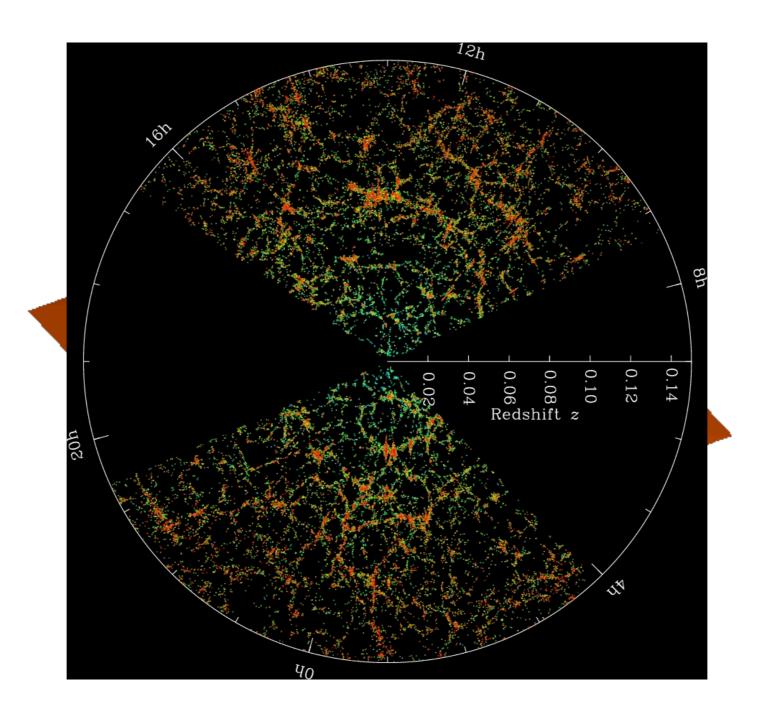
credit: Bassett & Hlozek



credit: Bassett & Hlozek

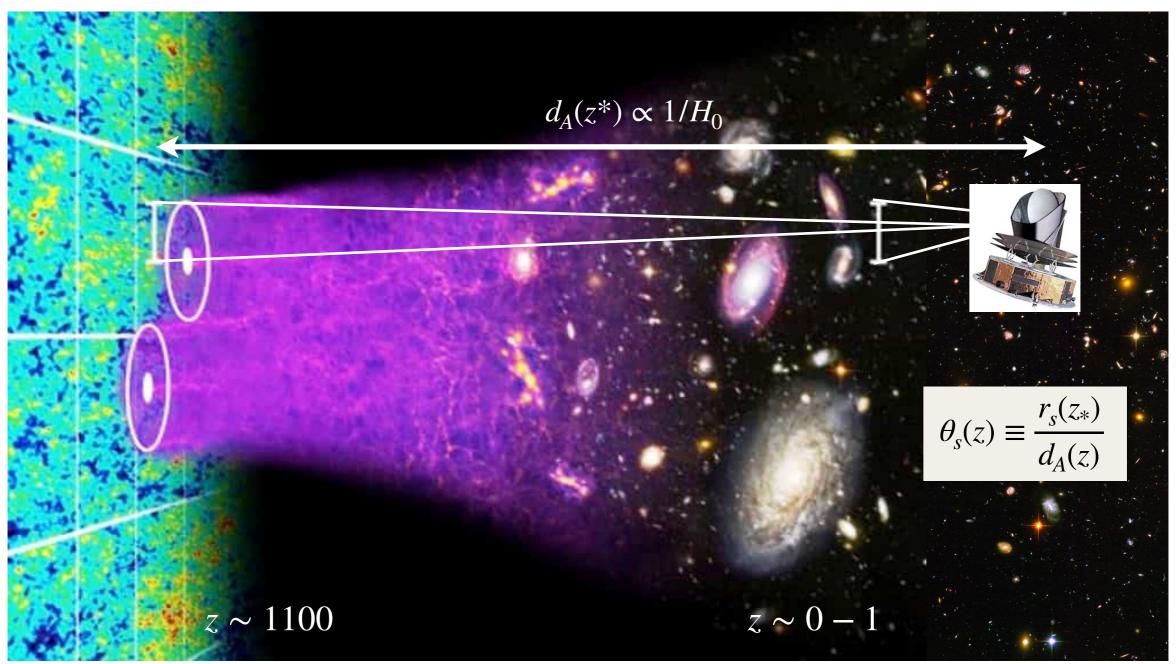


credit: Bassett & Hlozek



credit: SDSS collaboration credit: Bassett & Hlozek

The Baryonic Acoustic Oscillation: a standard ruler in the sky

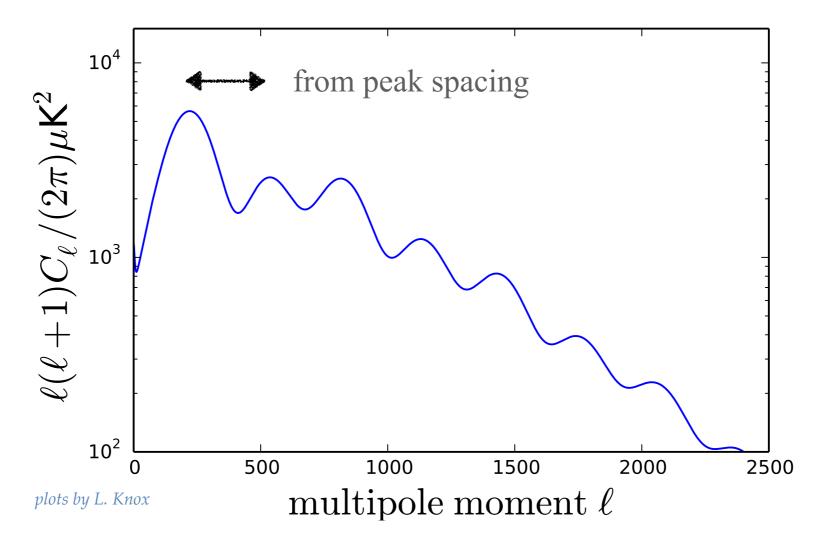


- Planck measures  $\theta_s$  at 0.04% precision but  $r_s \& d_A$  are model dependent.
- ullet  $H_0$  appears only in the angular diameter distance  $d_A$ . Summary of other measurements: Verde++ 2311.13305

#### How does CMB data measure H0?

• Inference of  $H_0$  comes from the measurement of three angular scales  $\theta_{s}$ ,  $\theta_{d}$ ,  $\theta_{eq}$ .

#### $\theta_s$ sound horizon at last scattering ~1.0404

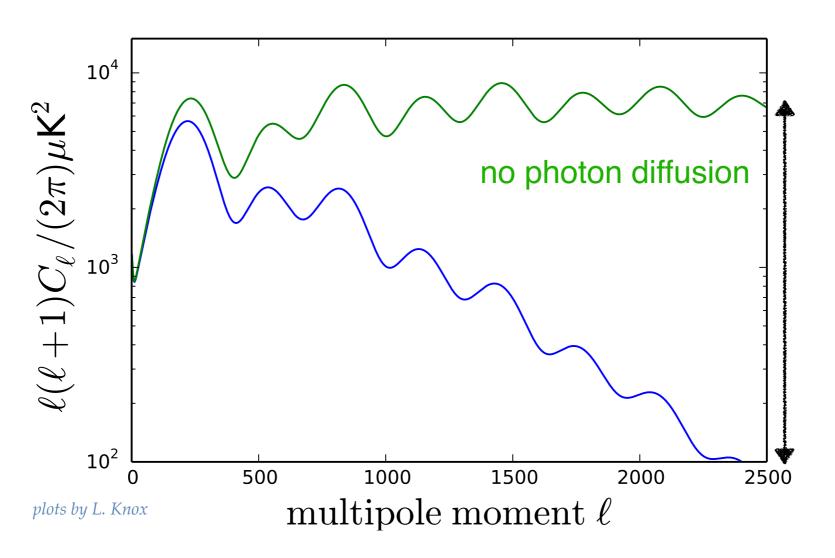


e.g. Hu&White astro-ph/9609079, Hu++astro-ph/0006436

#### How does CMB data measure H0?

• Inference of  $H_0$  comes from the measurement of three angular scales  $\theta_{s}$ ,  $\theta_{d}$ ,  $\theta_{eq}$ .

#### $\theta_d$ photon diffusion length at last scattering $\sim 0.1609$



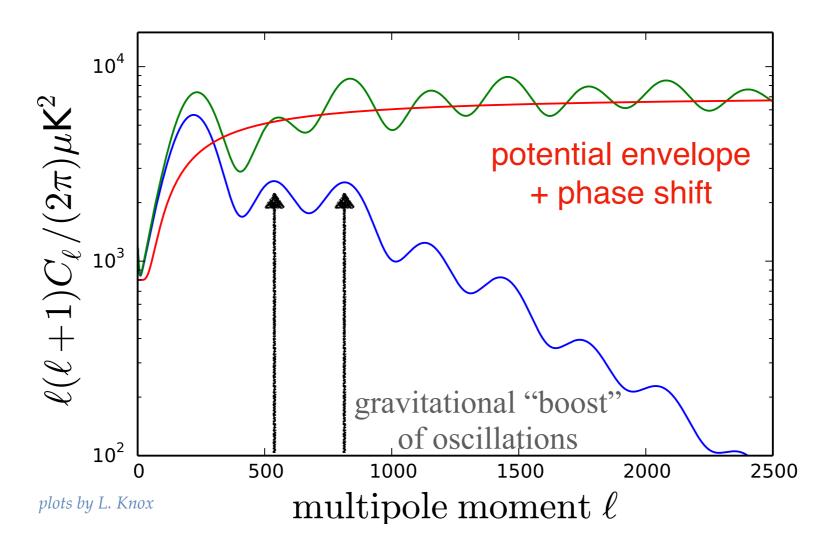
"Silk Damping"

e.g. Hu&White astro-ph/9609079, Hu++astro-ph/0006436

#### How does CMB data measure H0?

• Inference of  $H_0$  comes from the measurement of three angular scales  $\theta_{s}$ ,  $\theta_{d}$ ,  $\theta_{eq}$ .

#### $\theta_{eq}$ horizon size at matter-radiation equality $\sim 0.81$

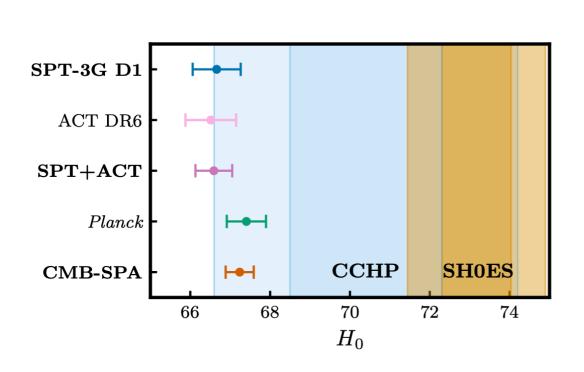


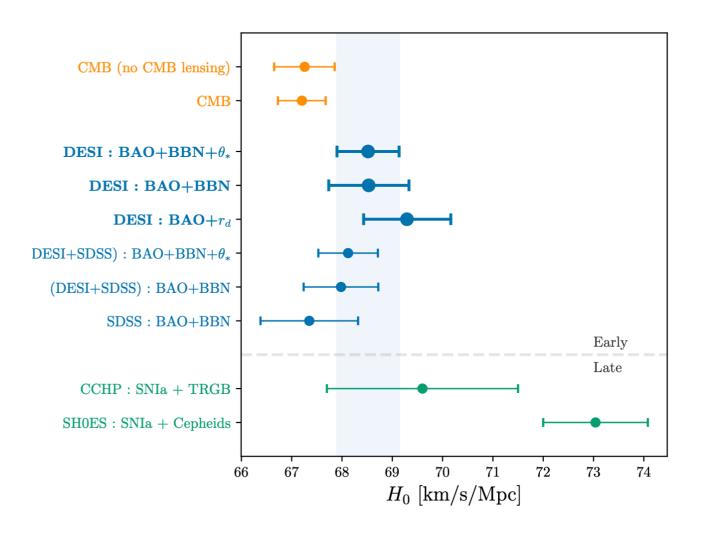
e.g. Hu&White astro-ph/9609079, Hu++astro-ph/0006436

## CMB and BAO measurement of $H_0$

#### **Comparing CMB measurements**

#### **Comparing BAO+BBN measurements**

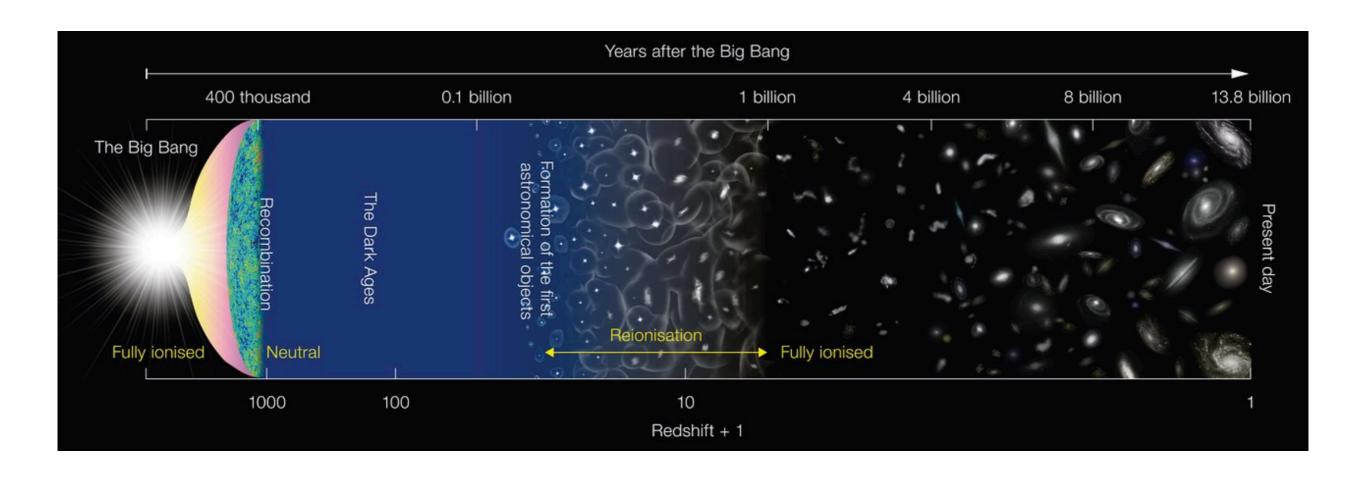




• The measurement of  $H_0$  in the  $\Lambda$ CDM model using the inverse-distance ladder is robust

## Two ways of solving the Hubble tension

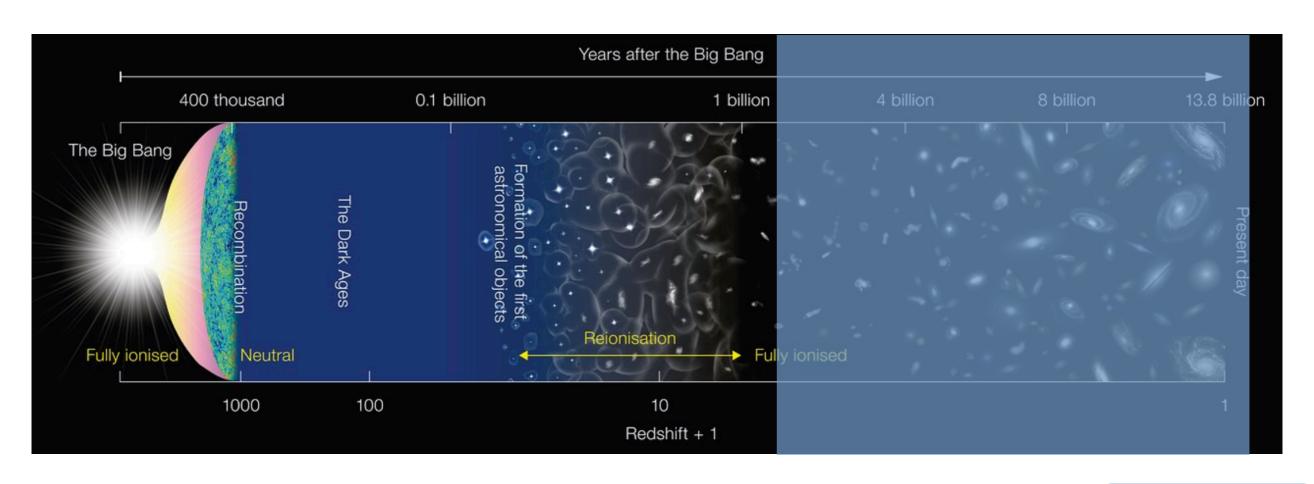
$$\theta_s = \frac{r_s}{r_A} = \frac{H_0 r_s}{\int_0^{z_*} 1/E(x) \mathrm{d}x}$$



## Two ways of solving the Hubble tension

$$\theta_s = \frac{r_s}{r_A} = \frac{H_0 r_s}{\int_0^{z_*} 1/E(x) \mathrm{d}x}$$

#### Late-universe models



$$\frac{H_0 \nearrow r_s}{\int_0^{z_*} 1/E(x) \searrow \mathrm{d}x}$$

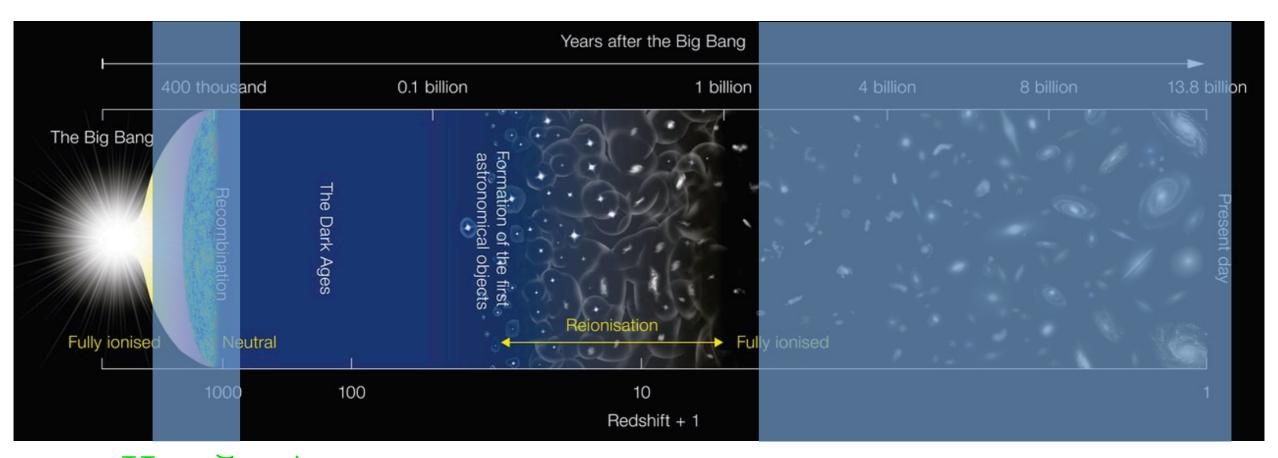
Change expansion history

## Two ways of solving the Hubble tension

**Early universe models** 

$$\theta_s = \frac{r_s}{r_A} = \frac{H_0 r_s}{\int_0^{z_*} 1/E(x) \mathrm{d}x}$$

#### Late-universe models



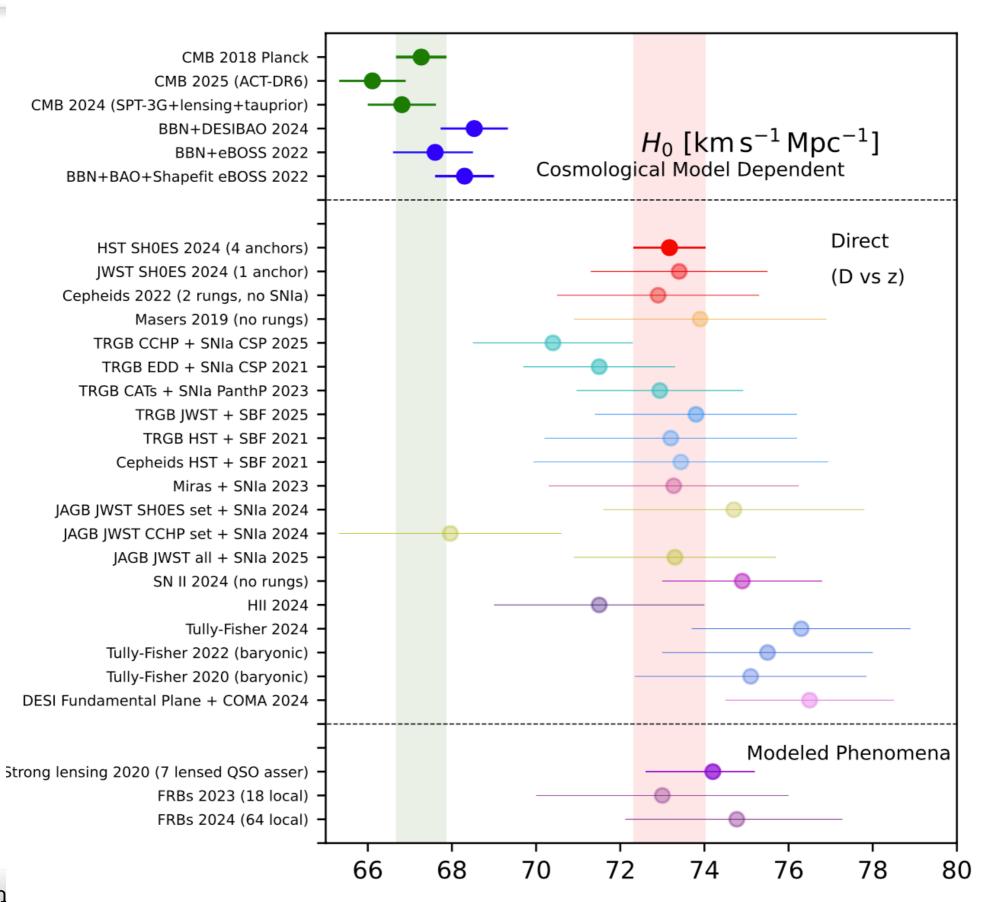
$$\frac{H_0 \nearrow r_s}{\int_0^z 1/E(x) dx}$$

**Change calibrator** 

$$\frac{H_0 \nearrow r_s}{\int_0^{z_*} 1/E(x) \searrow \mathrm{d}x}$$

Change expansion history

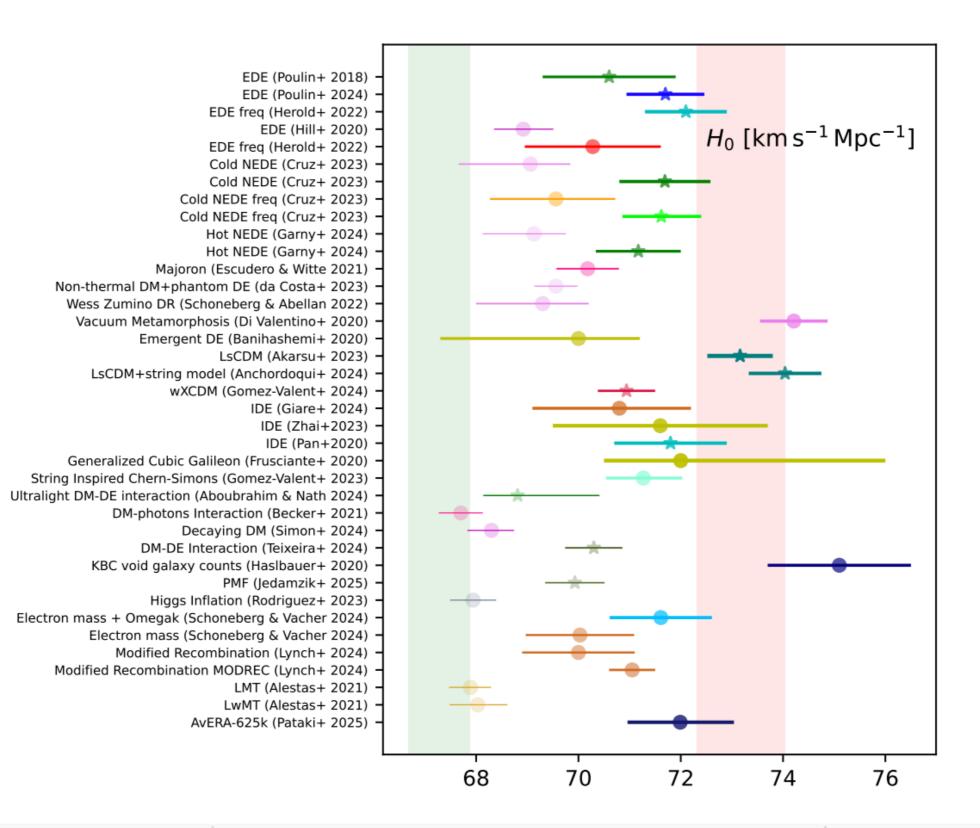
#### The Hubble tension



- ★ Planck+BAO+SNIa+SH0ES
- ★ Planck+DESI+SNIa PP+SH0ES
- ★ Planck+BAO+SH0ES
- Planck+BAO+SNIa+S8
- Planck+BAO+S8
- Planck+BAO+SNIa
- Planck+BAO+SNIa+BBN

- ★ Planck+BAO+BBN+SNIa+SH0ES
- Planck+BAO
- Planck
- ★ Planck+BAOtr+KiDS1000+SNIa+SH0ES
- ★ Planck+BAOtr+DESY5 SN+CC+fσ<sub>12</sub>+SH0ES
- Planck+DESI

- Planck+DESI+SNIa+SH0ES
- Planck+B2K+BAO+SNIa
- Planck+DESI+SNIa
- Planck+BAO+RSD+SNIaPlanck+BAO+CC+SNIa+SH0ES
- SNIa



#### The $H_0$ Olympics: A fair ranking of proposed models

Nils Schöneberg<sup>a,\*</sup>, Guillermo Franco Abellán<sup>b</sup>, Andrea Pérez Sánchez<sup>a</sup>, Samuel J. Witte<sup>c</sup>, Vivian Poulin<sup>b</sup>, Julien Lesgourgues<sup>a</sup>

#### Abstract

Despite the remarkable success of the  $\Lambda$  Cold Dark Matter ( $\Lambda$ CDM) cosmological model, a growing discrepancy has emerged (currently measured at the level of  $\sim 4-6\,\sigma$ ) between the value of the Hubble constant  $H_0$  measured using the local distance ladder and the value inferred using the cosmic microwave background and galaxy surveys. While a vast array of  $\Lambda$ CDM extensions have been proposed to explain these discordant observations, understanding the (relative) success of these models in resolving the tension has proven difficult – this is a direct consequence of the fact that each model has been subjected to differing, and typically incomplete, compilations of cosmological data. In this review, we attempt to make a systematic comparison of seventeen different models which have been proposed to resolve the  $H_0$  tension (spanning both early- and late-Universe solutions), and quantify the relative success of each using a series of metrics and a vast array of data combinations. Owing to the timely appearance of this article, we refer to this contest as the " $H_0$  Olympics"; the goal being to identify which of the proposed solutions, and more broadly which underlying mechanisms, are most likely to be responsible for explaining the observed discrepancy (should unaccounted for systematics not be the culprit). This work also establishes a foundation of tests which will allow the success of novel proposals to be meaningfully "benchmarked".

Keywords: Hubble Tension, Dark Energy, Dark Matter Phenomenology, Dark Radiation, Early Dark Energy, Varying fundamental constants

<sup>&</sup>lt;sup>a</sup>Institute for Theoretical Particle Physics and Cosmology (TTK), RWTH Aachen University, D-52056 Aachen, Germany.

<sup>&</sup>lt;sup>b</sup>Laboratoire Univers & Particules de Montpellier (LUPM), CNRS & Université de Montpellier (UMR-5299), Place Eugène Bataillon, F-34095 Montpellier Cedex 05, France.

<sup>&</sup>lt;sup>c</sup>GRAPPA Institute, Institute for Theoretical Physics Amsterdam and Delta Institute for Theoretical Physics, University of Amsterdam, Science Park 904, 1098 XH Amsterdam, The Netherlands

## The $H_0$ olympics: fairly ranking models

• We compare 17 different models suggested to resolve the Hubble tension.

Schöneberg (VP) ++ 2107.10291

$$\frac{H_0 \nearrow r_s}{\int_0^{z_*} 1/E(x) \searrow \mathrm{d}x}$$

$$\int_0^z 1/E(x)$$

Late-Universe models

- CPL Dark Energy
- Emergent Dark Energy
- Generalized Emergent Dark Energy
- Decaying Dark matter to massless particles
- Decaying Dark matter to massive particles

- Dark Radiation models
- Free-streaming  $N_{\text{eff}}$
- Self-Interacting  $N_{\rm DR}$
- Mixture of  $N_{\text{eff}} + N_{\text{DR}}$
- DM-DR interaction +  $N_{\rm DR}$
- Self interacting  $\nu + N_{\text{eff}}$

- Exotic Early Universe models
  - Early Dark Energy
  - New Early Dark Energy
  - Early modified Gravity
  - Primordial magnetic fields
  - Varying electron mass  $m_e$
  - Varying electron mass  $m_e + \Omega_k$

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$\frac{H_0 \nearrow r_s}{\int_0^{z_*} 1/E(x) \searrow \mathrm{d}x}$	$\frac{H_0 \nearrow r_s}{\int_0^z 1/E(x) \mathrm{d}x}$	
Late-Universe models	Dark Radiation models	Exotic Early Universe models
<ul> <li>CPL Dark Energy</li> </ul>	• Free-streaming $N_{\text{eff}}$	<ul><li>Early Dark Energy</li></ul>
<ul><li>Emergent Dark Energy</li></ul>	• Self-Interacting $N_{\rm DR}$	New Early Dark Energy
<ul><li>Generalized Emergent Dark Energy</li></ul>	• Mixture of $N_{\text{eff}} + N_{\text{DR}}$	<ul><li>Early modified Gravity</li></ul>
<ul><li>Decaying Dark matter to</li></ul>	• DM-DR interaction + $N_{\rm DR}$	<ul><li>Primordial magnetic fields</li></ul>
massless particles	• Self interacting $\nu + N_{\rm eff}$	• Varying electron mass $m_e$
<ul> <li>Decaying Dark matter to massive particles</li> </ul>		• Varying electron mass $m_e + \Omega_k$

- I/ Frequentist: Can a model give a good fit to all data including SH0ES and be favored over ΛCDM.
- II/ Bayesian: Can a model be favored over  $\Lambda$ CDM independently of SH0ES and "predict" a high  $H_0$ .

Assumed from LCDM

Measured

$$\theta_s \equiv \frac{r_s(z_*)}{d_A(z_*)}$$

$$\theta_{s} \equiv \frac{r_{s}(z_{*})}{d_{A}(z_{*})} \qquad d_{A}(z) \equiv \int_{0}^{z} \frac{dz'}{H_{0}\sqrt{\Omega_{m}(1+z)^{3} + \Omega_{\Lambda}(1+z)^{3(1+w)} + \cdots}}$$

Assumed from LCDM Measured

$$\theta_{s} \equiv \frac{r_{s}(z_{*})}{d_{A}(z_{*})} \qquad d_{A}(z) \equiv \int_{0}^{z} \frac{dz'}{H_{0}\sqrt{\Omega_{m}(1+z)^{3} + \Omega_{\Lambda}(1+z)^{3(1+w)} + \cdots}} \qquad H_{0} \uparrow \Rightarrow \Omega_{X}(z) \downarrow$$

• 'phantom dark energy' w < -1, DE-DM interactions, decaying DM, and many more...

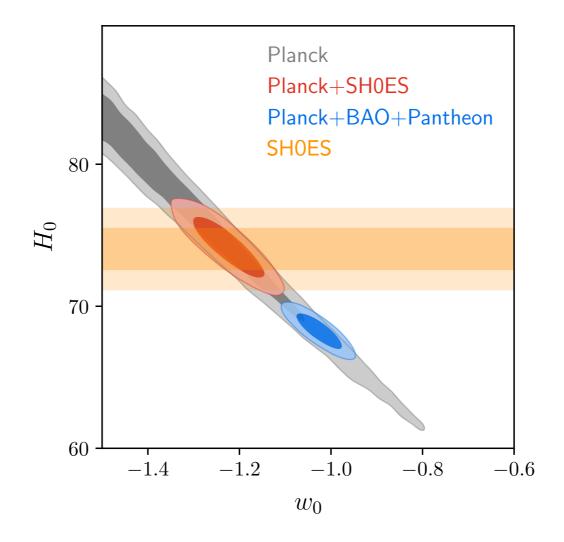
[http://arxiv/insert\_your\_favorite\_ model\_here.com]

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• Planck data can easily accommodate a higher  $H_0$ : problem with BAO and Pantheon

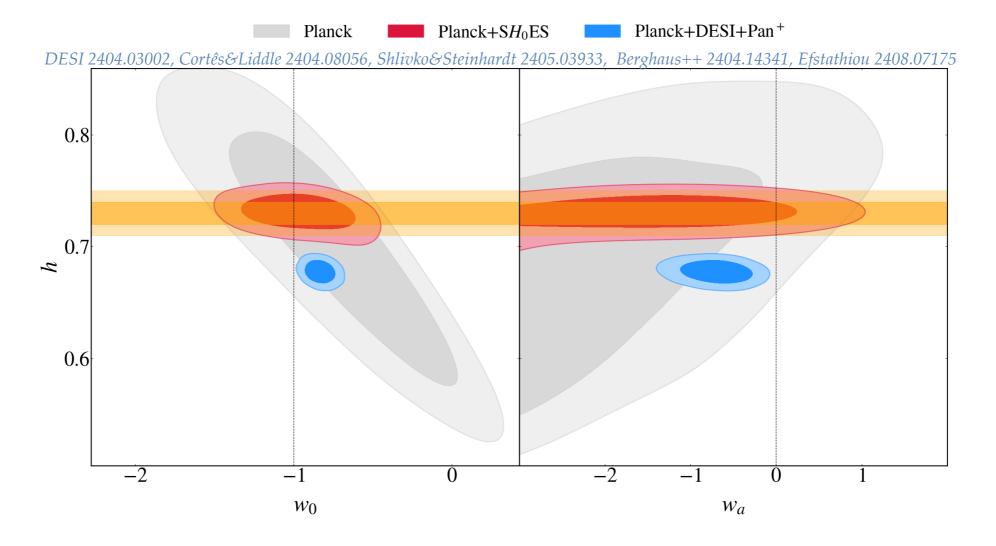


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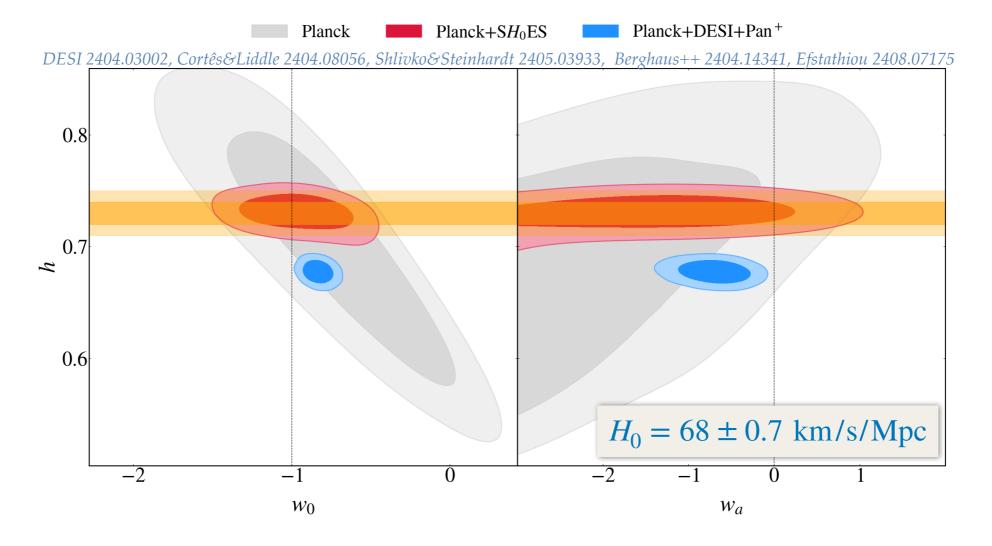


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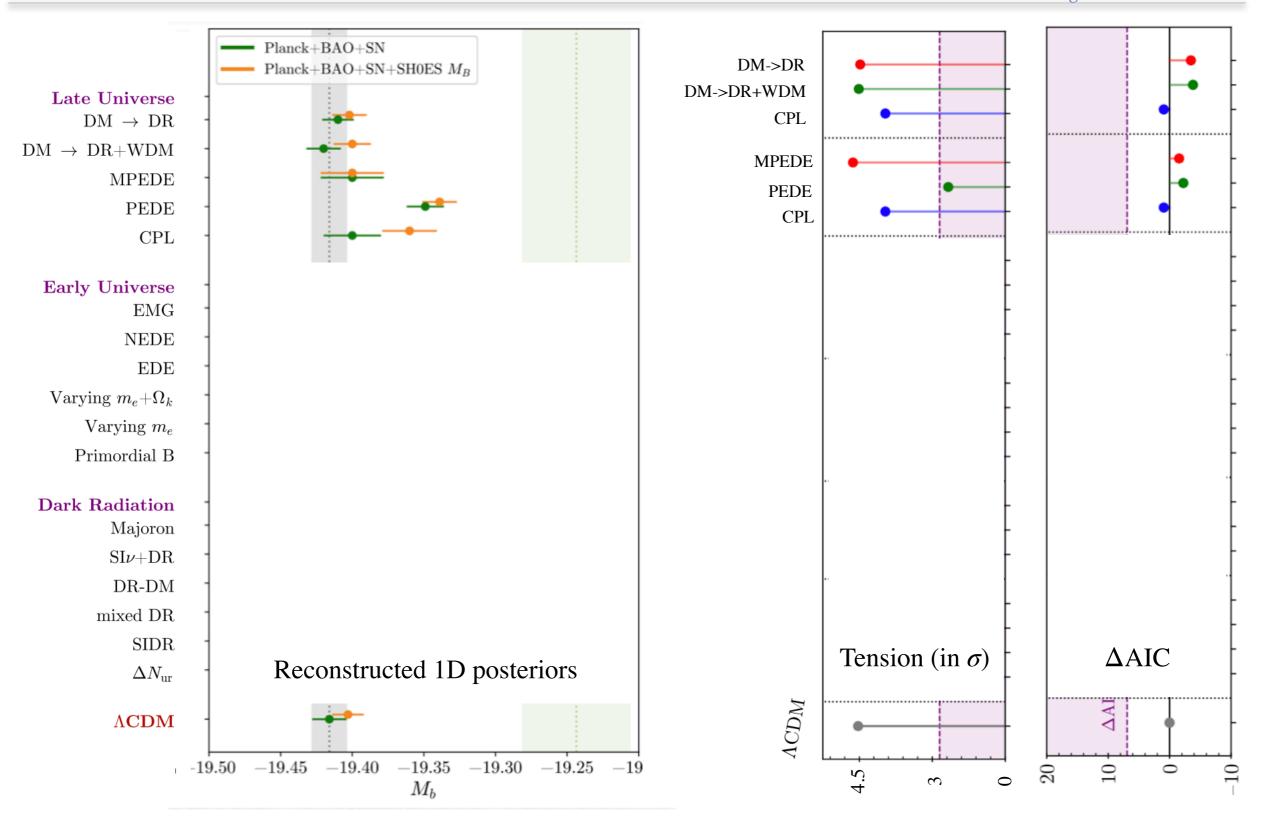
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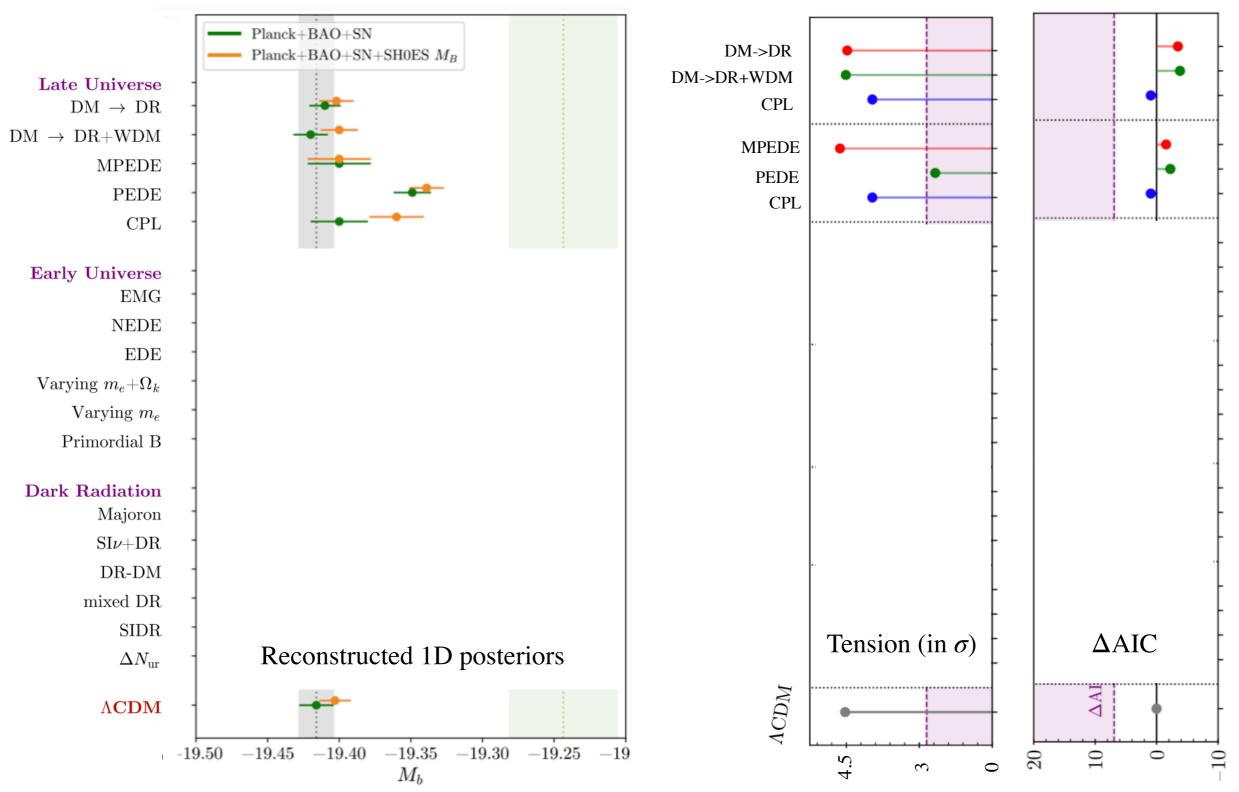
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# The $H_0$ olympics: fairly ranking models Schöneberg (VP) ++ 2107.10291



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BAO: 
$$\theta_d(z) = \frac{r_s(z_{\text{drag}})}{D_A(z)}$$

SN1a: 
$$m(z) = 5 \log_{10}(D_L(z)) + M_b$$

• GR + photon conservation imposes the "distance-duality relation":  $D_A(z) = D_L(z)/(1+z)^2$ 

BAO: 
$$\theta_d(z) = r_s(z_{\text{drag}})$$
 Planck

SN1a: 
$$m(z) = 5 \log_{10}(D_L(z)) + M_b$$

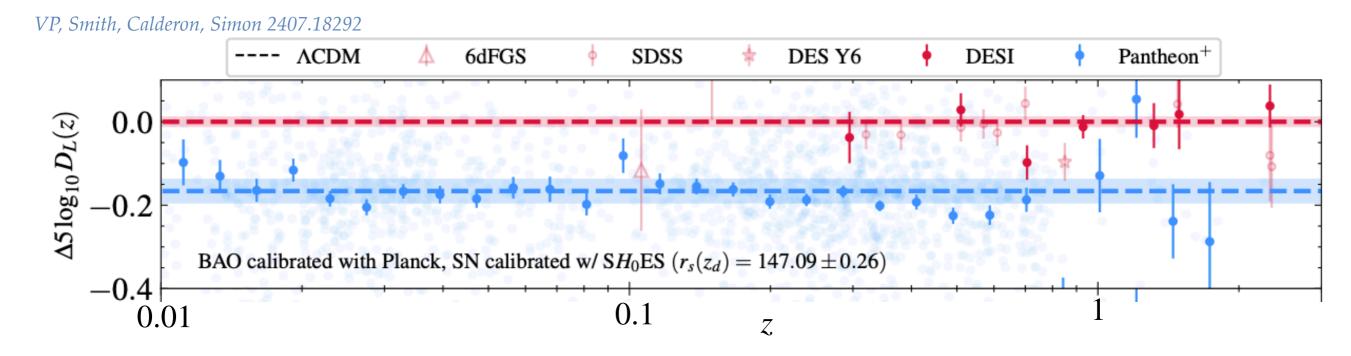
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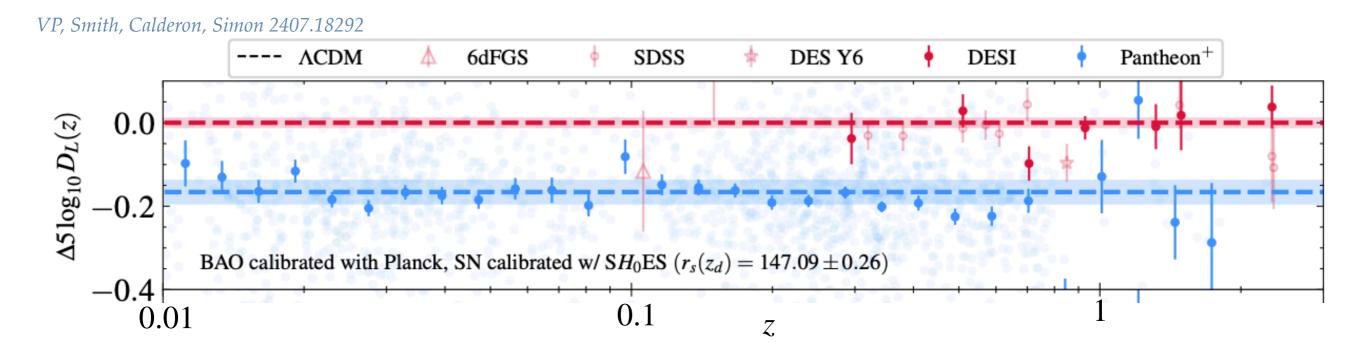


• Assuming  $r_s \sim 147$  Mpc and  $M_b \sim -19.25$ ,  $D_A(z)$  and  $D_L(z)$  are incompatible! Still true with DESI Camarena&Marra 2101.08641, Efstathiou 2103.08723, Raveri 2309.06795

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- Solving the tension require to either change calibrators or break the DDR relation

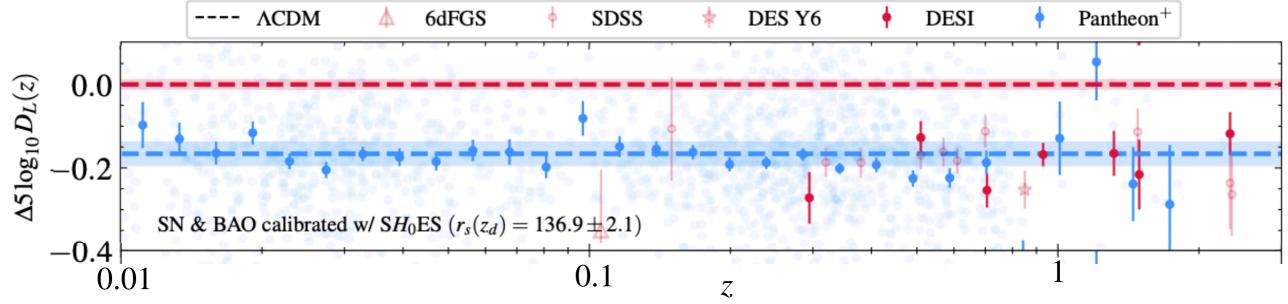
BAO: 
$$\theta_d(z) = r_s(z_{\text{drag}})$$
 Planck

SN1a: 
$$m(z) = 5 \log_{10}(D_L(z)) + M_b$$

GR + photon conservation imposes the "distance-duality relation":  $D_A(z) = D_L(z)/(1+z)^2$ 

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- Assuming  $r_s \sim 147$  Mpc and  $M_b \sim -19.25$ ,  $D_A(z)$  and  $D_L(z)$  are incompatible! Still true with DESI Camarena&Marra 2101.08641, Efstathiou 2103.08723, Raveri 2309.06795
- Solving the tension require to either change calibrators or break the DDR relation

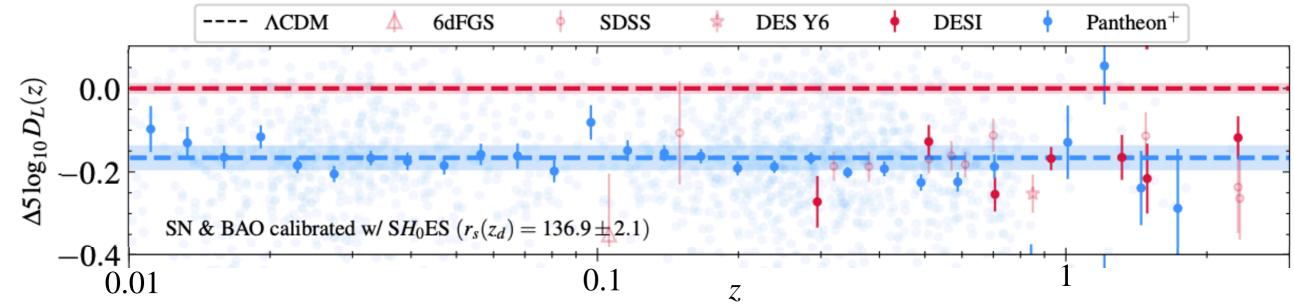
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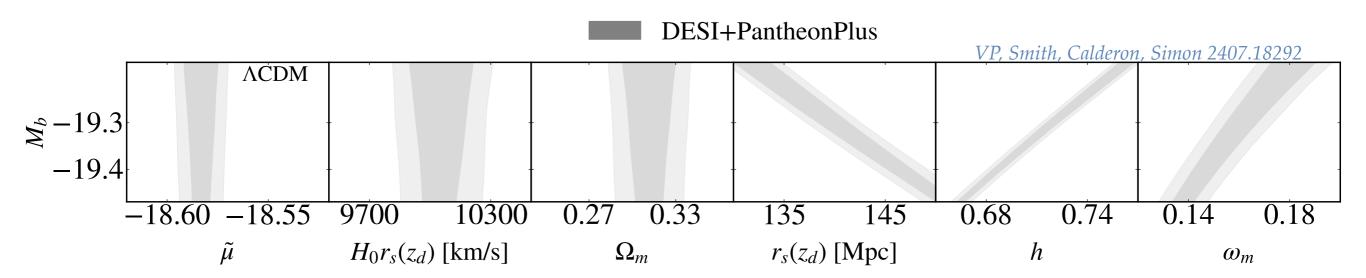
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VP, Smith, Calderon, Simon 2407.18292



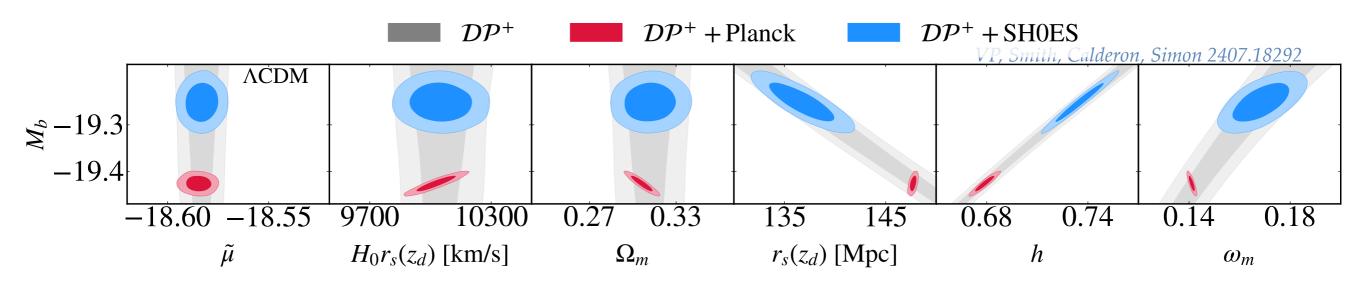
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- Solving the tension require to either change calibrators or break the DDR relation
- A single "constant" shift is currently sufficient  $\Rightarrow$  changing calibrators favored! *Teixeira* (VP) ++ 2504.10464

• What is the impact of calibrating the BAO+SN1a Hubble diagram with either Planck or SH0ES?



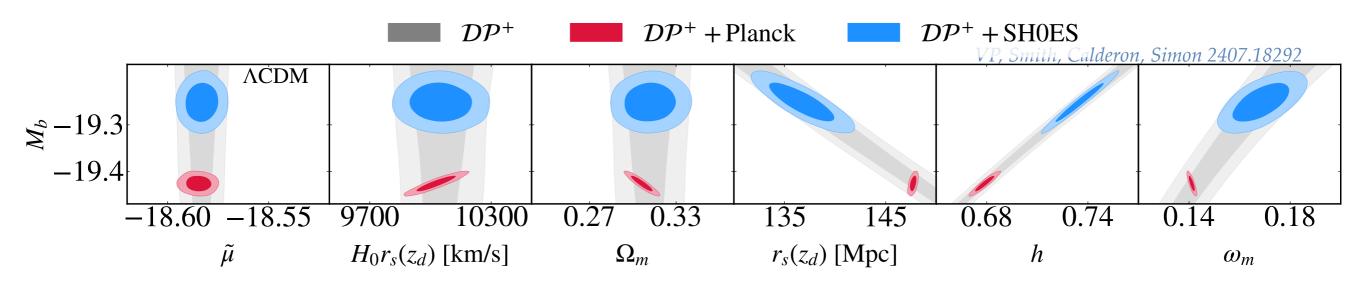
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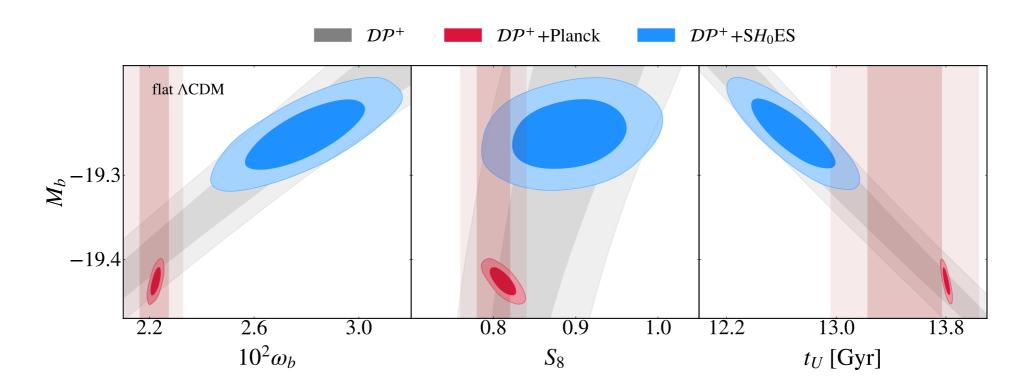
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- Challenge for new physics: Reduce the sound horizon and compensate the larger  $\omega_m$  on the CMB

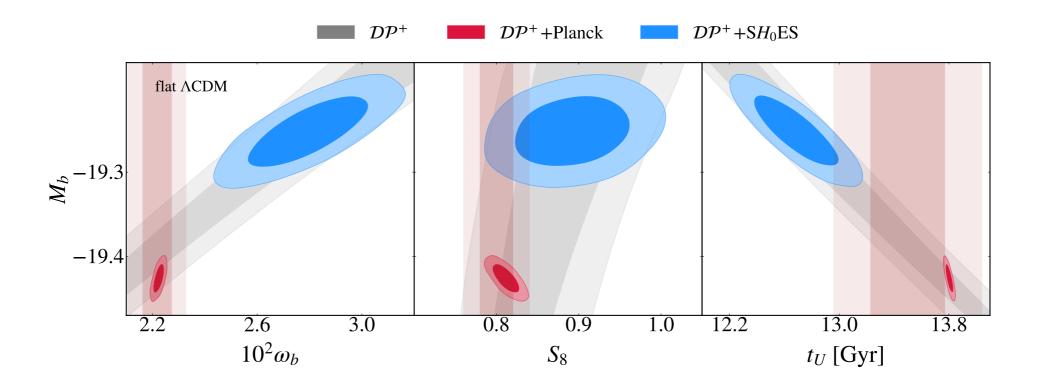
See also Jedamzik++ 2010.04158, Blanchard++ 2205.05017, Pedrotti++ 2408.04530

• Under  $\Lambda$ CDM, the S $H_0$ ES calibration has further implications



VP, Smith, Calderon, Simon 2407.18292

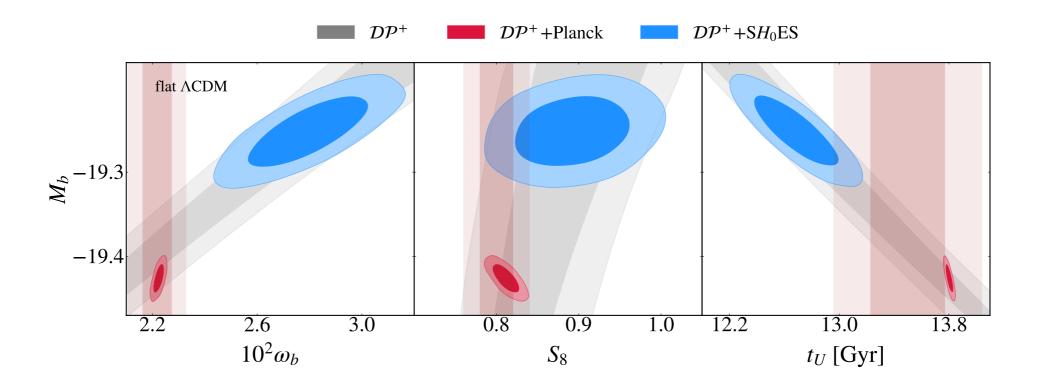
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VP, Smith, Calderon, Simon 2407.18292

**BBN** tension:  $\omega_b$  is much larger in the "S $H_0$ ES cosmology" in order to adjust the low  $r_s$ 

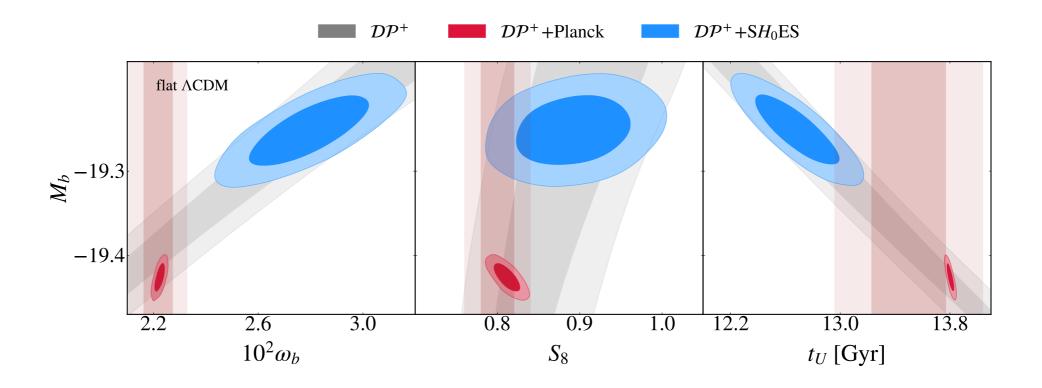
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VP, Smith, Calderon, Simon 2407.18292

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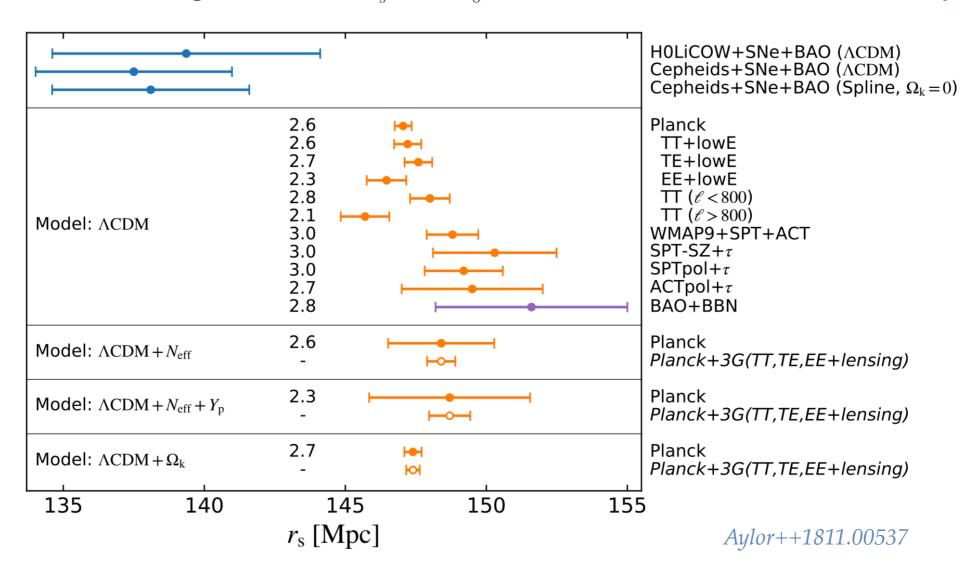


VP, Smith, Calderon, Simon 2407.18292

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- Age of the universe tension:  $t_U$  is younger by about 1 Gyr than in Planck/ $\Lambda$ CDM!

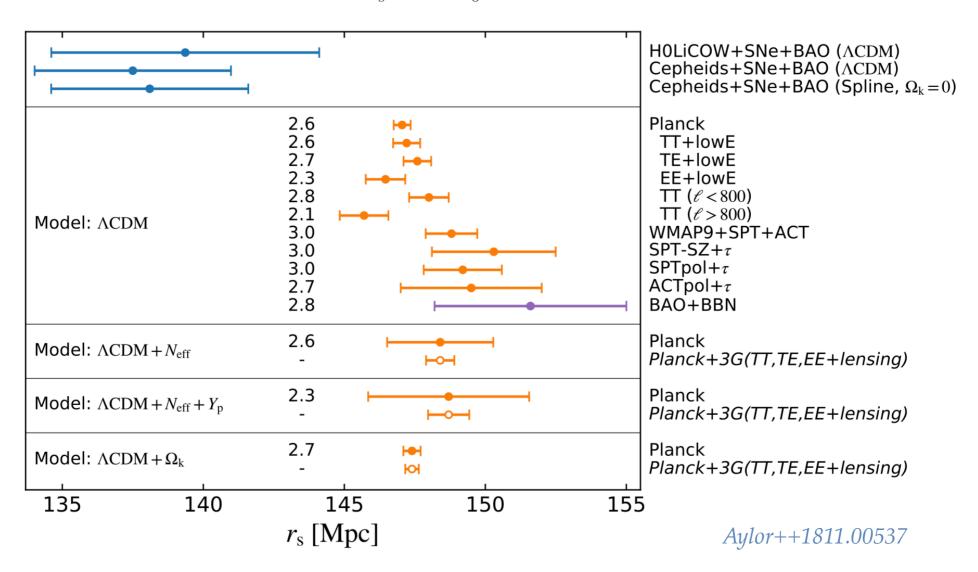
Bernal++ 2102.05066, Boylan-Kolchin 2103.15824, Vagnozzi 2105.10425

• One can deduce the co-moving sound horizon  $r_s$  from  $H_0$  and BAO: CMB estimate must decrease by  $\sim 10 \text{ Mpc}$ 



$$r_s = \int_{\infty}^{z_*} dz \frac{c_s(z)}{8\pi G/3\sqrt{\rho_{\text{tot}}(z)}}$$

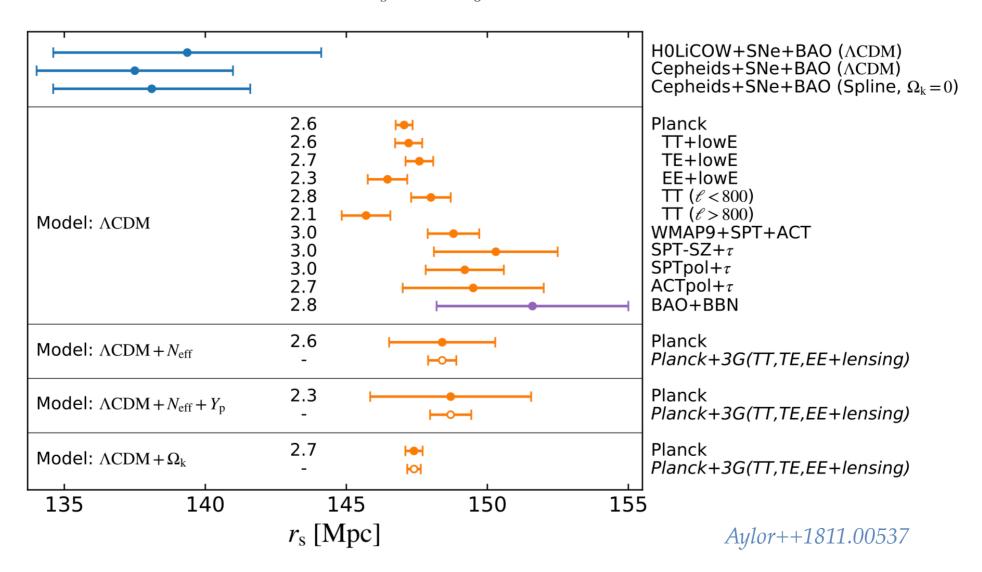
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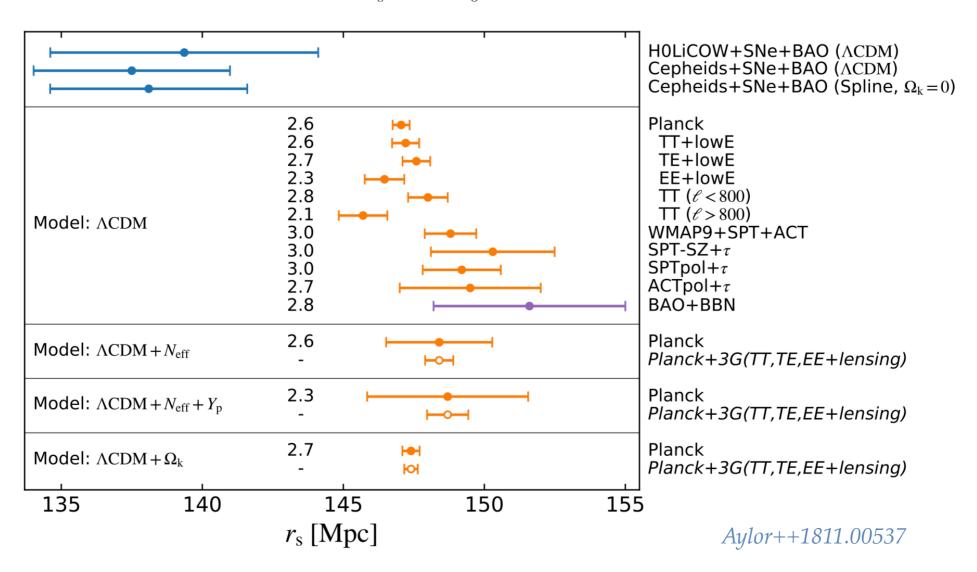


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$$r_s = \int_{-\infty}^{z_*} dz \frac{c_s(z)}{8\pi G/3\sqrt{\rho_{\text{tot}}(z)}}$$

increase  $\rho(z)$ : Neff? Early Dark Energy?

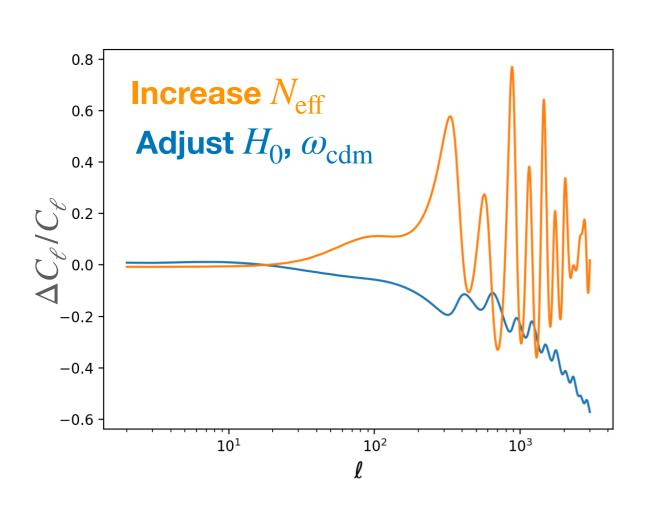
Modified Gravity?

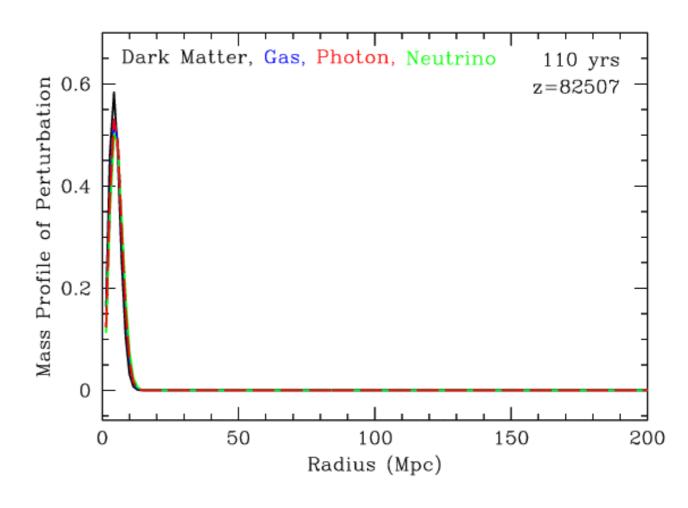
#### Relativistic species $N_{\text{eff}}$

ullet Additional relativistic degrees of freedom can be parametrized by  $N_{\rm eff}$ 

$$\rho_R = \rho_\gamma \left( 1 + \frac{7}{8} \left( \frac{4}{11} \right)^{4/3} N_{\text{eff}} \right)$$

Standard Model neutrinos behave as *free-streaming radiation* since T~1MeV with  $N_{\text{eff}} = 3.044$ 





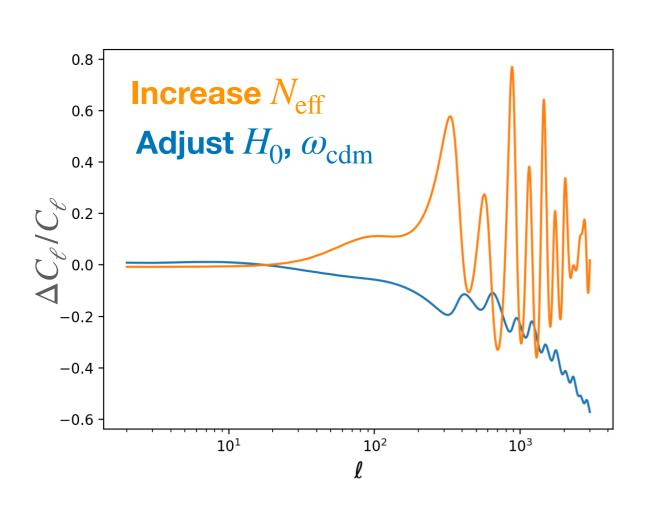
- Most of the "background" effects can be compensated for by changing  $\omega_{\rm cdm}$ ,  $H_0!$
- Standard  $\nu$ -perturbations phase-out and reduce  $\gamma$ -perturbations (" $\nu$ -drag").

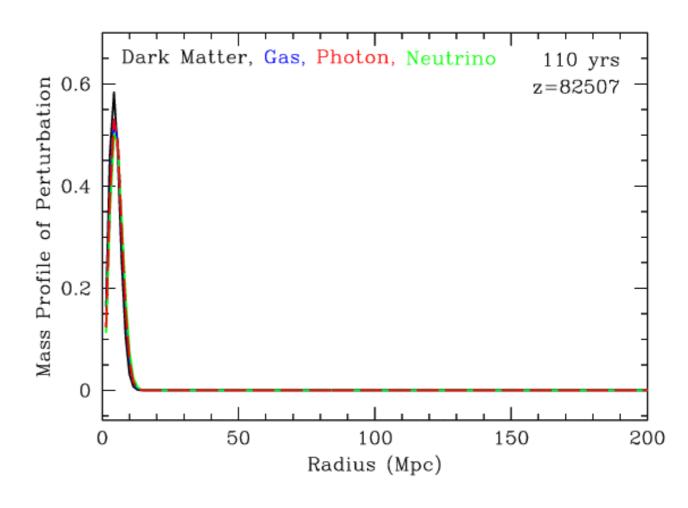
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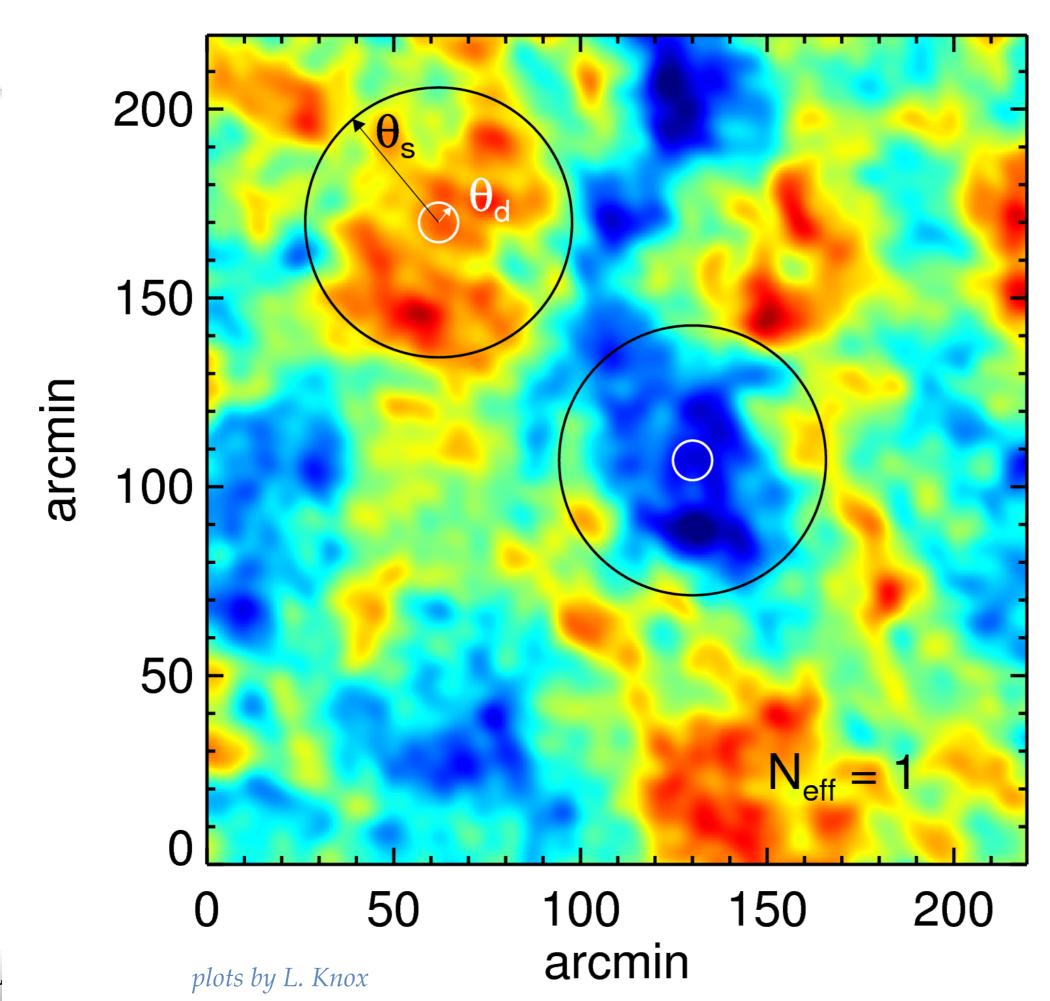
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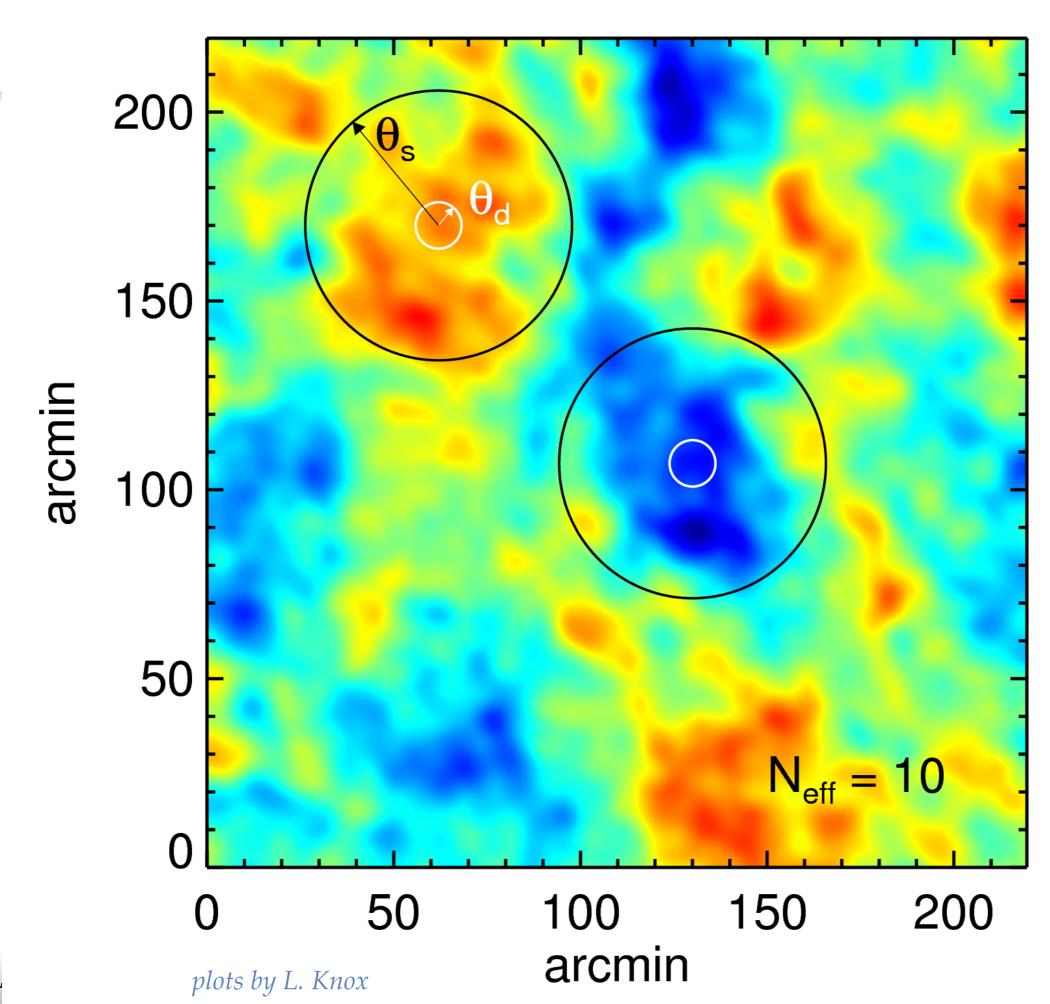
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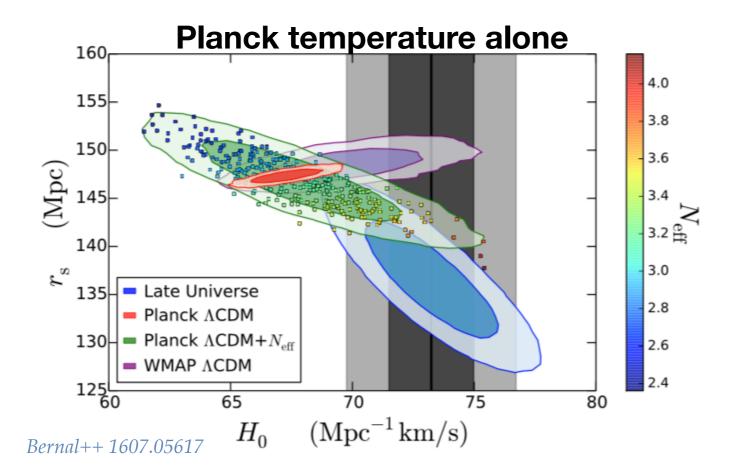
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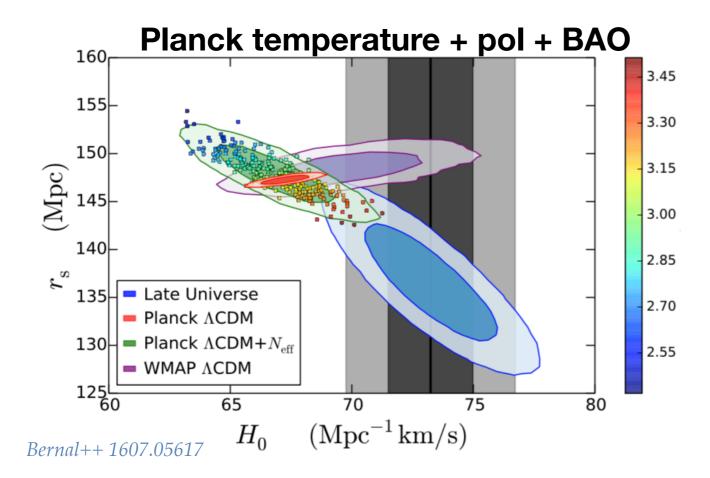




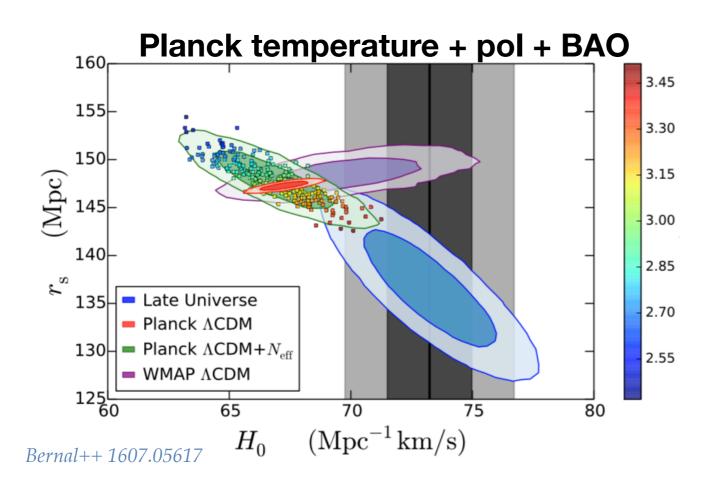
•  $\Delta N_{\rm eff}$  (free-streaming) ~0.5 – 1 is needed

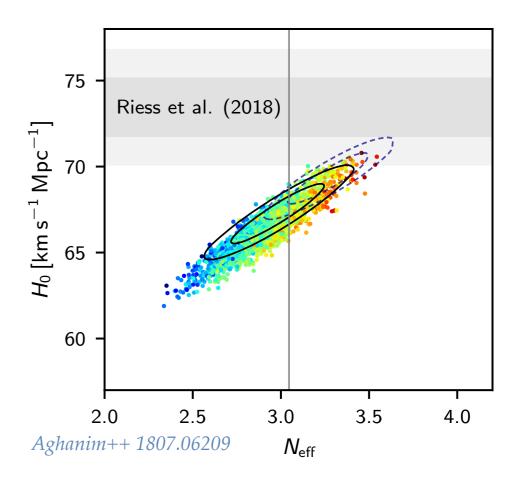


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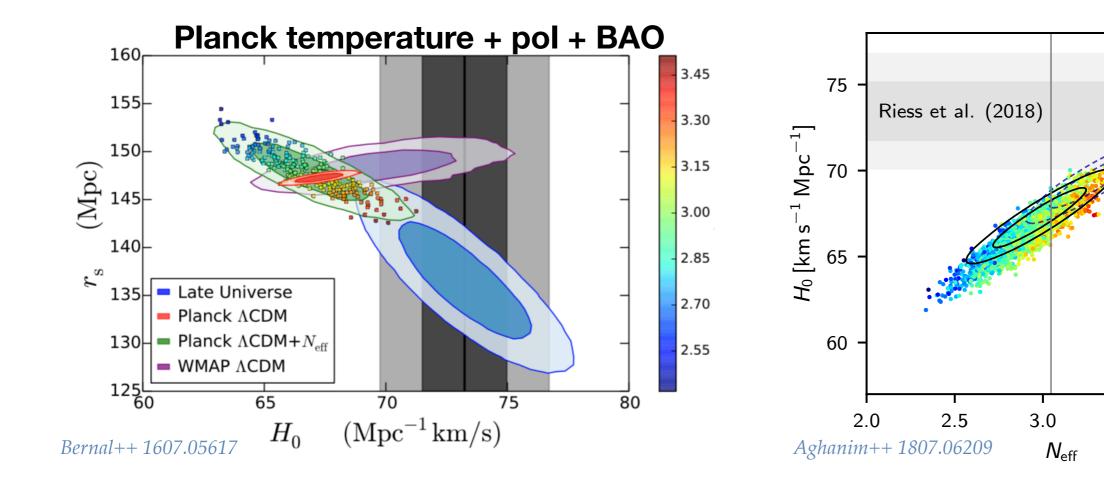


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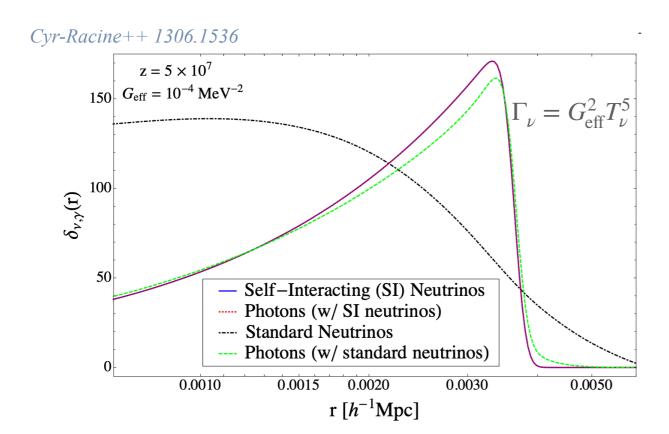
Planck+BAO constrains  $N_{\rm eff}=2.99\pm0.17$  and  $H_0=67.3\pm1.1$  km/s/Mpc Is it possible to relax this bound?

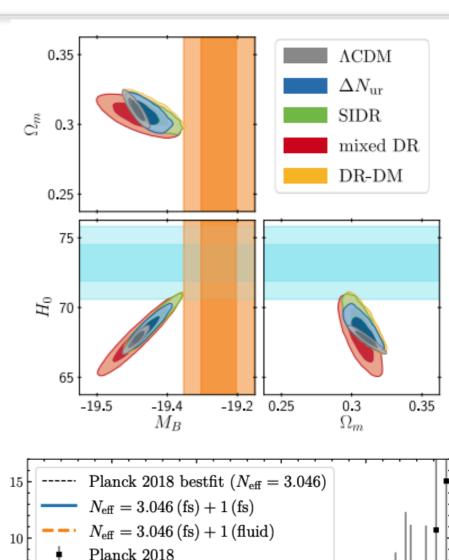
3.5

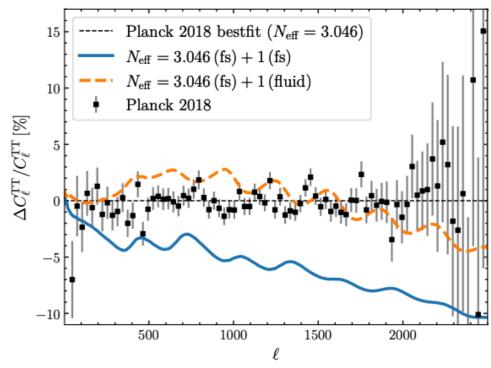
4.0

## Introducing exotic $\nu$ -interactions

- Exotic  $\nu$  self-interaction can erase free-streaming and make  $\nu$  behaves as perfect fluid.
- Once interactions are switched on:  $\gamma$ -perturbations are enhanced / phase is washed-out.
- This helps a bit but not enough to explain  $H_0$







#### Interacting neutrinos could resolve $H_0$ tension

see also Ghosh ++1908.09843

Free-streaming neutrinos lead to a phase shift:

\*\*Bashinsky&Seljak, astro-ph/0310198, Baumann++ 1508.06342\*\*

$$\theta_{\text{peak}} = \theta_s + \delta\theta \sim 0.6 \left(\frac{\rho_{\nu}}{\rho_g}\right)$$

• Self-interacting neutrinos: no phase shift,  $\theta_s$  is larger than in  $\Lambda$ CDM => larger  $H_0$ 

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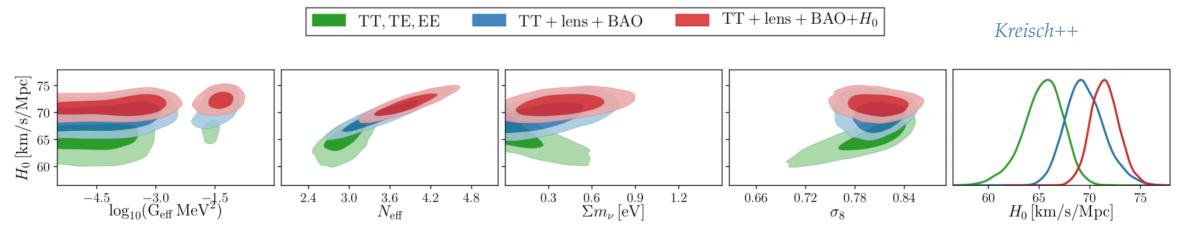
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- 2015 data: solution requires 4 strongly interacting neutrinos with  $M \sim 0.4 eV$
- "For free": solve  $S_8$  tension and reactor anomalies!
- BBN & Lab. requires majorana neutrinos and a heavy mediator coupled to  $\nu_{\tau}$

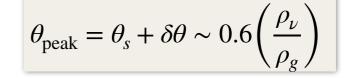
Blinov++ 1905.02727

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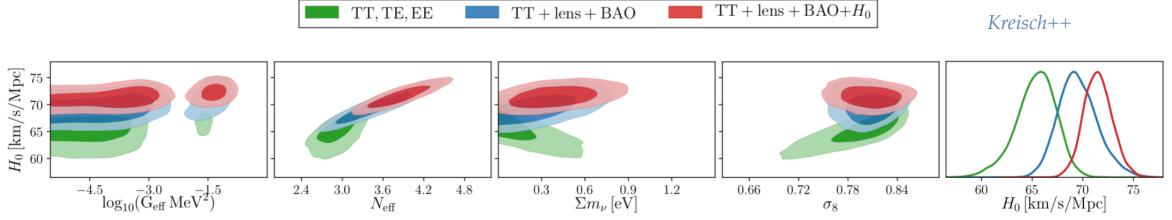
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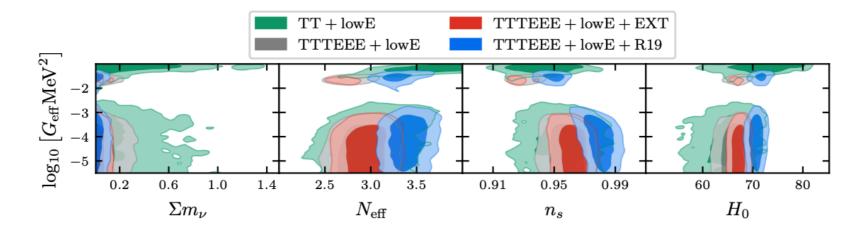
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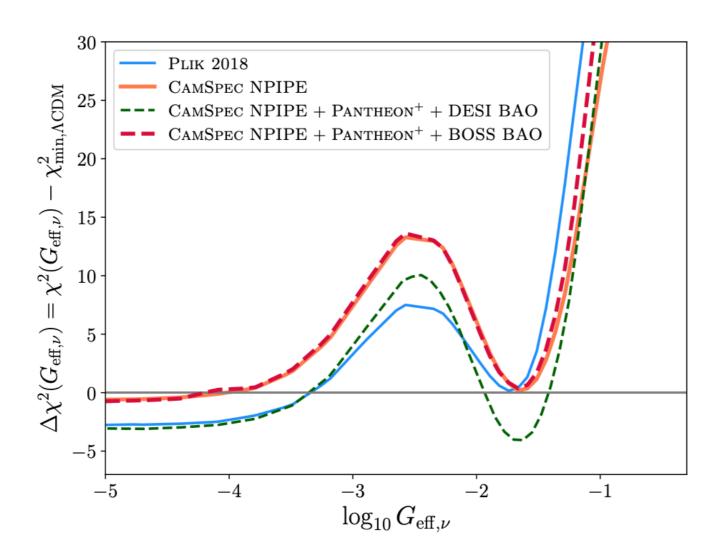
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Blinov++ 1905.02727

2018 Planck polarization data disfavor the strongly interacting mode.

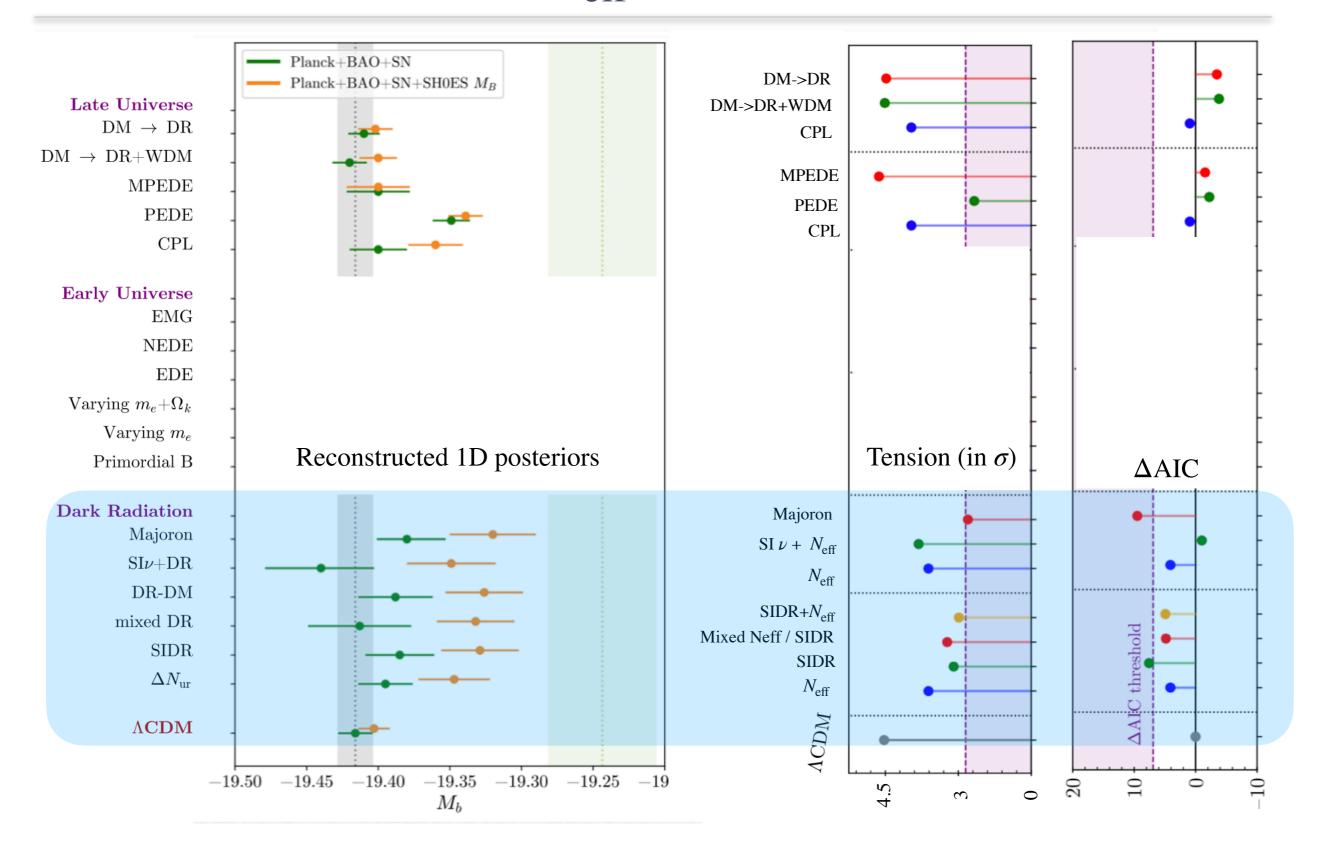


## Update with NPIPE & DESI

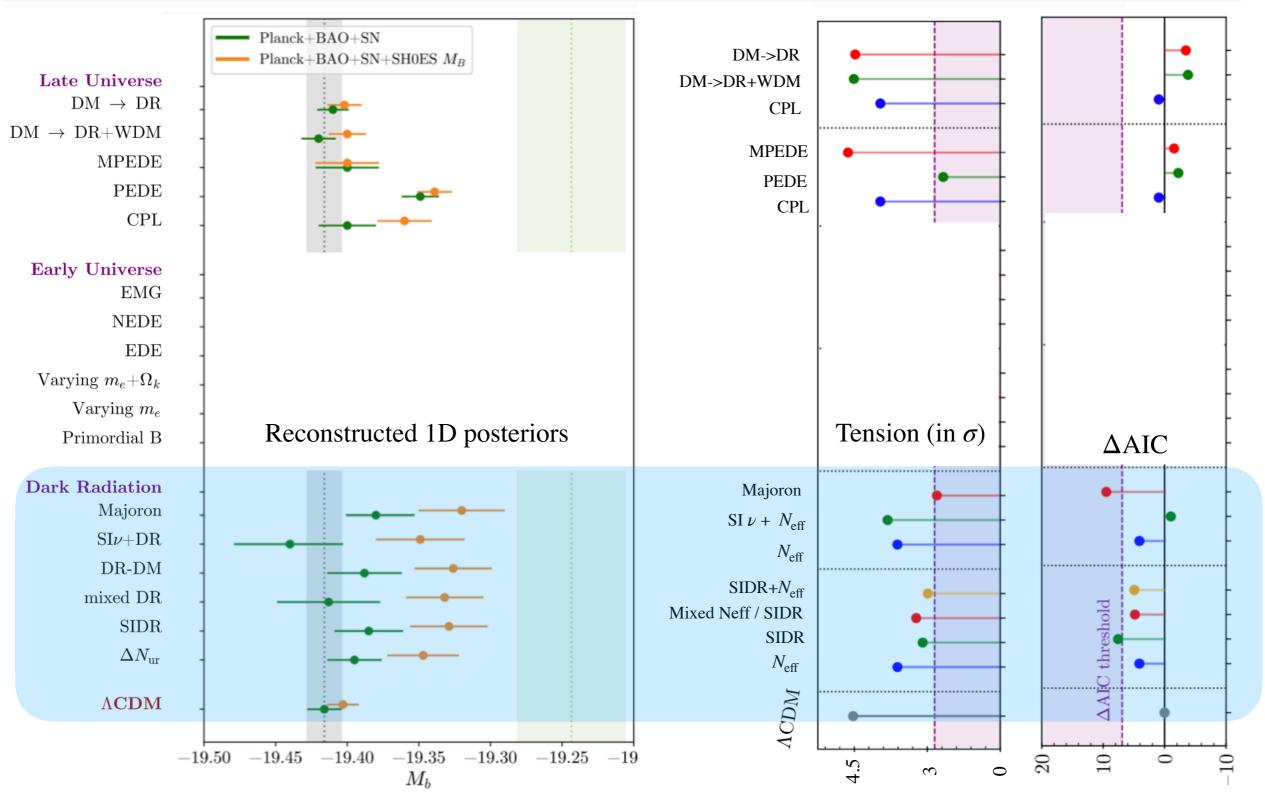


	$\mathrm{SI} u$	
	BOSS BAO	DESI BAO
$ m log_{10}G_{eff, u}$	$-1.63(-1.66)^{+0.13}_{-0.10}$	$-1.61(-1.66)^{+0.14}_{-0.11}$
$-2\log\mathcal{L}(\log_{10}G_{ ext{eff}, u})$	$-1.66^{+0.14}_{-0.13}$	$-1.66^{+0.11}_{-0.13}$
$H_0$	$67.58(67.2)_{-1.4}^{+0.91}$	$68.6(68.0)^{+1.2}_{-1.4}$

## Status of " $N_{\text{eff}}$ -like" solutions

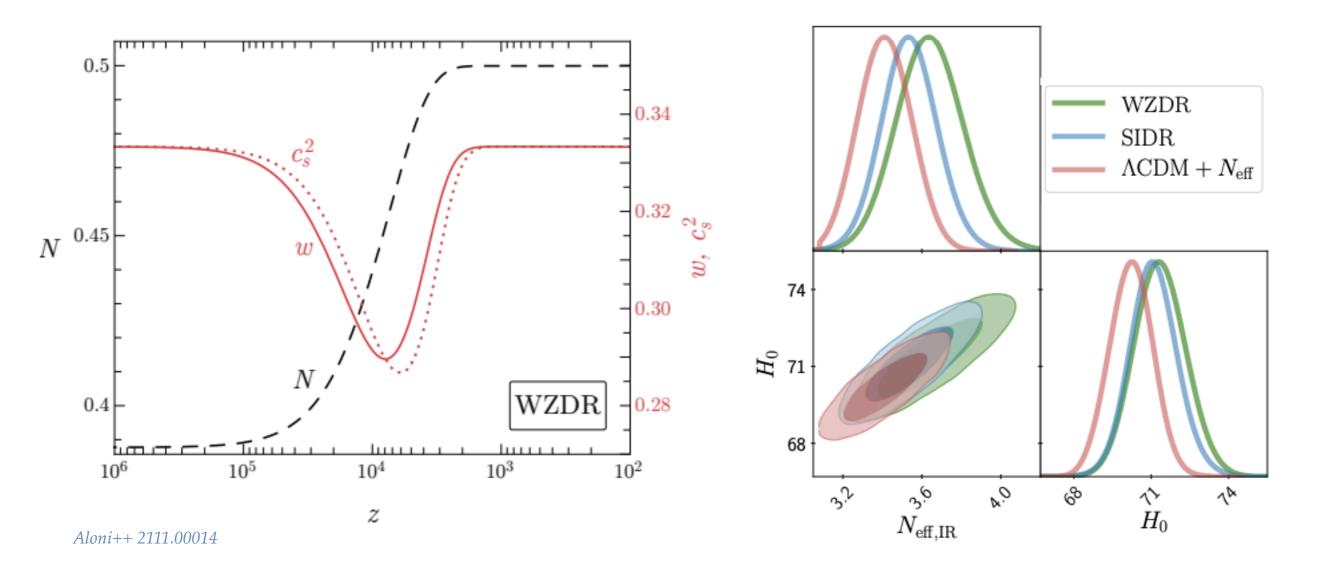


## Status of " $N_{\text{eff}}$ -like" solutions



Extra-radiation can bring the tension down to the  $3\sigma$  level but are not favored without  $H_0$  prior

## A step in resolving the H0 tension

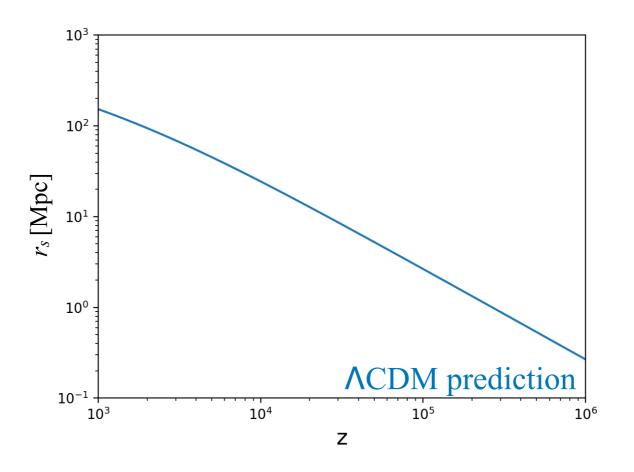


• If  $N_{\rm eff}$  has a step around  $z \sim 10^4$  e.g. from decoupling of a scalar  $\phi$  with  $m \sim 1 \, \rm eV$  reheating a bath of fermion  $\psi$ .

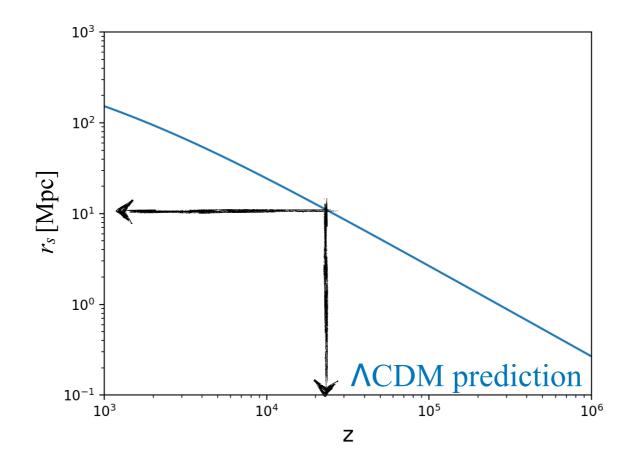
$$rac{N_{
m IR}}{N_{
m UV}} = \left(rac{g_*^\phi + g_*^\psi}{g_*^\psi}
ight)^{1/3} = \left(rac{15}{7}
ight)^{1/3} \simeq 1.29$$

ullet WZDR model provides a significantly better fit than standard  $N_{\rm eff}$  or SIDR

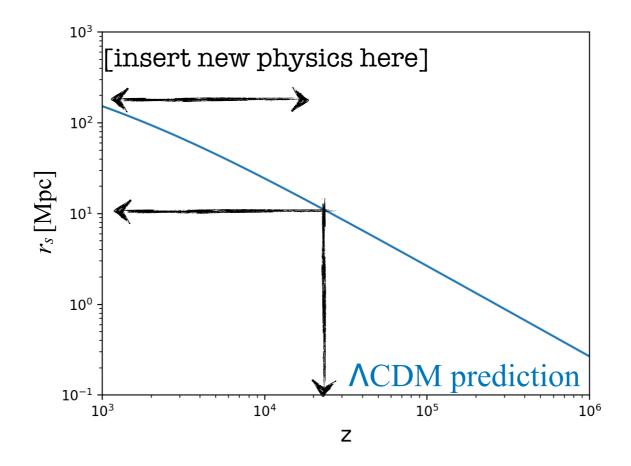
•  $r_s$  must decrease by  $\sim 10 \mathrm{Mpc}$ .



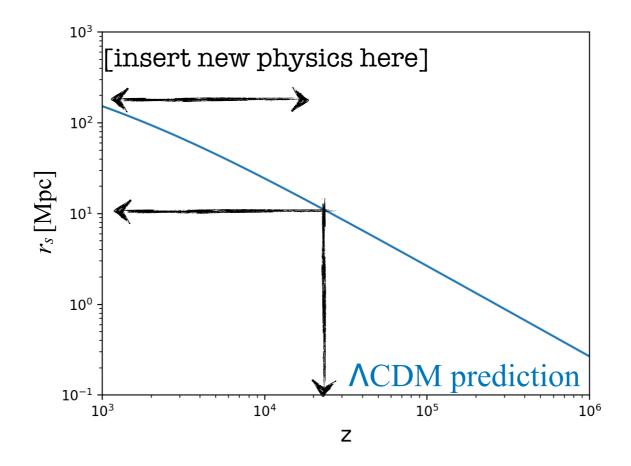
- $r_s$  must decrease by  $\sim 10 \mathrm{Mpc}$ .
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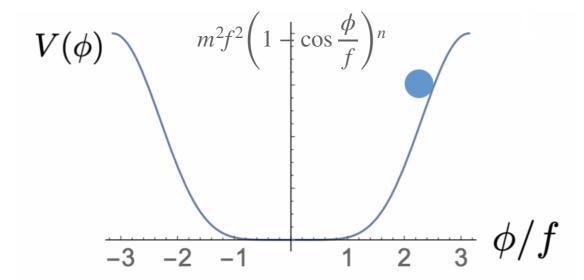


• Inject energy between recombination and  $z \sim 2 \times 10^4$  to reduce  $r_s$ 

## What is Early Dark Energy?

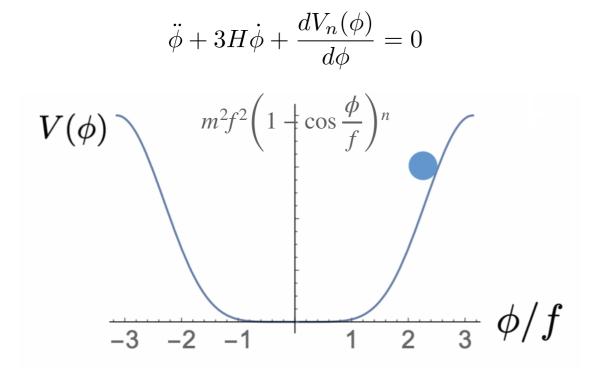
Initially slowly-rolling field (due to Hubble friction) that later dilutes faster than matter

$$\ddot{\phi} + 3H\dot{\phi} + \frac{dV_n(\phi)}{d\phi} = 0$$

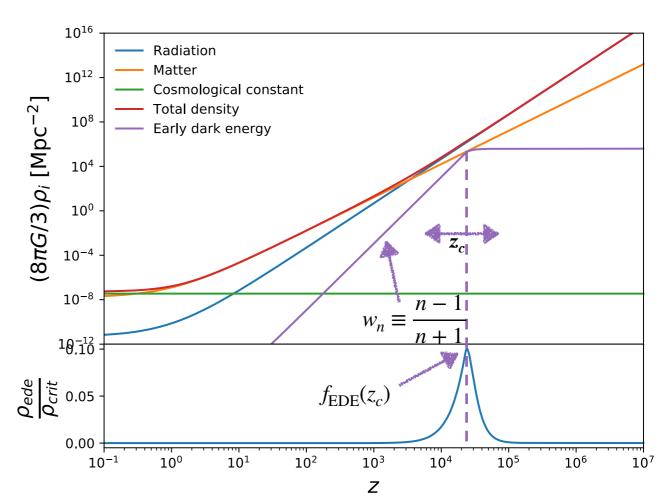


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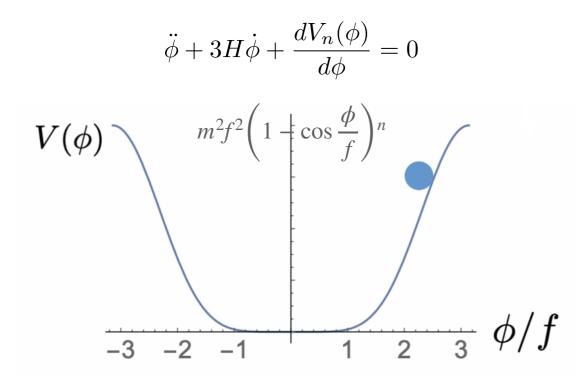


Specified by  $f_{\text{EDE}}(z_c)$ ,  $z_c$ , w(n),  $c_s^2(k, \tau)$ n = 1: matter, n = 2: radiation, etc.

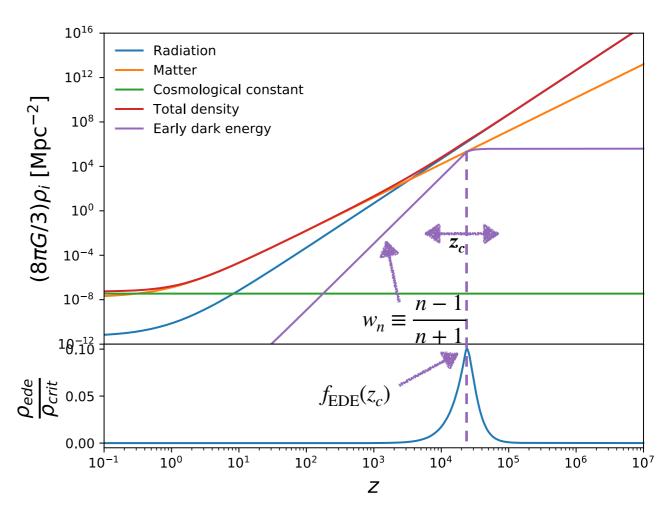


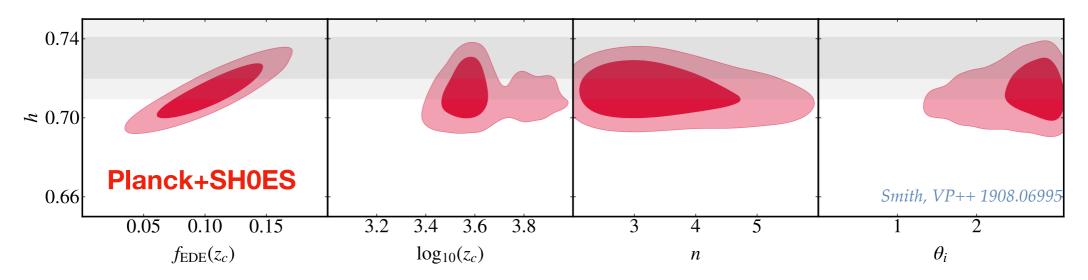
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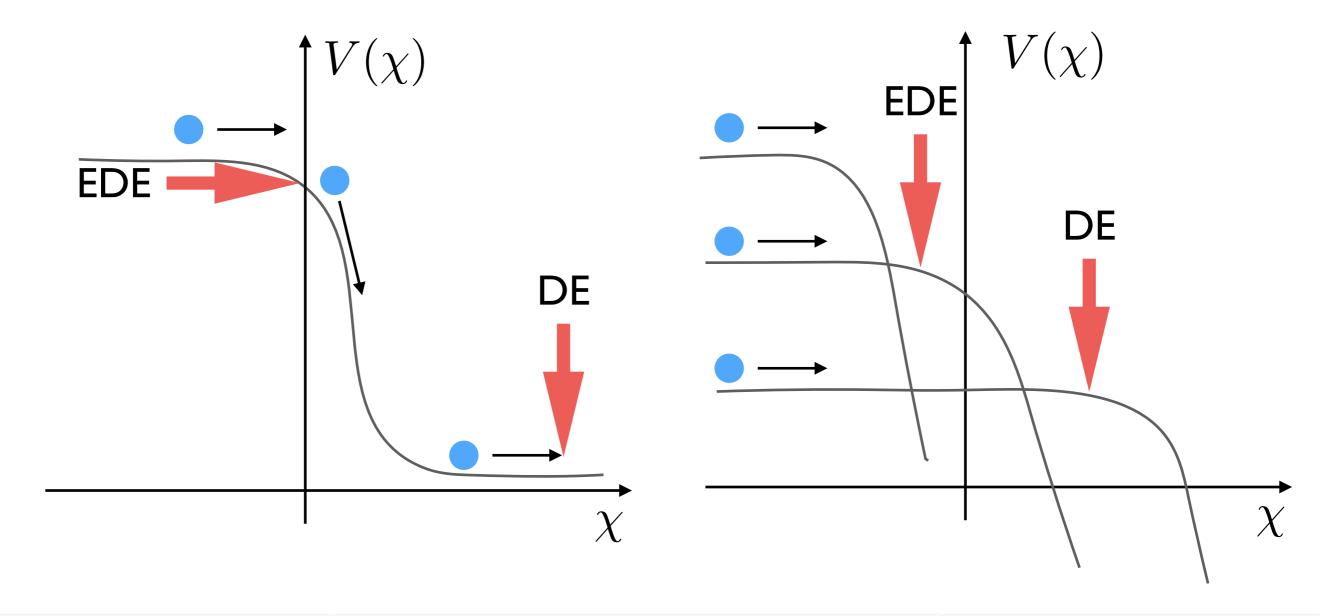


# A New Understanding Of Dark Energy?

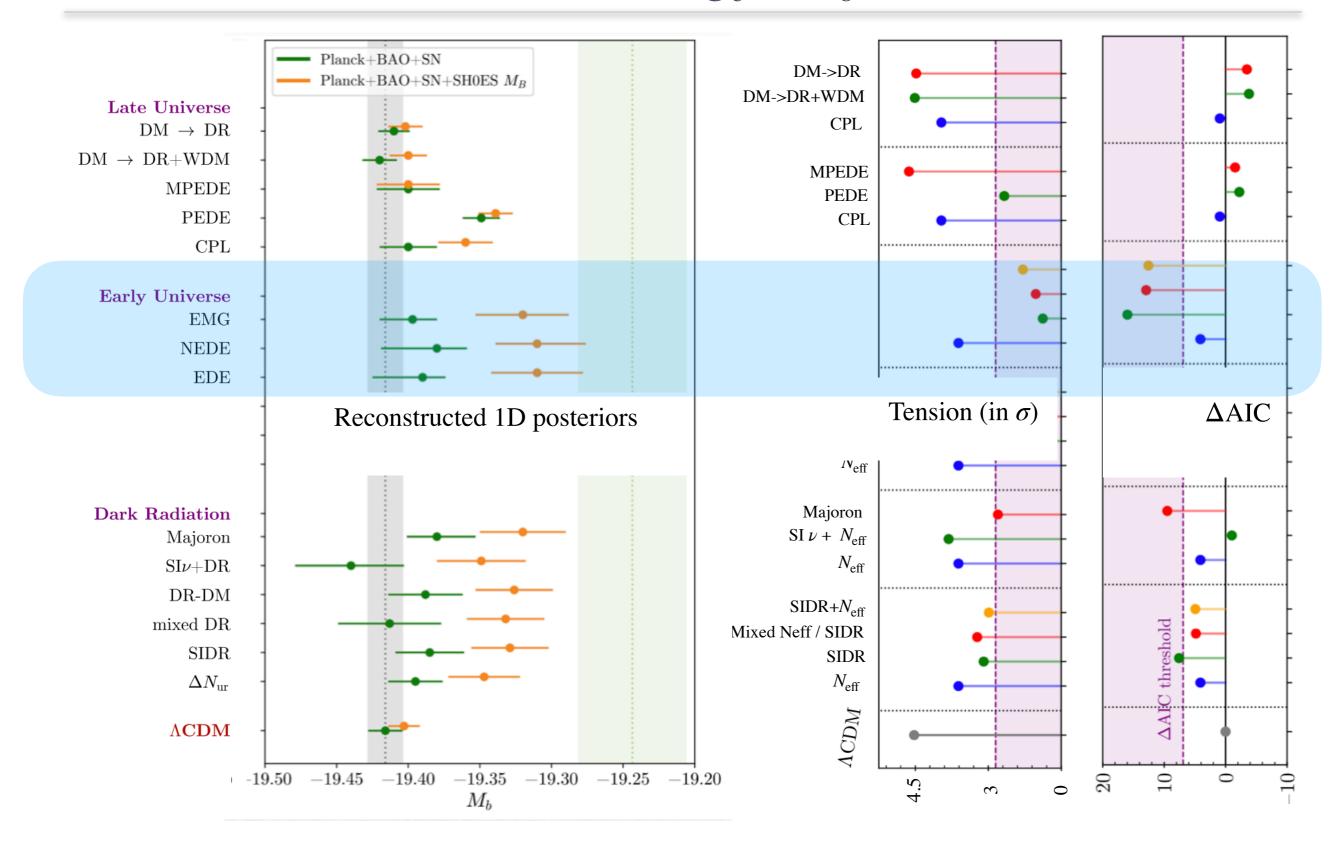
The idea that DE altered the early universe is not new: explains the DE coincidence problem

e.g. Dodelson++astro-ph/0002360, Griest astro-ph/0202052, Bielefeld (Caldwell)++ 1305.2209, Kamionkowski++1409.0549

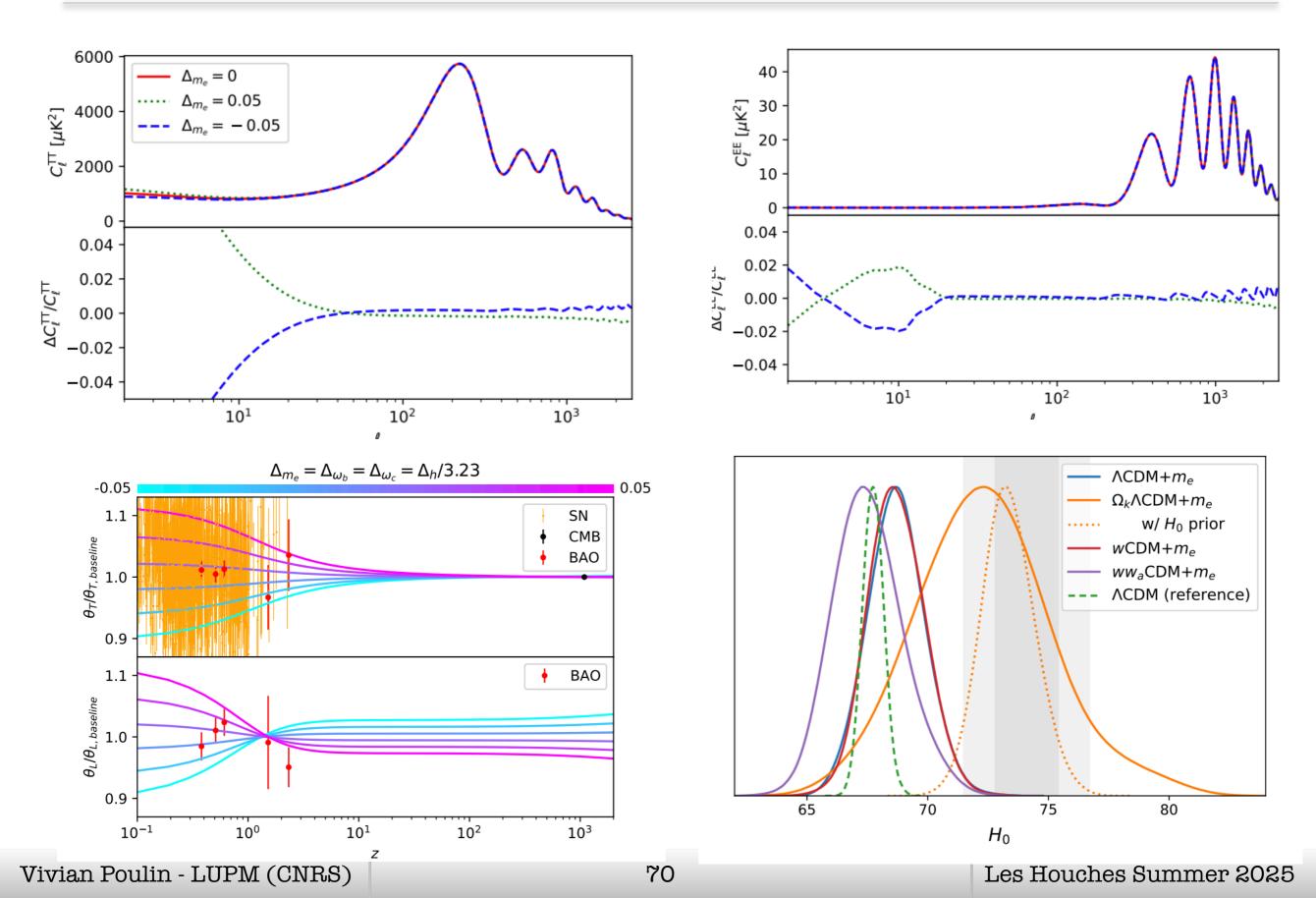
- Is their one field with a complicated potential or many fields with simple potentials?
- We know of Inflation, late-time Dark Energy: could there be more of such era to be discovered?



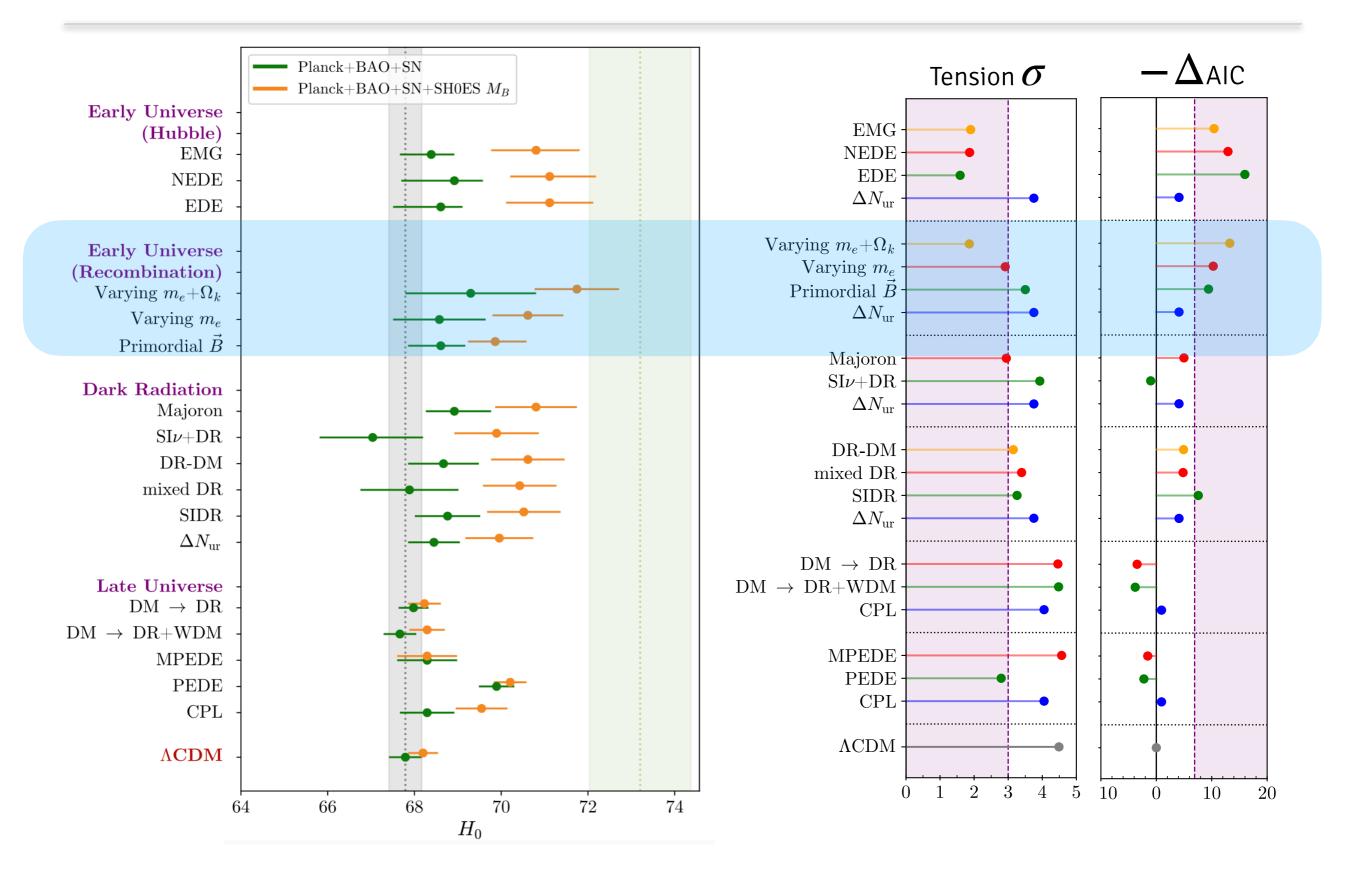
## Status of "Exotic Energy Injection" solutions



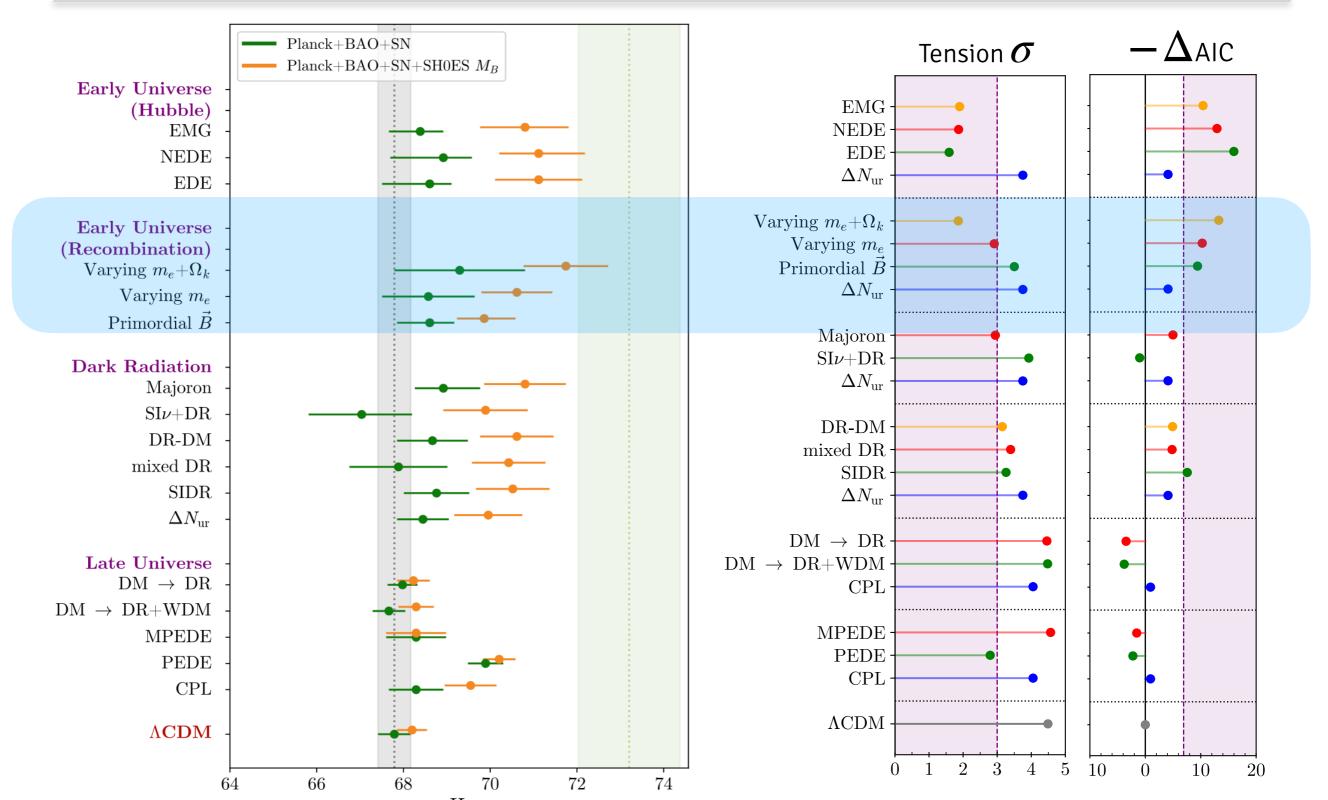
#### A higher $m_e$ could resolve the $H_0$ tension



#### Status of "modified recombination" solutions

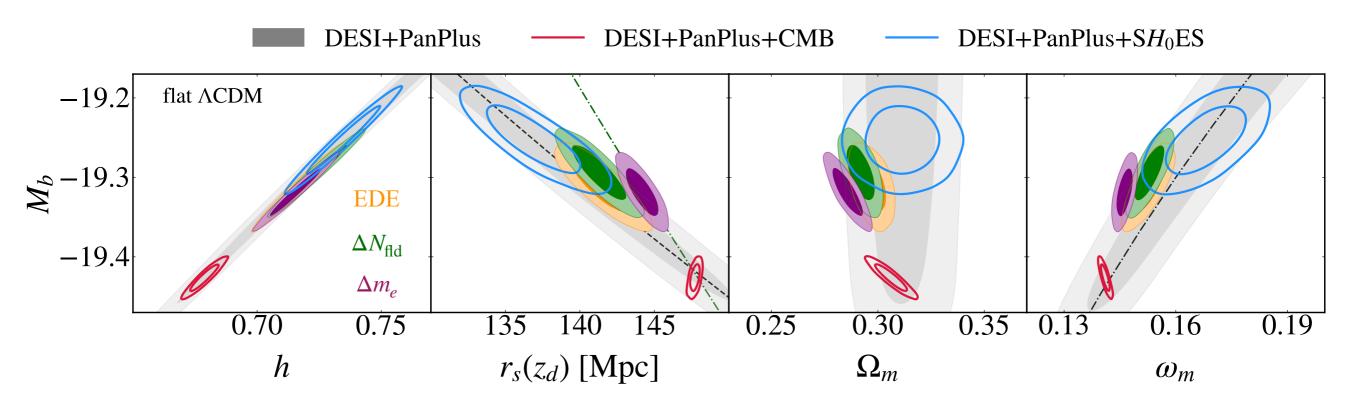


#### Status of "modified recombination" solutions

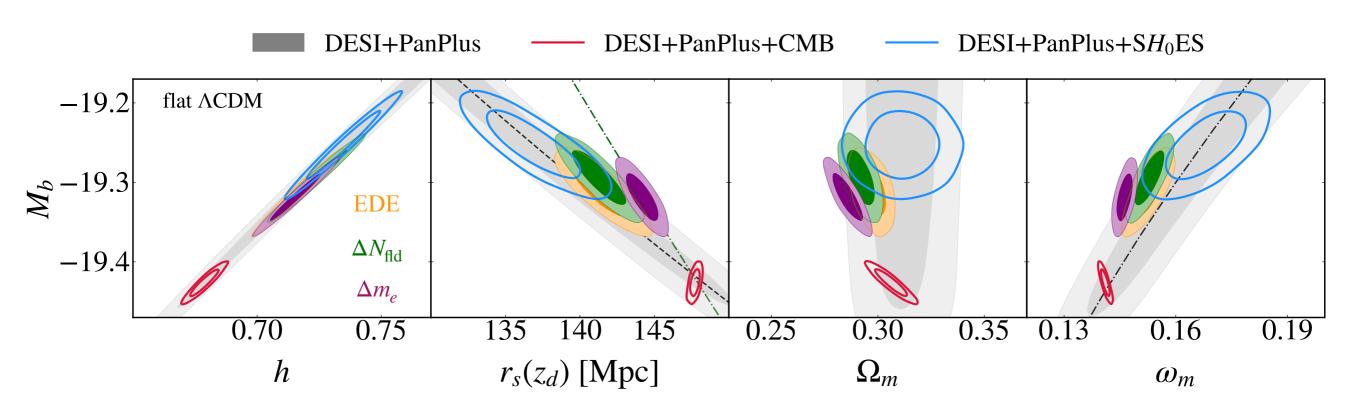


 $m_e + \Omega_k$  can bring the tension down to the 1.8 $\sigma$  level and is slightly favored without  $H_0$  prior

#### Can models respect degeneracies imposed by the distance ladder?

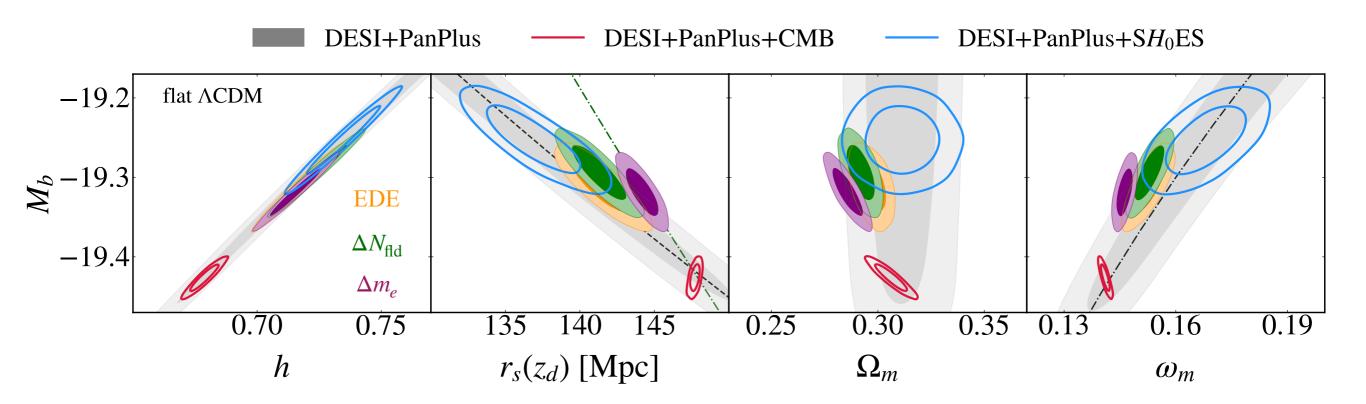


#### Can models respect degeneracies imposed by the distance ladder?



• Models affecting expansion history can reduce tension to  $\sim 2-3\sigma$  level

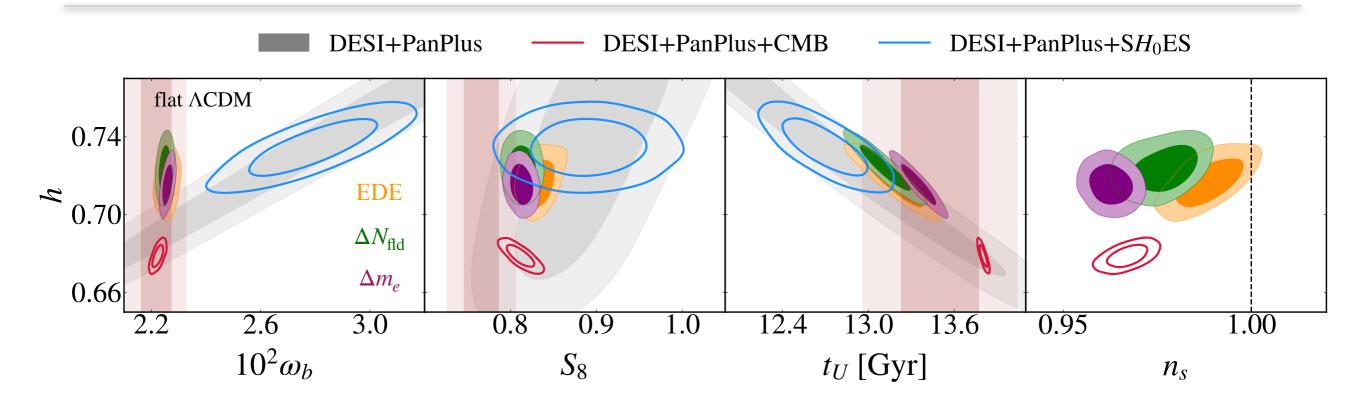
#### Can models respect degeneracies imposed by the distance ladder?



- Models affecting expansion history can reduce tension to  $\sim 2-3\sigma$  level
- Models affecting solely the way recombination proceeds are disfavored: they lead to a low  $\Omega_m$

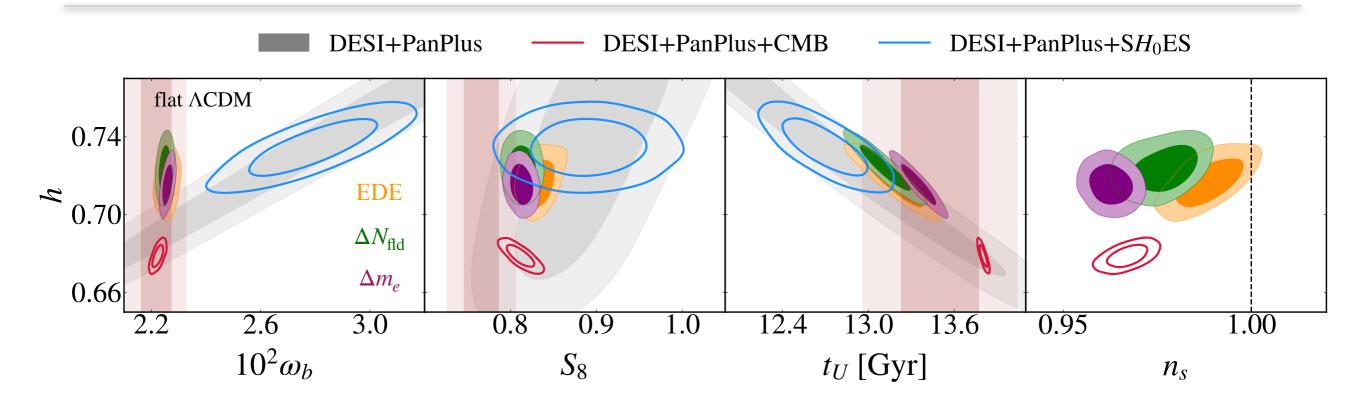
*Lee (VP)++ PRL 2022, Lynch++ 2404.05715* 

#### Can models resolve the other tensions?



- No more tension with BBN but tension with weak lensing measurements at the  $3 3.5\sigma$  level
- Age of the universe  $\sim 0.7$  Gyr younger: problem with old objects? JWST?

#### Can models resolve the other tensions?

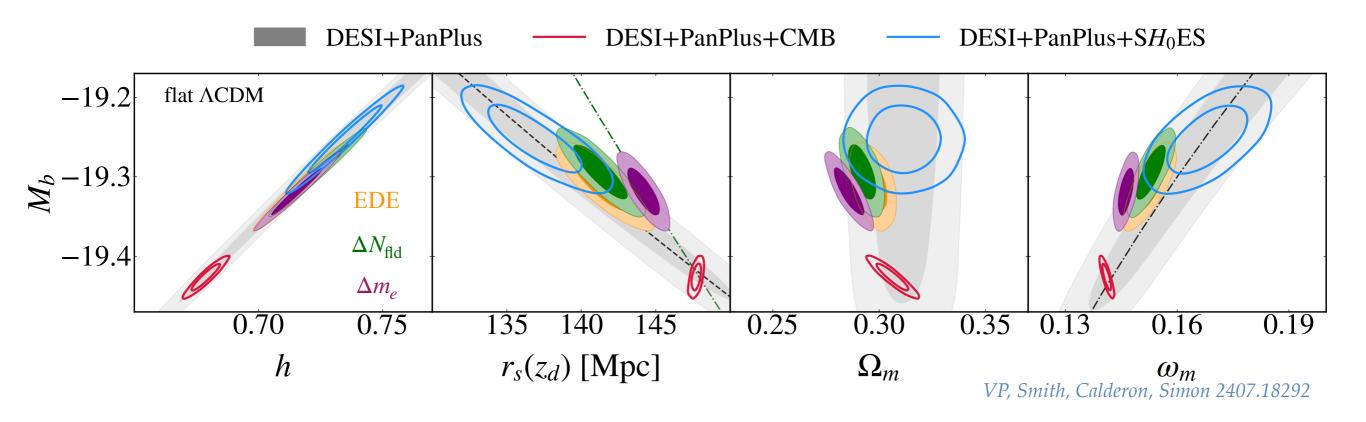


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- Age of the universe  $\sim 0.7$  Gyr younger: problem with old objects? JWST?
- $n_s$  is degenerate with the angular damping scale  $\frac{\delta\theta_D}{\theta_D} \simeq 0.2 \frac{\delta n_s}{n_s}$ . Smith&VP, 2309.03265
- $\theta_D \sim \sqrt{H_0}$  while  $\theta_s \propto H_0$ :  $n_s$  must increase to compensate a higher  $H_0$ .

- Constraints to  $\Omega_m$  are an important part of the problem: how to relax them?
- Hints of dynamical dark energy: could it play a role in the tensions?

DESI 2404.03002, Cortês&Liddle 2404.08056, Shlivko&Steinhardt 2405.03933, Berghaus++ 2404.14341, Efstathiou 2408.07175

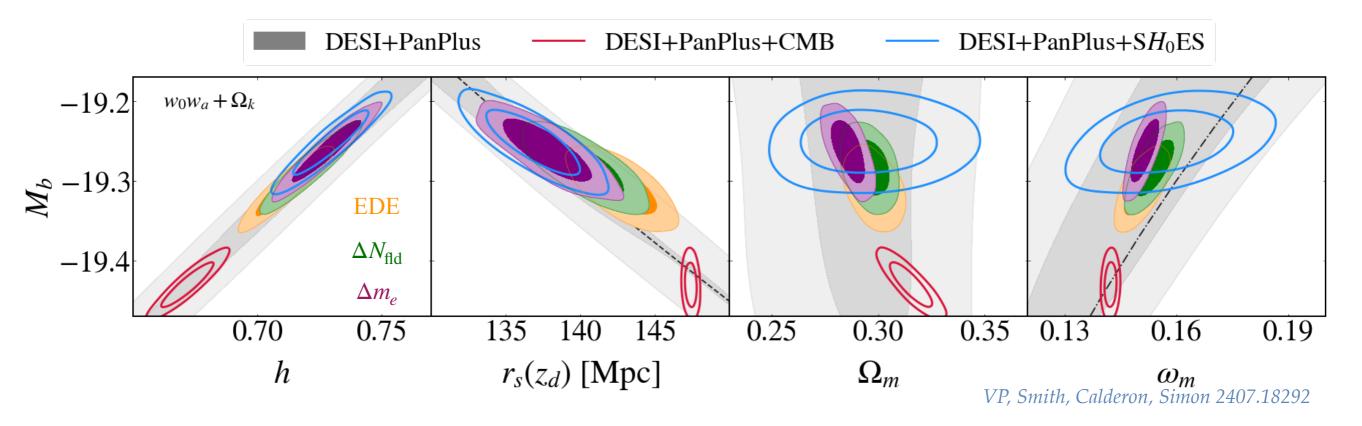
• Test: Open the late-time expansion history  $(w_0, w_a, \Omega_k)$ 



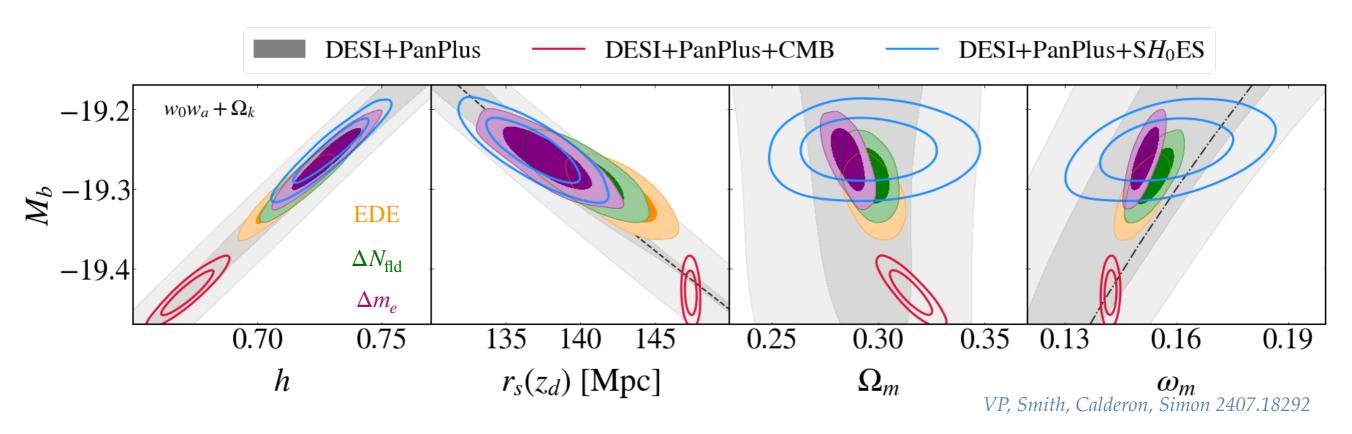
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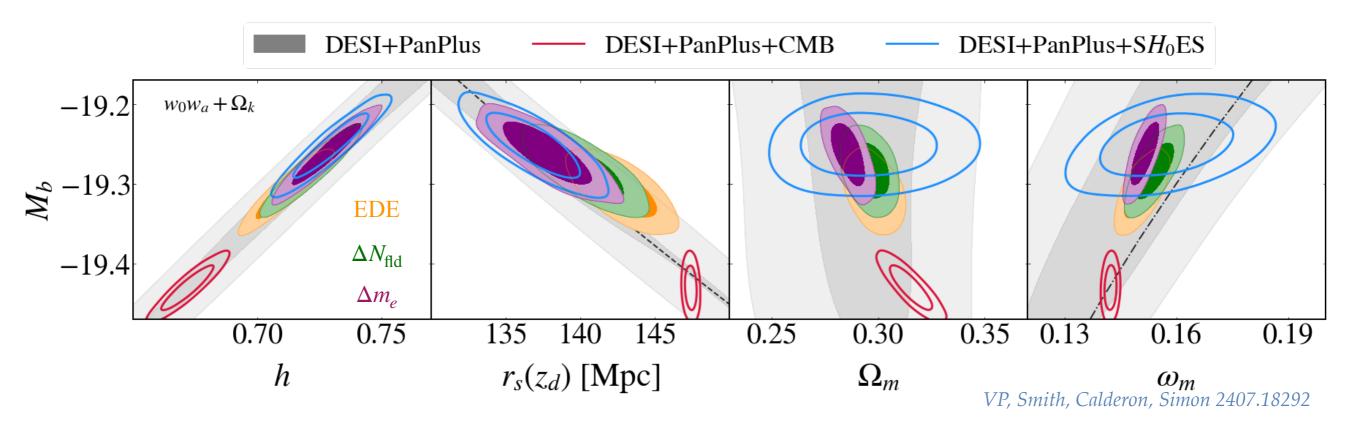
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Sekiguchi++2007.03381

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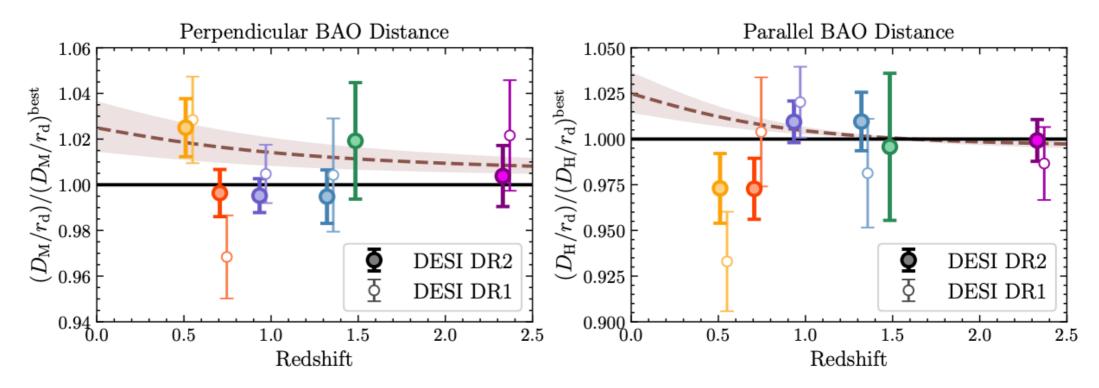


- $\Delta m_e$  model has almost no tension with  $\Omega_k \simeq -0.0114 \pm 0.0031$ ; No impact for EDE and  $\Delta N_{\rm fld}$ .

  Sekiguchi++2007.03381
- Localized energy injection boosting H(z) or a broad change affecting several cosmic epochs?

# DESI BAO measurements: hint of new physics?

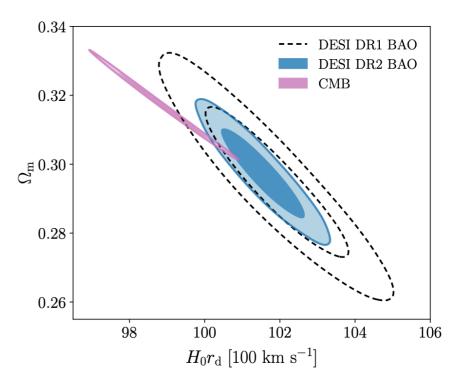
• Under  $\Lambda$ CDM, 2.3 $\sigma$  tension between CMB and BAO data



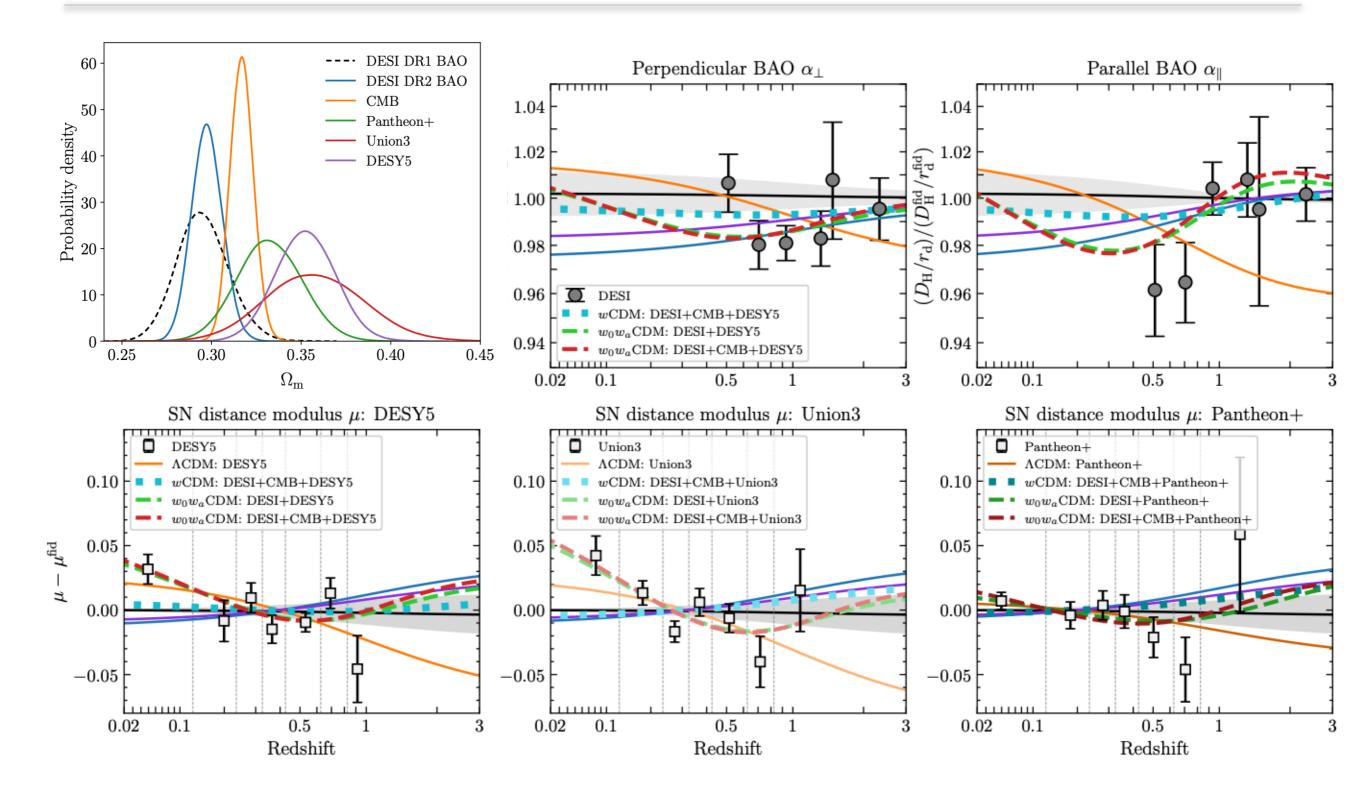
• Under  $\Lambda$ CDM, the BAO allows to measure  $\Omega_m$  and  $H_0r_d$ .

$$\frac{r_d}{D_M} \equiv \frac{H_0 r_s(z_d)}{\int_0^z dz (\Omega_m [(1+z)^3 - 1] + 1)^{-1/2}}$$

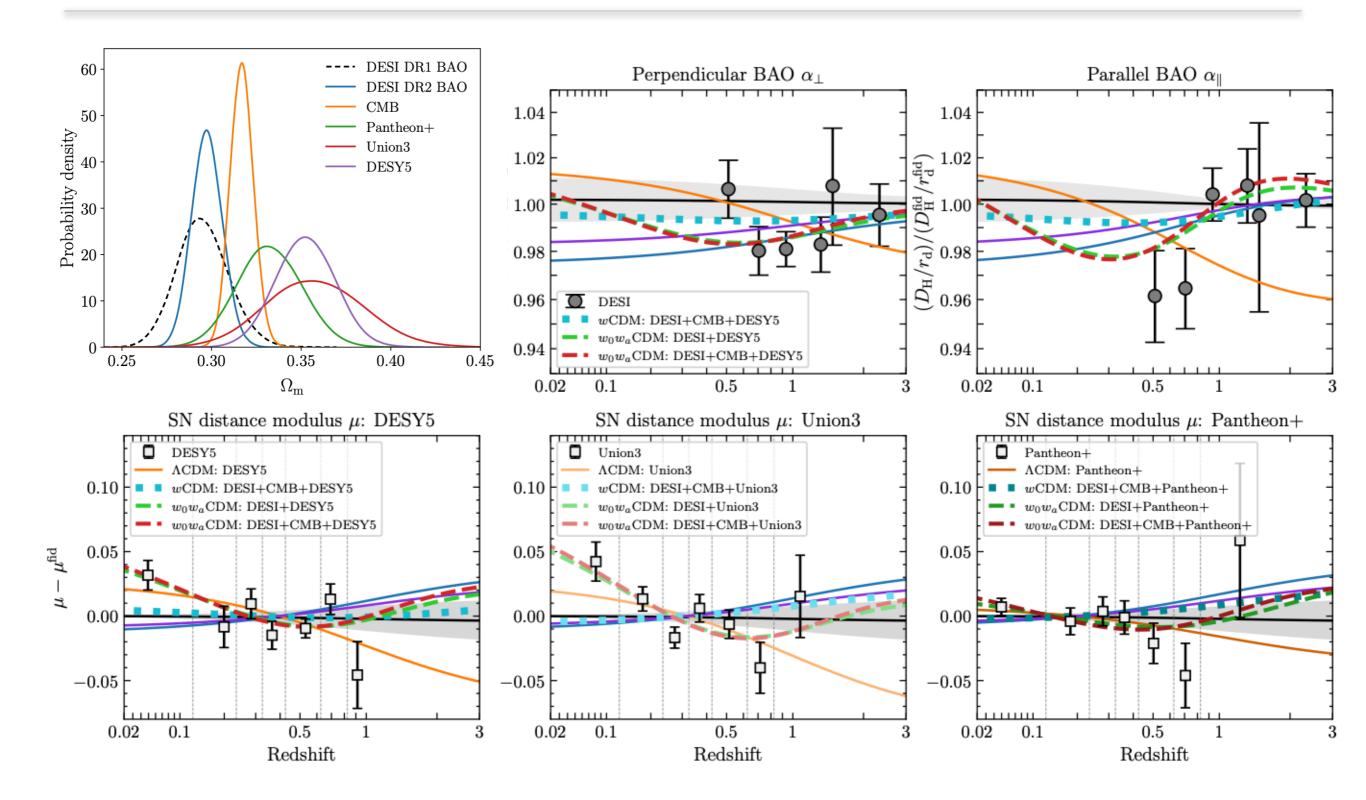
$$\frac{r_d}{D_H} \equiv H_0 r_s(z_d) \sqrt{\Omega_m [(1+z)^3 - 1] + 1}$$



#### DESI vs SNIa

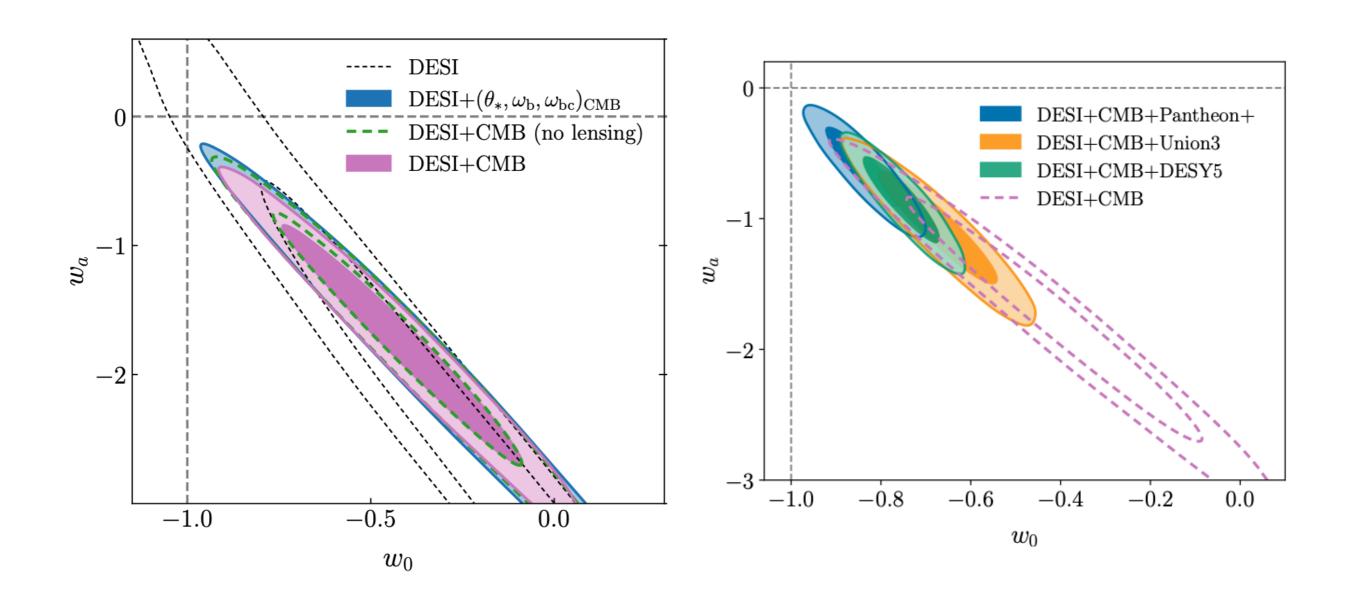


#### DESI vs SNIa

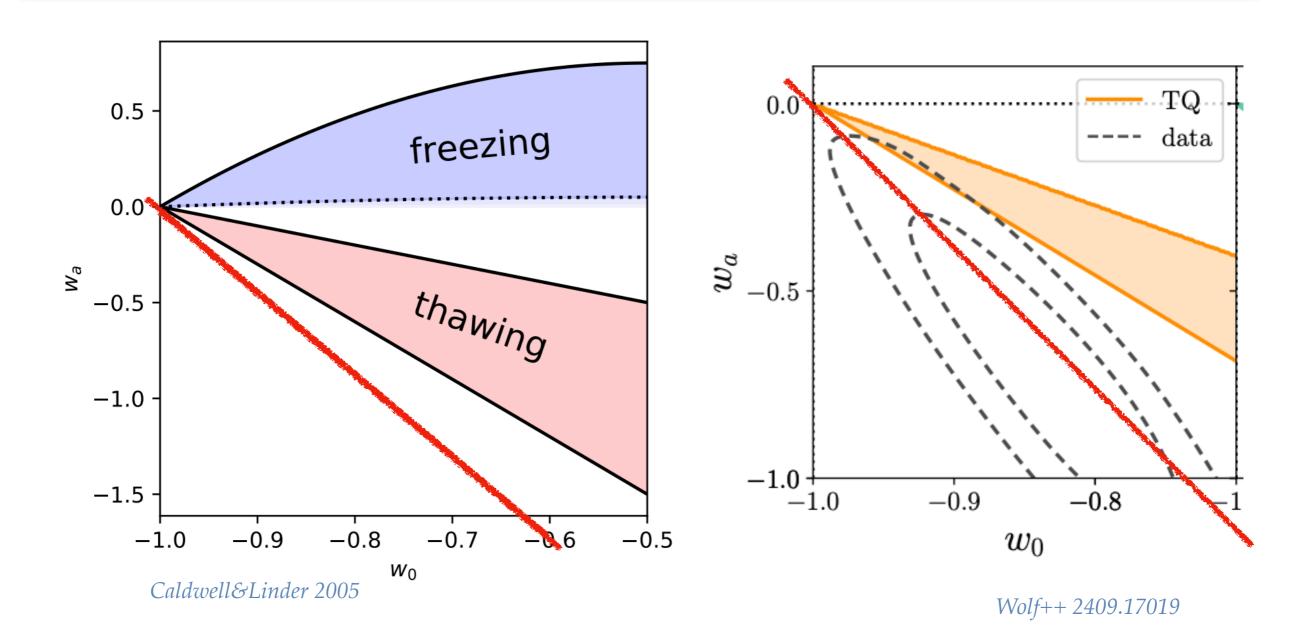


• DESI+CMB in tension at the  $\sim 2-3\sigma$  level with SN1a in the determination of  $\Omega_m$ 

# Evidence for dynamical dark energy



### Evidence for non canonical quintessence



$$\langle w \rangle = -1 \Rightarrow w_a \approx -3.66(1 + w_0)$$

Linder 0708.0024

# A mirage of dynamical dark energy?

DESI+CMB:	DESI+CMB: +PantheonPlus		+Union3		+DESY5	
DE classes		$\Delta { m DIC}$	$(\Delta\chi^2)$			
Thaw. (Cal.)	+0.4 (-1.6)	-0.6	(-2.5)	-5.8	(-7.1)	
Thaw. (Alg.)	-1.0 (-2.9)	-4.6	(-6.9)	-10.1	(-13.2)	
Emergent	+2.1 (-0.05)	+1.8	(-0.1)	+0.2	(-1.5)	
Mirage	-9.1 (-10.5)	-13.8	(-16.2)	-18.7	(-20.7)	
$w_0w_a$	-6.8 (-10.7)	-13.5	(-17.4)	-17.2	(-21.0)	

$$\langle w \rangle = -1 \Rightarrow w_a \approx -3.66(1 + w_0)$$

Linder 0708.0024

WEIRD??

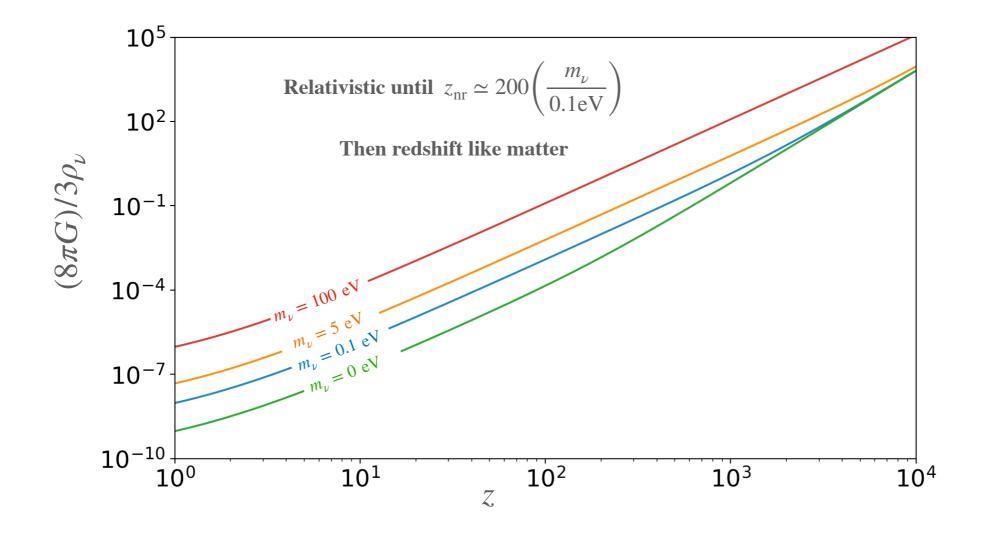
# $m_{\nu}$ can affect cosmological distances

•  $\nu$  can affect the homogeneous expansion

$$H(z) \equiv \frac{\dot{a}}{a} = H_0 \sqrt{\Omega_m (1+z)^3 + \Omega_r (1+z)^4 + \Omega_\Lambda + \dots}$$

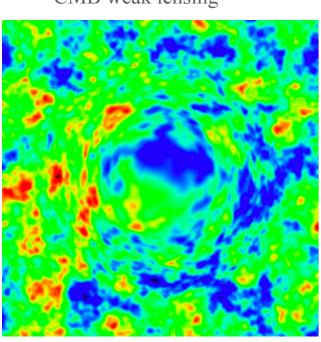
• Affect distance measurements via SN1a and BAO

$$D_M(z) = \int_0^z \frac{cdz}{H(z)}$$

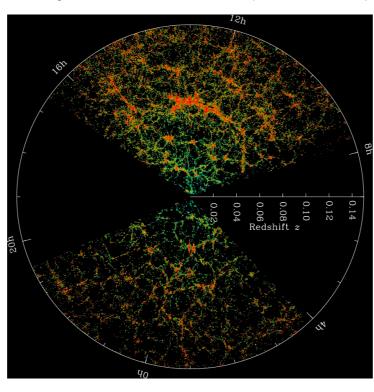


# $m_{\nu}$ can affect probes of matter structuring

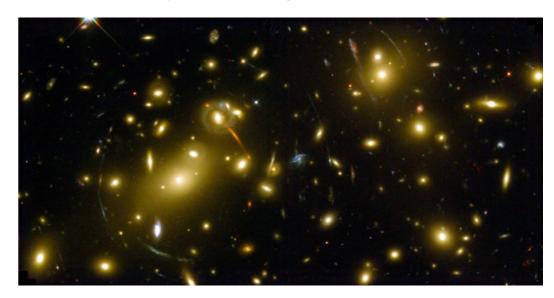
CMB weak lensing



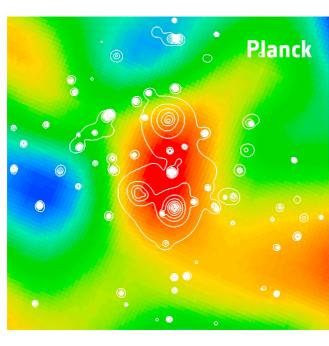
Galaxy correlation function (SDSS/BOSS)

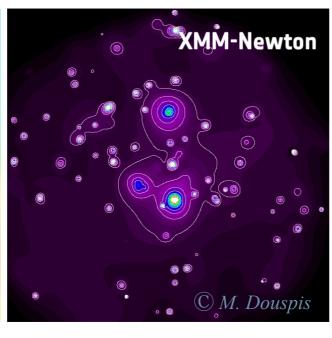


Galaxy weak lensing (KiDS / DES)

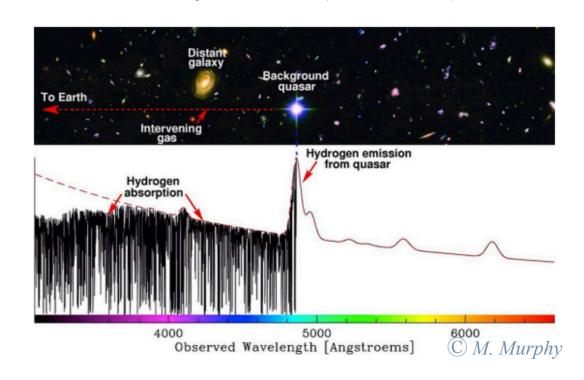


Galaxy cluster mass function

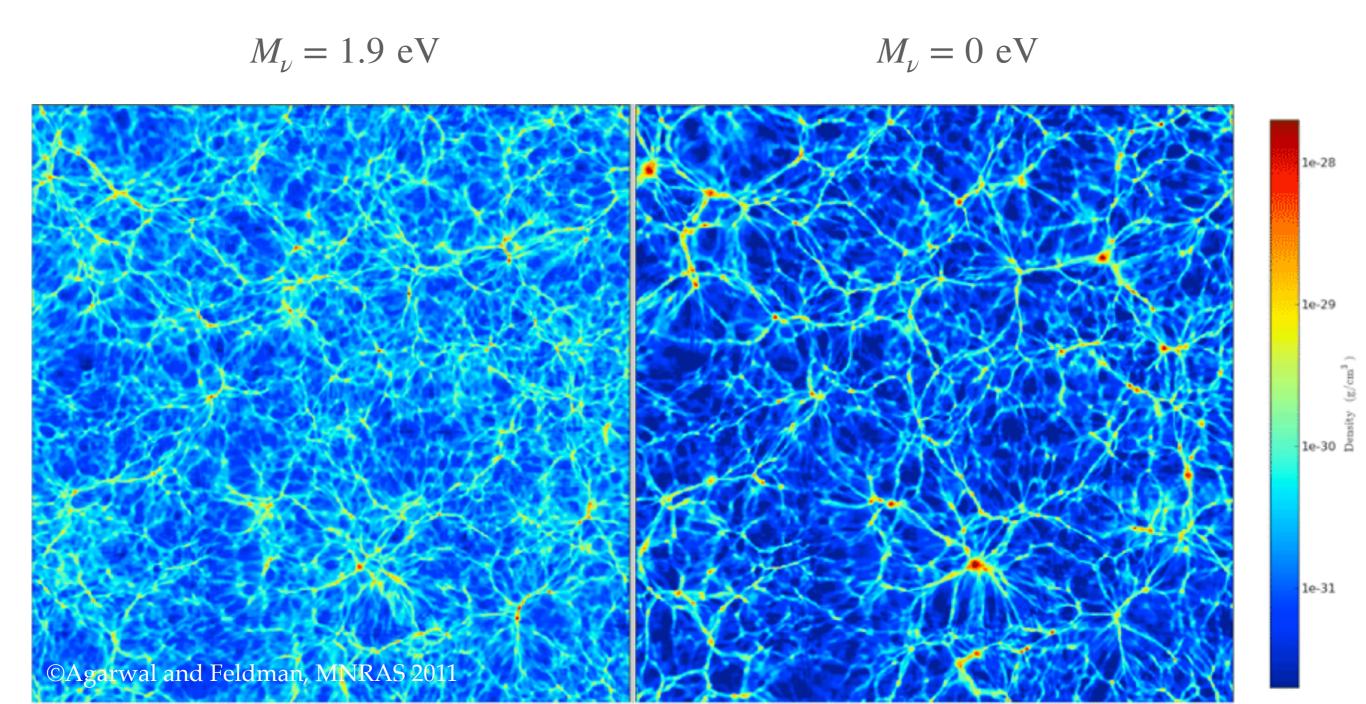




The Lyman- $\alpha$  forest (MIKE/HIRES)



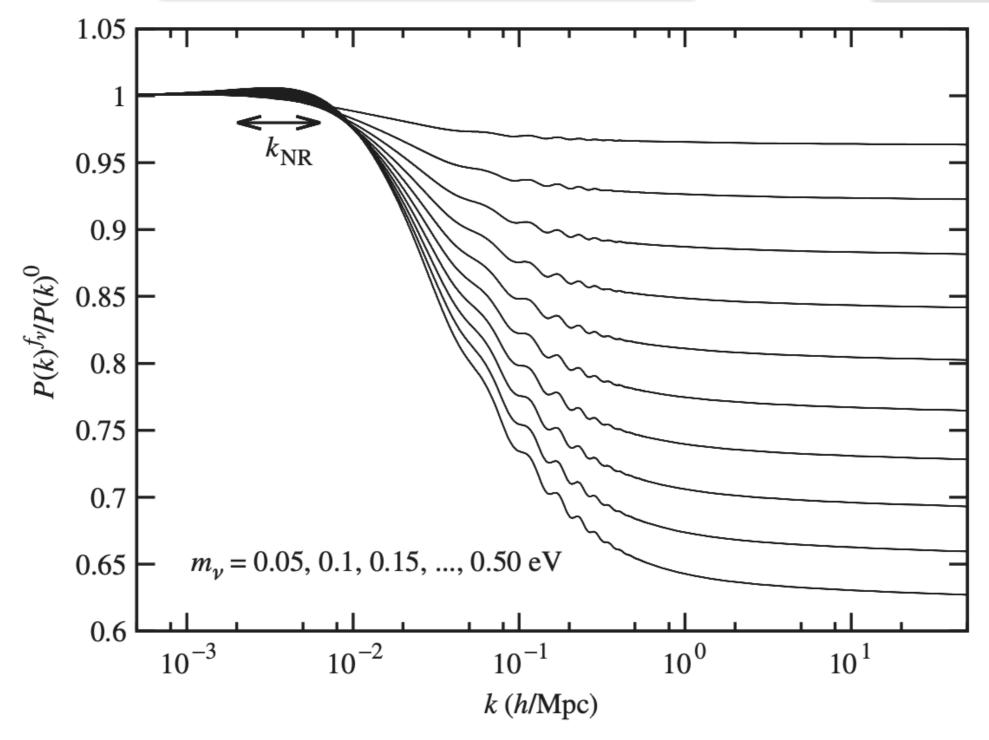
# $m_{\nu}$ suppresses the clustering of matter



The density field is "smoother" on small scales

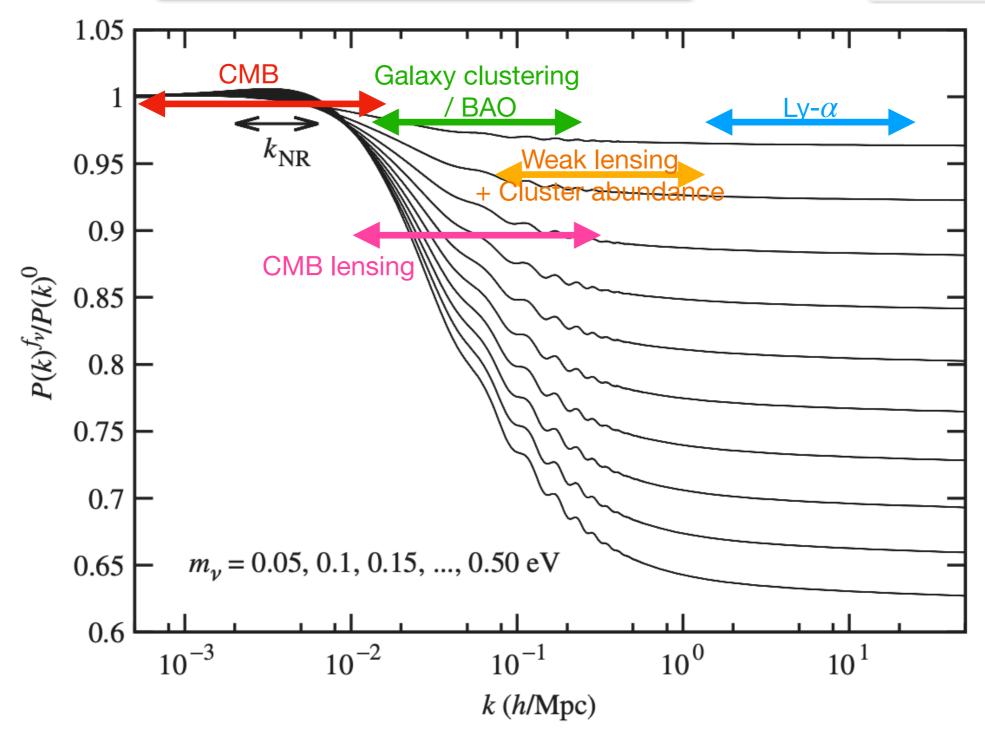
### $m_{\nu}$ suppresses the matter power spectrum

For 
$$\nu$$
's suppression at  $k \ge k_{\rm nr} \equiv 0.01 \left(\frac{m_{\nu}}{1 {\rm eV}}\right)^{1/2} \left(\frac{\Omega_{\rm m}}{0.3}\right)^{1/2} h {\rm Mpc}^{-1}$  and amplitude  $\frac{\Delta P}{P} \simeq -8 \frac{\omega_{\nu}}{\omega_{\rm m}}$ 



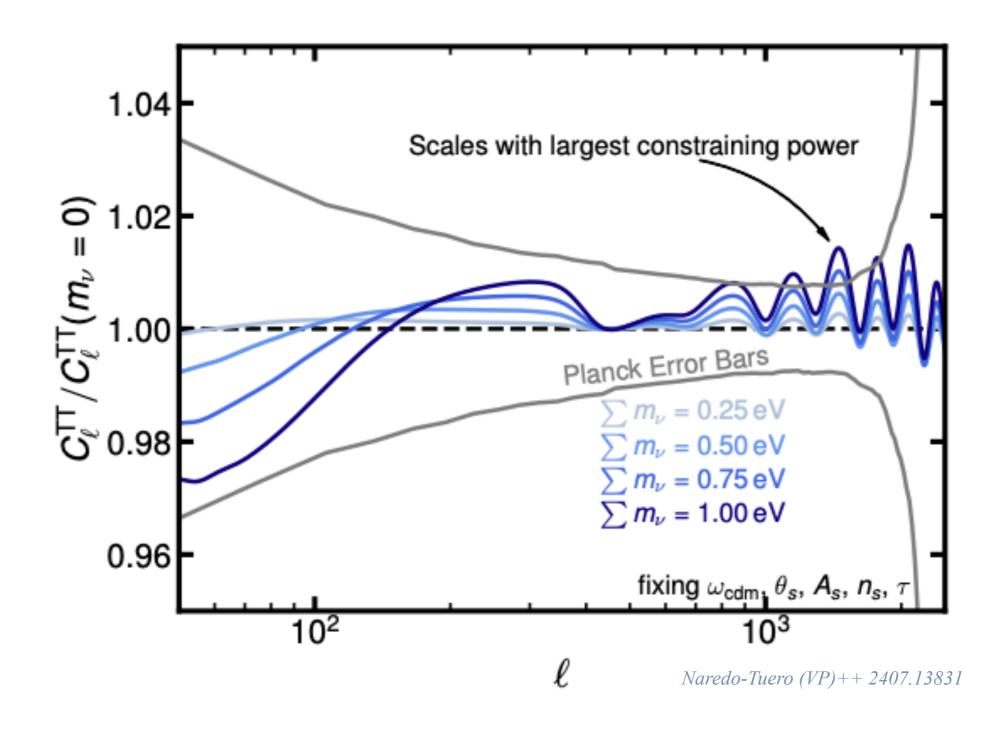
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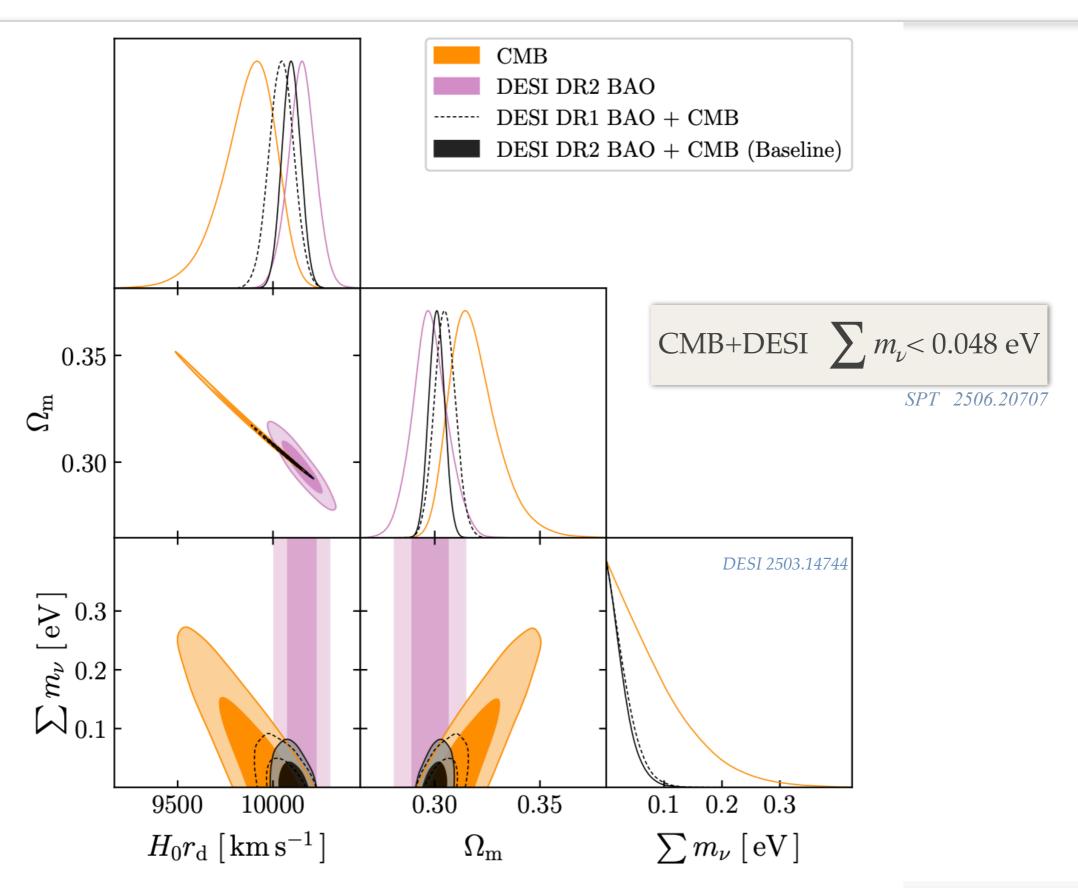


# $m_{\nu}$ affects the CMB power spectrum

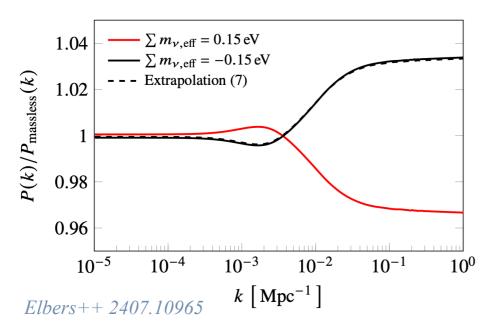
For  $\sum m_{\nu} < 1$  eV, the main effect of neutrinos is to alter the lensing in the CMB



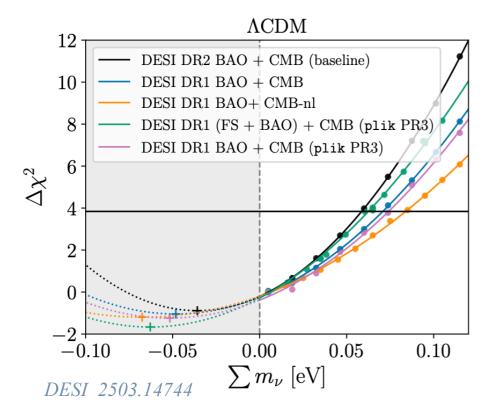
#### Strong constraints to $M_{ u}$ from background effect



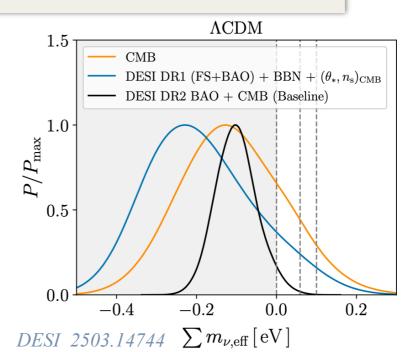
• When allowing for (effective)  $m_{\nu} < 0 \text{ eV}$ 



Also seen in Frequentist analyses



$$\sum m_{\nu,\text{eff}} = -0.101^{+0.047}_{-0.056} \text{ eV}$$



Planck + DESI-Y1

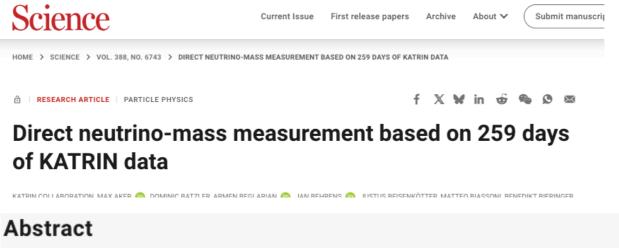
CamSpec22-PR4

Planck18-PR3

Hillipope Hilli

*Naredo-Tuero (VP)++ 2407.13831* 

#### A statistical fluke?



That neutrinos carry a nonvanishing rest mass is evidence of physics beyond the Standard Model of elementary particles. Their absolute mass holds relevance in fields from particle physics to cosmology. We report on the search for the effective electron antineutrino mass with the KATRIN experiment. KATRIN performs precision spectroscopy of the tritium  $\beta$ -decay close to the kinematic endpoint. On the basis of the first five measurement campaigns, we derived a best-fit value of  $m_{\nu}^2 = -0.14^{+0.13}_{-0.15}$  eV<sup>2</sup>, resulting in an upper limit of  $m_{\nu}$  < 0.45 eV at 90% confidence level. Stemming from 36 million electrons collected in 259 measurement days, a substantial reduction of the background level, and improved systematic uncertainties, this result tightens KATRIN's previous bound by a factor of almost two.

- Within Katrin, it is attributed to under fluctuations compatible with statistical fluke
- There are two main reasons driving this preference in CMB data, both still compatible with fluke

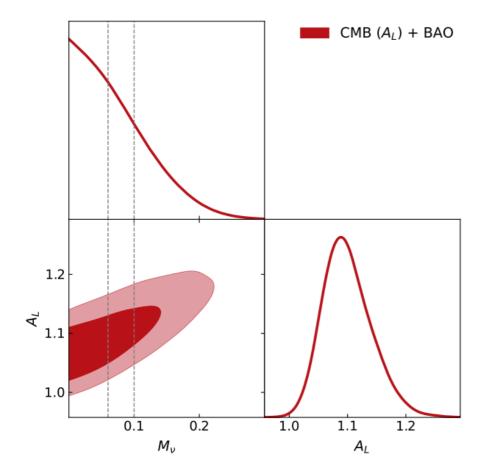
# $m_{\nu} < 0$ is due to a "lensing anomaly" in Planck

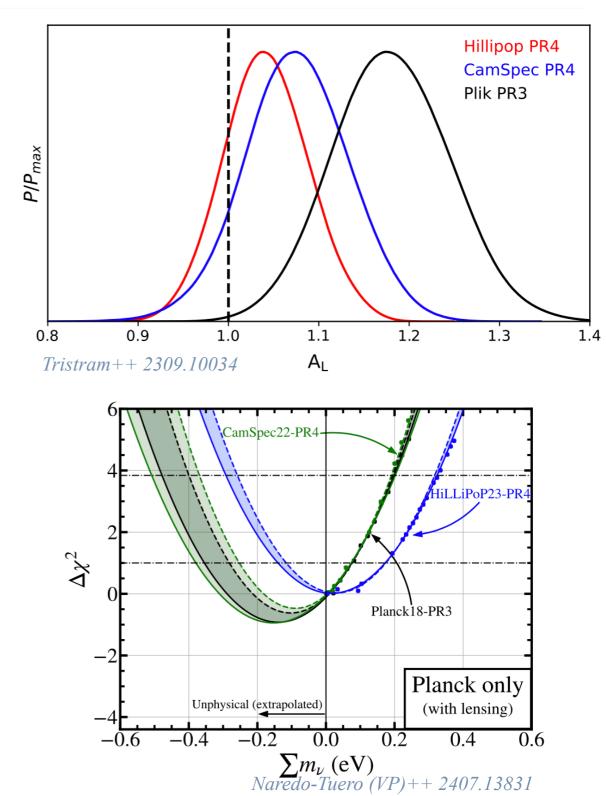
• Planck 2018 has an anomalous amount of lensing

$$C_{\ell}^{\phi\phi} \to A_L C_{\ell}^{\phi\phi}, \ A_L = 1.180 \pm 0.065$$

• This anomaly went down with PR4

$$A_L = 1.039 \pm 0.052$$





• The preference for  $m_{\nu} < 0$  goes away when using a newer version *Planck* data (NPIPE)

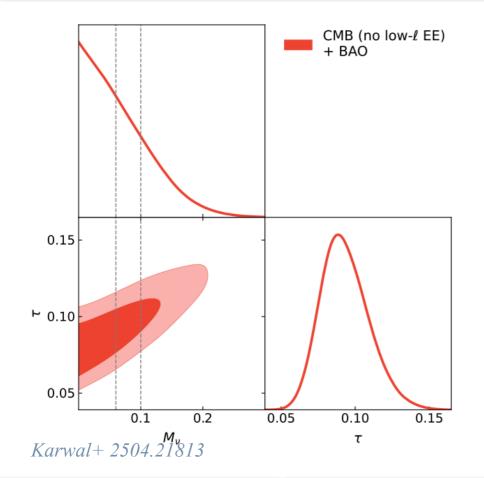
### The role of the optical depth $\tau$

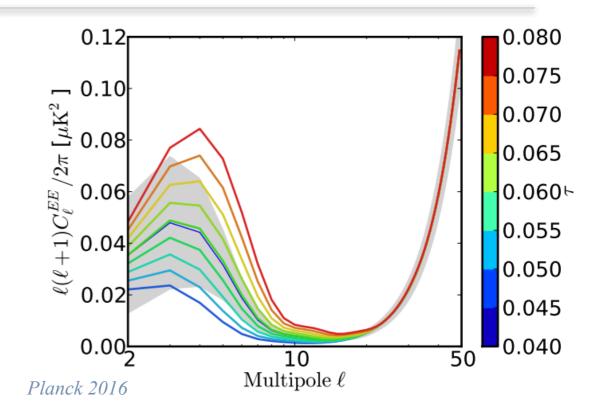
- The amplitude is correlated with the optical depth to reionization  $\tau$   $C_{\ell} \propto A_{s} \exp(-2\tau)$ 
  - The low- $\ell$  EE spectrum constrains

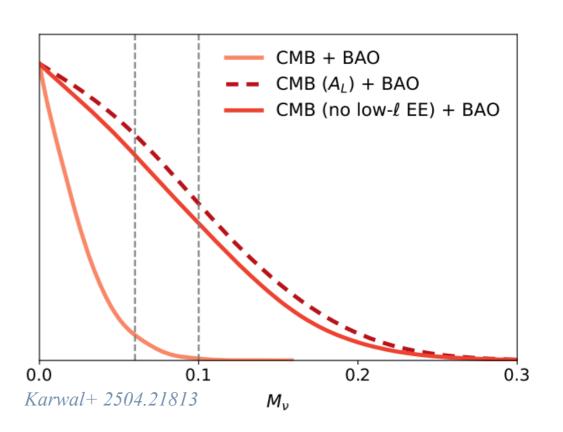
$$\tau = 0.0592 \pm 0.0062$$

• Without low- $\ell$  EE information,

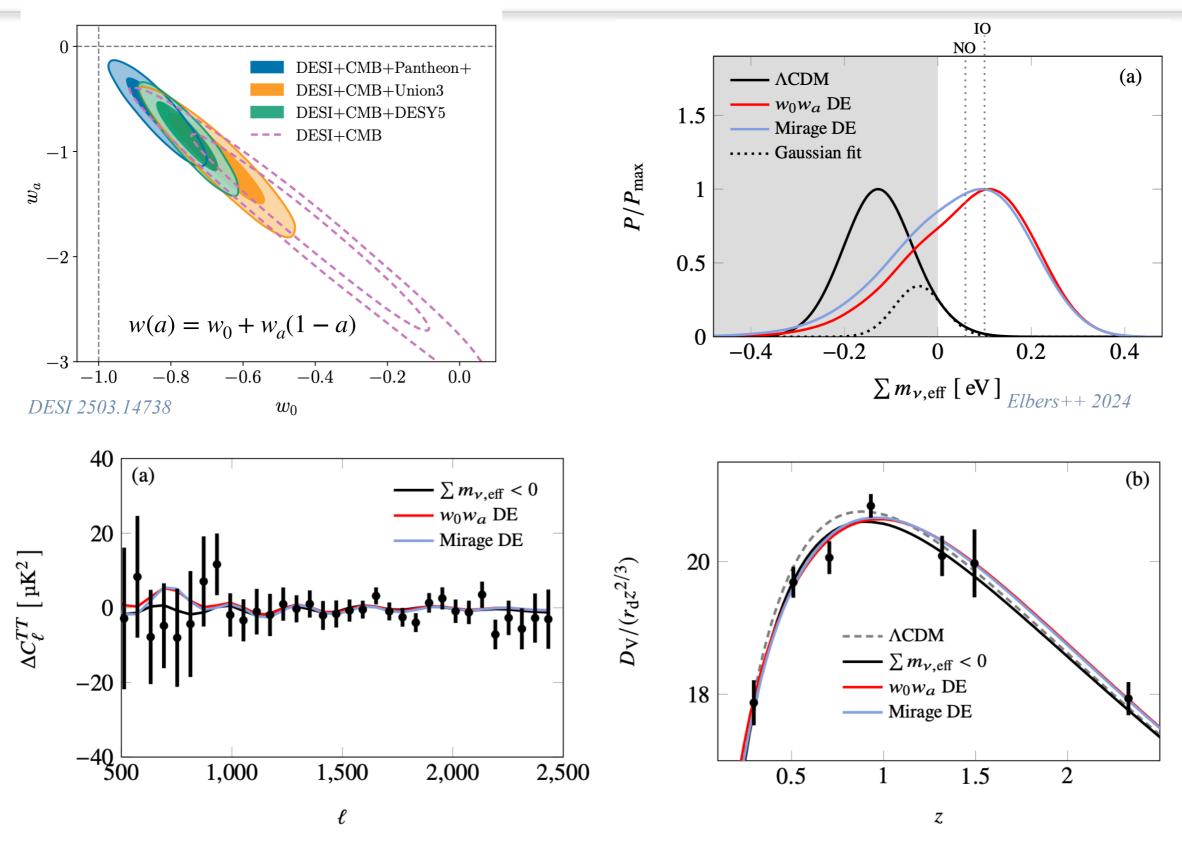
$$\tau = 0.09 \pm 0.012$$
,  $\sum m_{\nu} < 0.2 \text{ eV}$ 







### $m_{\nu}$ is degenerate with the effect of dark energy



# Summary of neutrino mass bounds

DESI 2503.14744, 2504.18464

• Despite anomalies, it is hard to get rid of the neutrino mass bound from cosmology

Flat 
$$\Lambda$$
CDM (3 deg  $\nu$ 's)

Planck+DESI  $\sum m_{\nu} < 0.064 \text{ eV}$ 

Planck+ ACT +DESI  $\sum m_{\nu} < 0.06 \text{ eV}$ 

P+ ACT +SPT +DESI  $\sum m_{\nu} < 0.048 \text{ eV}$ 

$$w_0 w_a \text{CDM}$$
 (3 deg  $\nu'$ s)

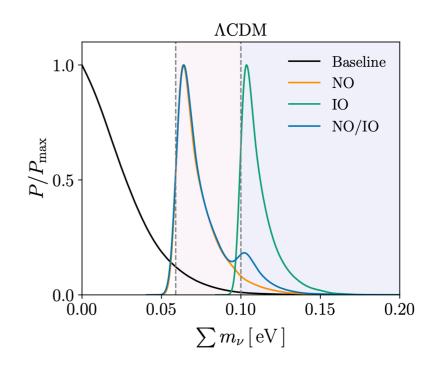
Planck+DESI  $\sum m_{\nu} < 0.163 \text{ eV}$ 

Planck+ ACT +DESI  $\sum m_{\nu} < 0.152 \text{ eV}$ 

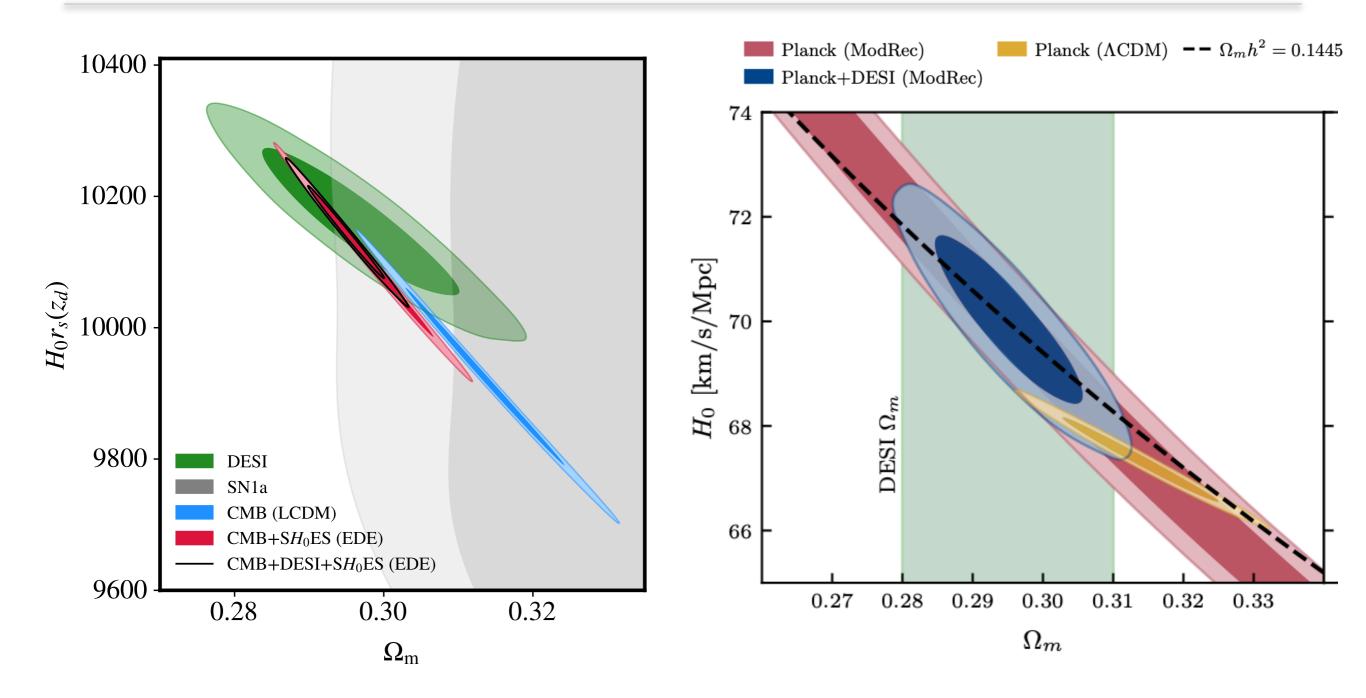
Planck+ ACT + SN1a  $\sum m_{\nu} < 0.117$ -0.139 eV

• Enforcing neutrino mass ordering leads to a preference for NO over IO at 10:1

Planck + DESI + NuFIT6.0 (
$$\Lambda$$
CDM)  
NO  $\sum m_{\nu} < 0.101 \text{ eV} \Rightarrow m_{l} < 0.023 \text{ eV}$   
IO  $\sum m_{\nu} < 0.133 \text{ eV} \Rightarrow m_{l} < 0.024 \text{ eV}$ 



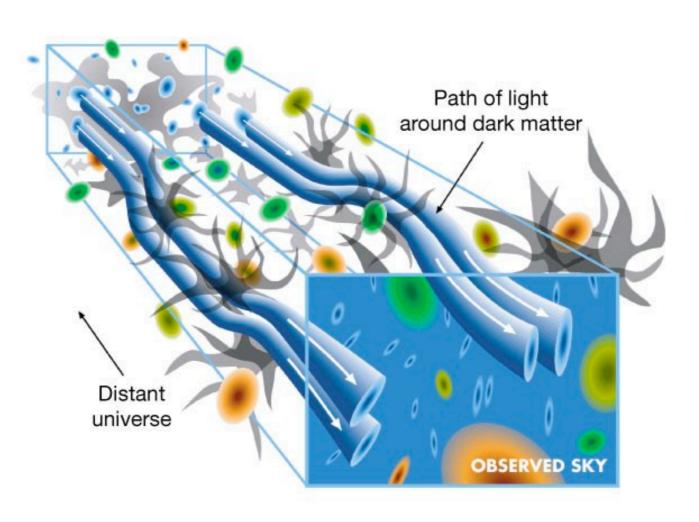
### Early universe solution to the BAO tension

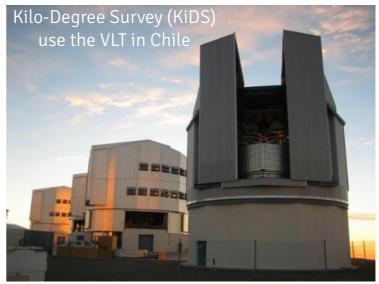


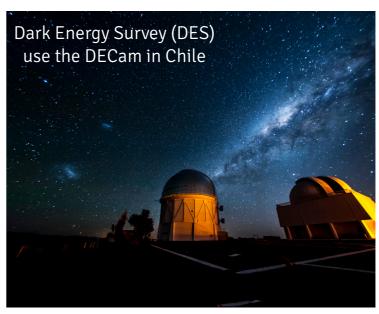
- Early universe solution can reduce  $H_0 r_d$  and  $\Omega_m$
- An alternative explanation to DESI results? Lynch&Chluba 2406.10202, Chaussidon++ 2503.24343

### Galaxy weak lensing

"Weak gravitational lensing" observations measure the distortions in the shape of galaxies



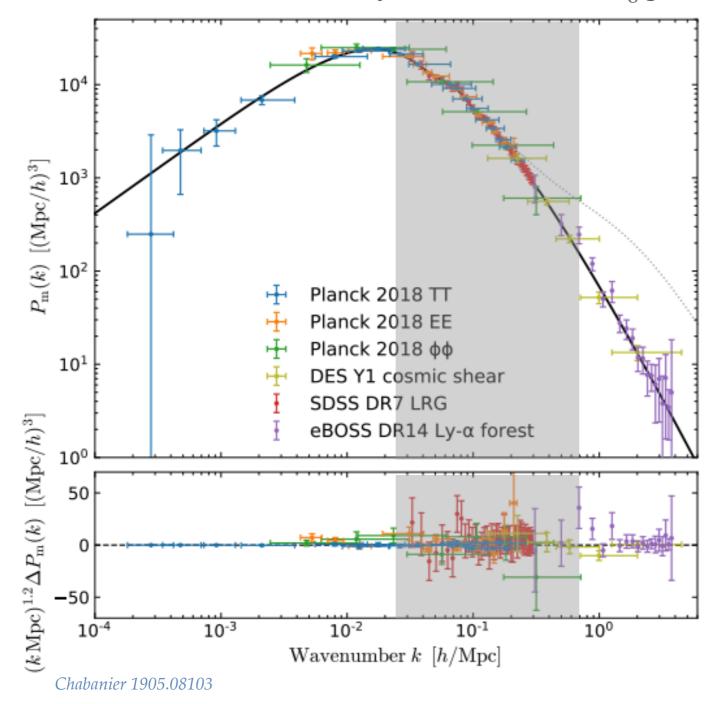




- There are three main WL observatory: DES, KiDS and HSC (Hawaii)
- New: Euclid!

### The $S_8$ parameter

• WL observations are mostly sensitive to the ' $S_8$  parameter'.

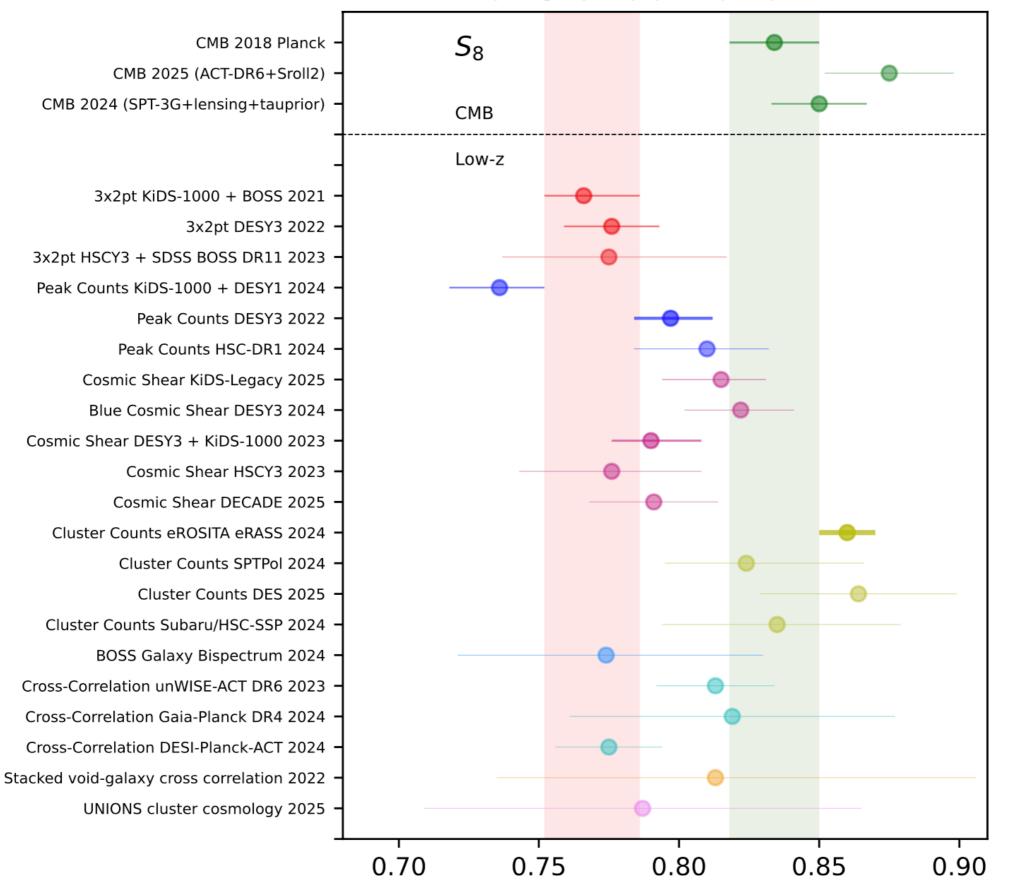


$$S_8 \equiv \sigma_8 \left(\frac{\Omega_m}{0.3}\right)^{0.5}$$

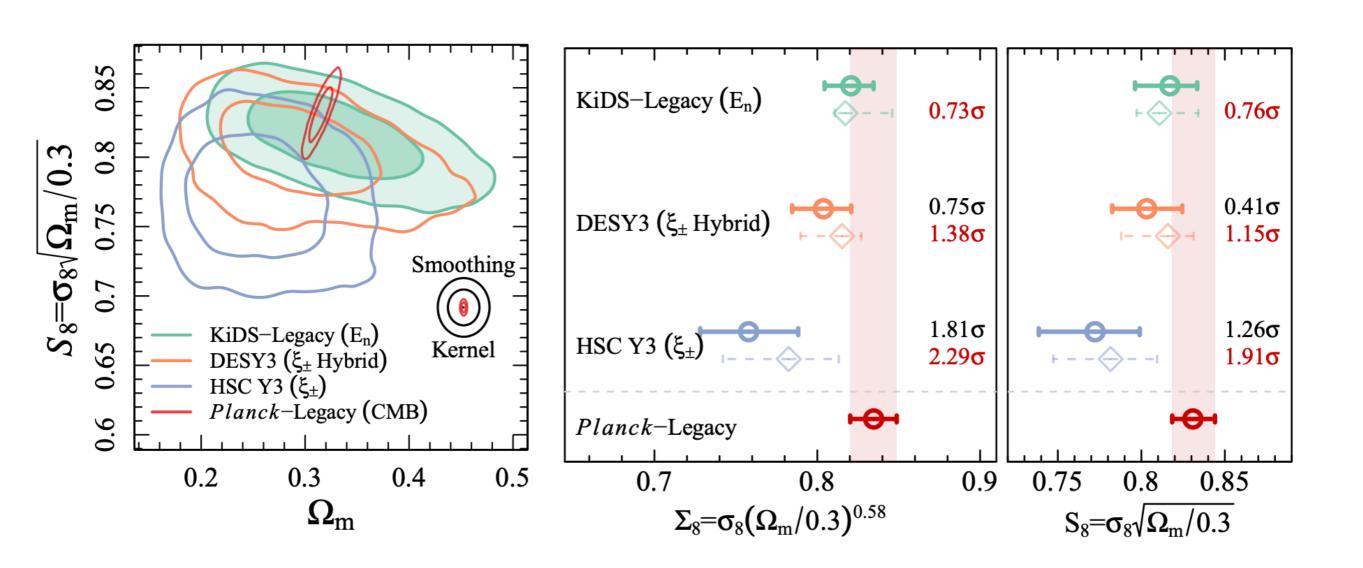
$$\sigma_8^2 = \int_0^\infty \frac{k^3}{2\pi^2} P_{\text{lin}}(k) W^2(kR) d\ln k$$

• The  $S_8$  parameter quantifies how "clumpy" the universe is on scales of  $\sim 30$  million-ly

#### The S8 tension

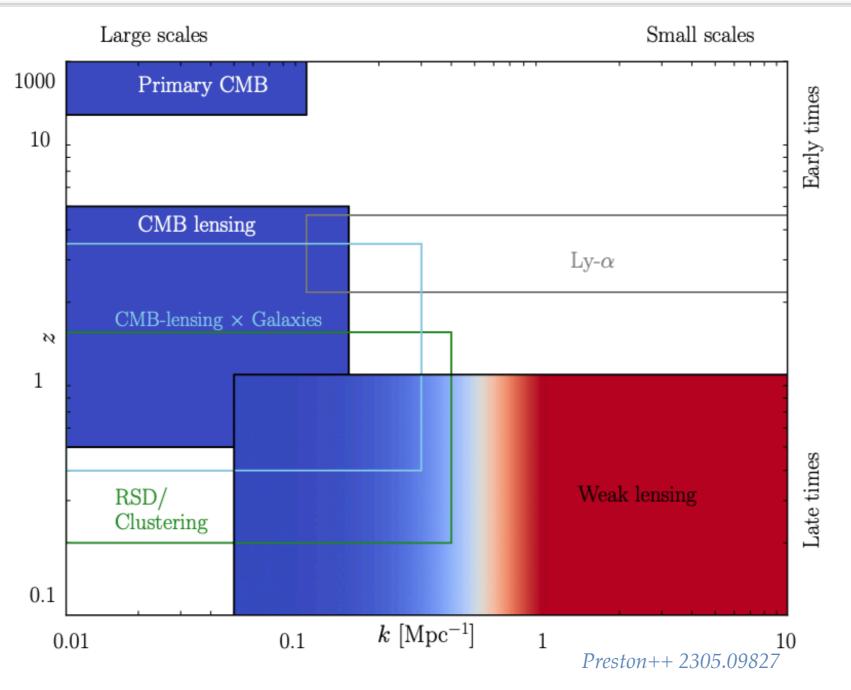


#### The $S_8$ tension revisited

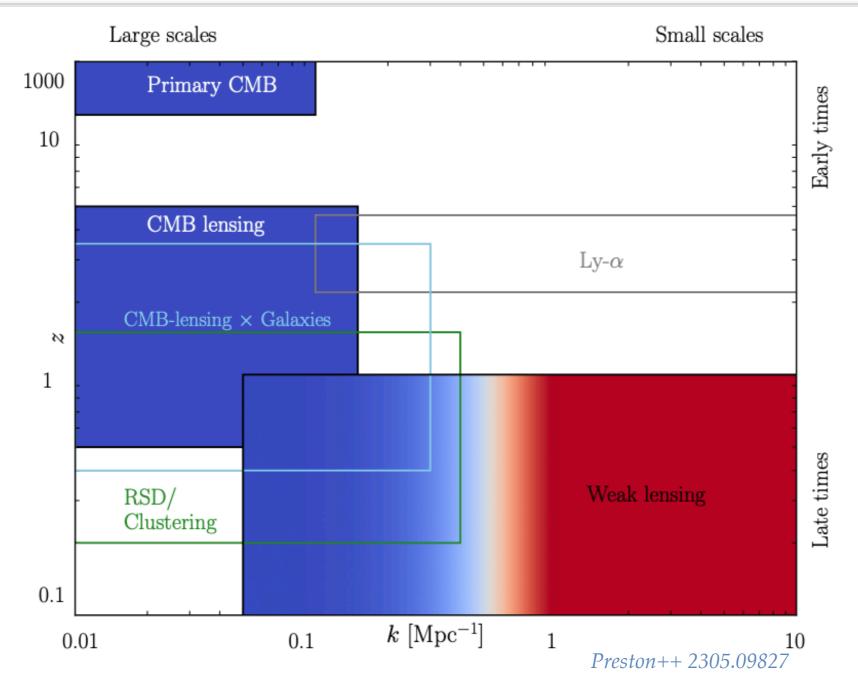


Latest results from KIDS Legacy: improved redshift calibration has resolved the tension

#### The general picture: what do we know about P(k,z)?



#### The general picture: what do we know about P(k,z)?



- Weak lensing measure smaller scales than galaxy cluster number counts! Power suppression at  $k \gtrsim 0.5$  h/Mpc?
- Lyman- $\alpha$  data may or may not favor a power suppression at  $z \sim 3$  and  $k \sim 0.7$  Mpc  $^{-1}$

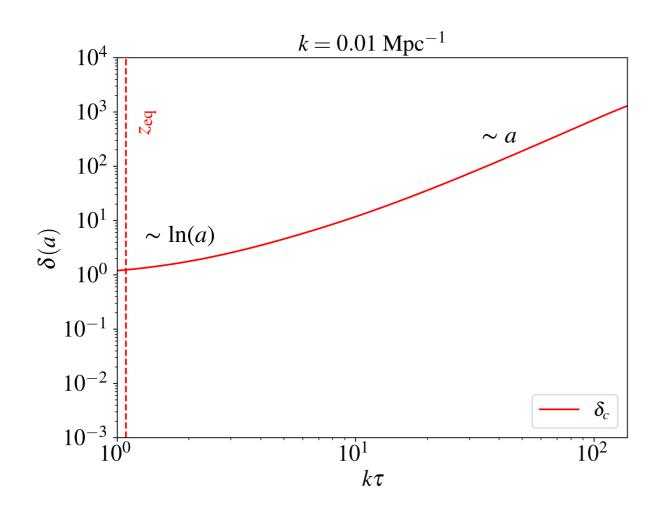
Goldstein++ 2303.00746 , K. Rogers & VP 2311.16377

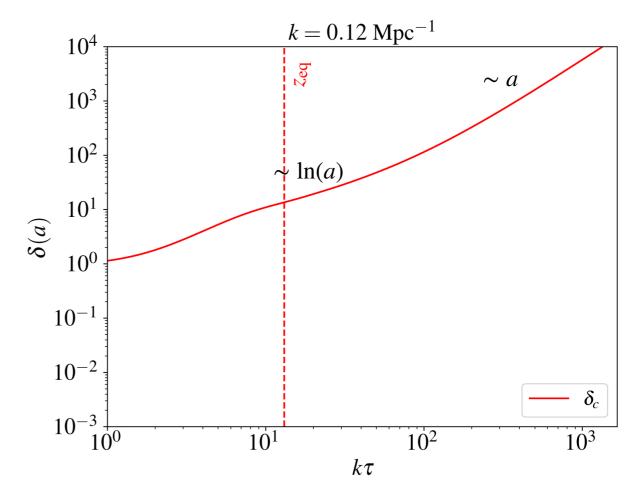
### The growth of CDM density perturbations

• Continuity & Euler equations for CDM perturbations (synchronous gauge):

$$\delta'_c = -\theta_c - \frac{1}{2}h', \qquad \theta'_c = -\frac{a'}{a}\theta_c \qquad \delta \equiv \delta \rho/\rho, \quad \theta \equiv kv \equiv 0$$

• For pressure-less, collision-less fluid: within the horizon, scale-independent growth.





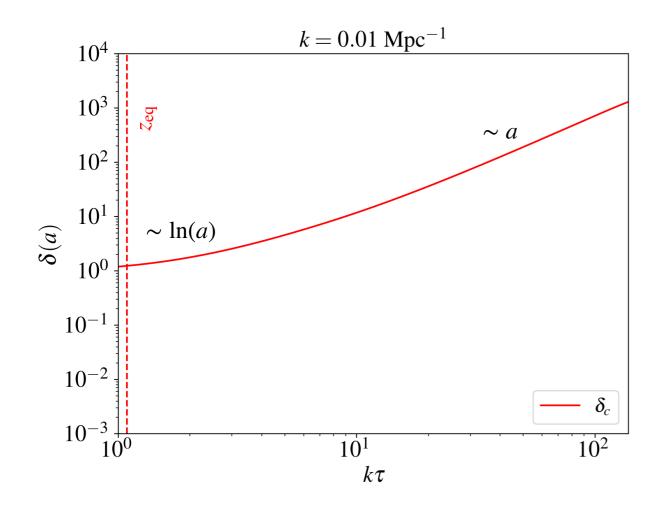
#### How to decrease power at small scales

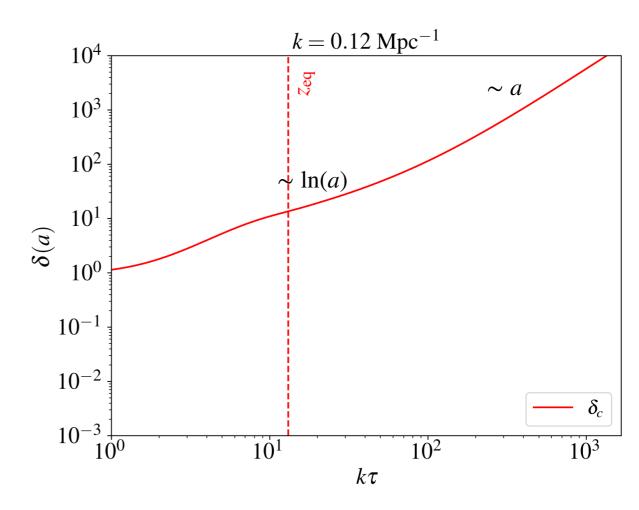
• If non-zero pressure:  $w = \frac{p}{\rho}$ ,  $c_a^2 = \frac{p'}{\rho'}$ ,  $c_s^2 = \frac{\delta p}{\delta \rho}$  can be non-zero

$$\delta'_{\text{ncdm}} = -(1+w)\left(\theta_{\text{ncdm}} - \frac{1}{2}h'\right) - 3\frac{a'}{a}(c_s^2 - w)\delta_{\text{ncdm}}, \quad \theta'_{\text{ncdm}} = -\frac{a'}{a}(1 - 3c_a^2)\theta_{\text{ncdm}} + \frac{c_s^2}{1+w} - k^2\sigma_{\text{ncdm}}$$

• One can write a formal equation driving the growth of perturbation on small scales:

$$k\tau \gg 1 \Rightarrow \delta_{\text{ncdm}}'' + \frac{a'}{a}\delta_{\text{ncdm}}' + (k^2 - k_J^2)c_s^2\delta_{\text{ncdm}} = 0$$
  $k_J \equiv \sqrt{\frac{3(a'/a)}{c_s^2}}$ 





#### The case of neutrinos

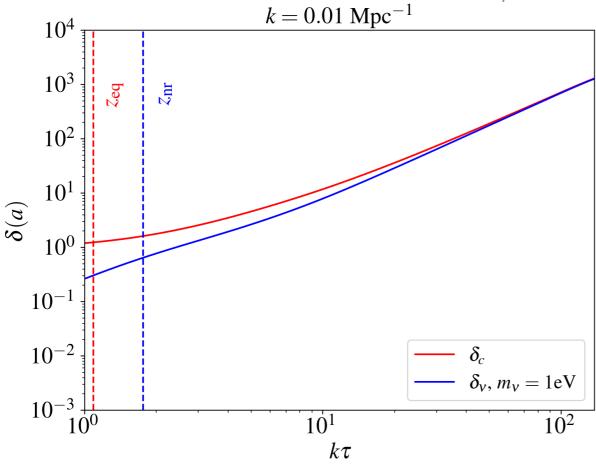
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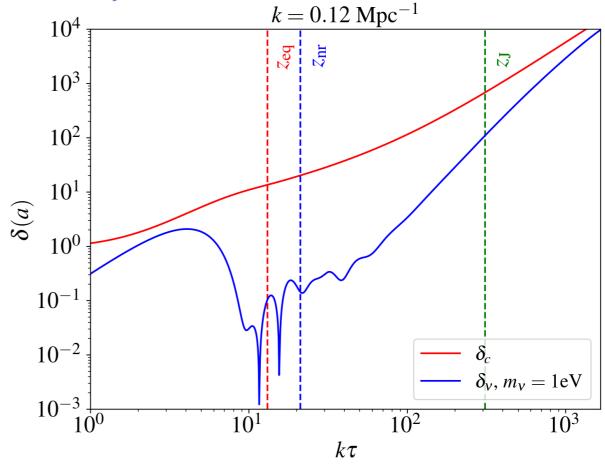
$$\delta'_{\text{ncdm}} = -(1+w)\left(\theta_{\text{ncdm}} - \frac{1}{2}h'\right) - 3\frac{a'}{a}(c_s^2 - w)\delta_{\text{ncdm}}, \quad \theta'_{\text{ncdm}} = -\frac{a'}{a}(1 - 3c_a^2)\theta_{\text{ncdm}} + \frac{c_s^2}{1+w} - k^2\sigma_{\text{ncdm}}$$

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  $k_J \equiv \sqrt{\frac{3(a'/a)}{c_s^2}}$ 

Example: thermal neutrinos with  $m_{\nu} = 1 \text{ eV}$ 



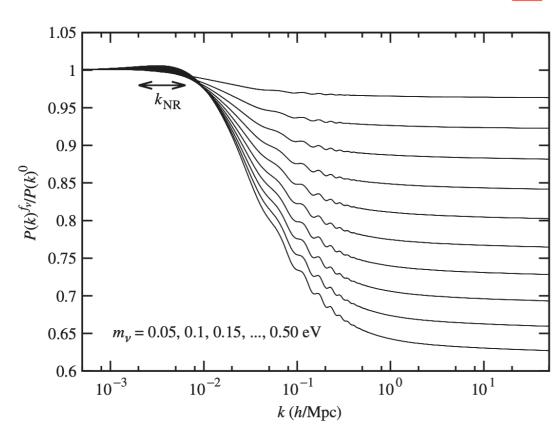


### Could $\nu$ 's explain the $S_8$ tension?

Power suppression:

$$k \ge k_{\rm nr} \equiv 0.01 \left(\frac{m_{\nu}}{1 {\rm eV}}\right)^{1/2} \left(\frac{\Omega_{\rm m}}{0.3}\right)^{1/2} h {\rm Mpc}^{-1}$$
 with amplitude  $\frac{\Delta P}{P} \simeq -8 \frac{\omega_{\nu}}{\omega_{\rm m}}$ 

Need  $\sum m_{\nu} \sim 0.2 \text{ eV to explain } S_8$ 

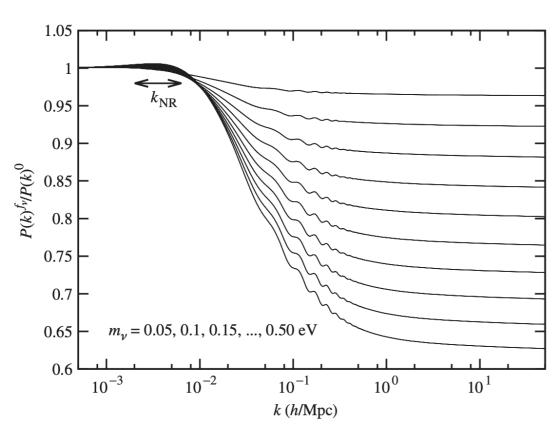


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Power suppression:

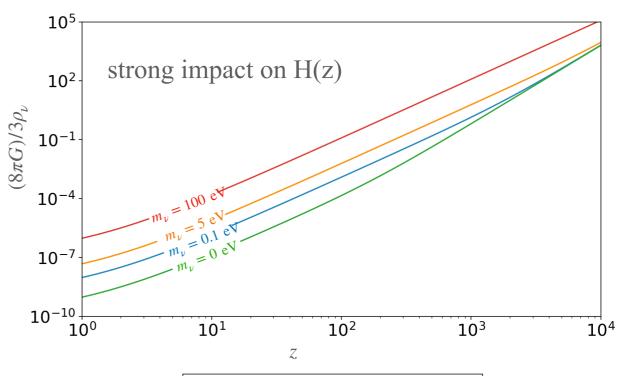
$$k \ge k_{\rm nr} \equiv 0.01 \left(\frac{m_{\nu}}{1 \, {\rm eV}}\right)^{1/2} \left(\frac{\Omega_{\rm m}}{0.3}\right)^{1/2} h {\rm Mpc}^{-1}$$
 with amplitude  $\frac{\Delta P}{P} \simeq -8 \frac{\omega_{\nu}}{\omega_{\rm m}}$ 

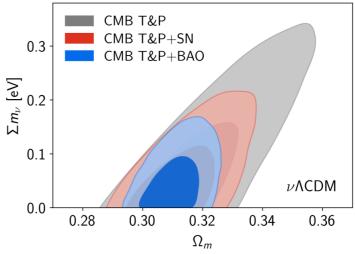
#### Need $\sum m_{\nu} \sim 0.2 \text{ eV to explain } S_8$



*Planck* 2018 + BAO < 0.12eV

Planck 1807.06205



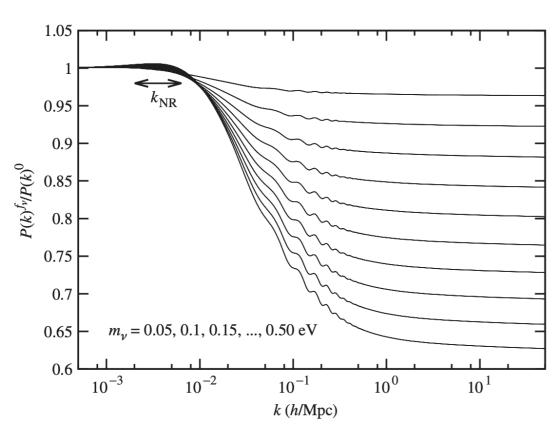


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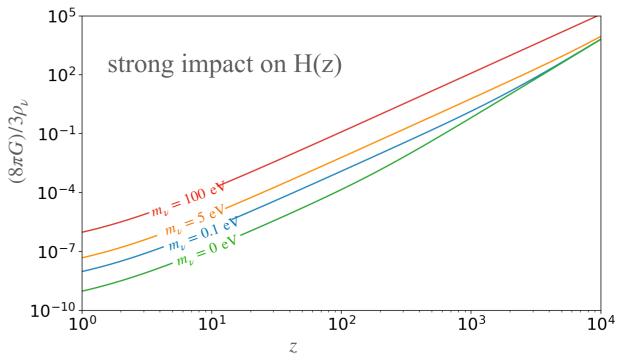
Planck 1807.06205

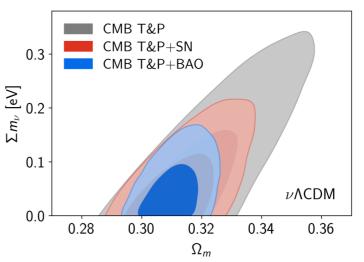
 $Planck\ 2018 + BAO + Ly - \alpha < 0.089eV$ 

 $Palanque ext{-}Delabrouille++ 1911.09073$ 

 $Planck\ 2018 + BOSS + eBOSS < 0.082eV$ 

Brieden++ 2204.11868, Simon++ 2210.14931



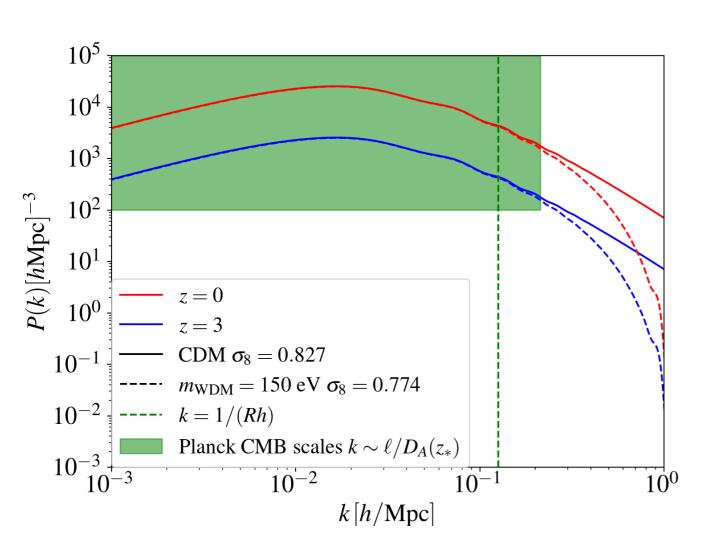


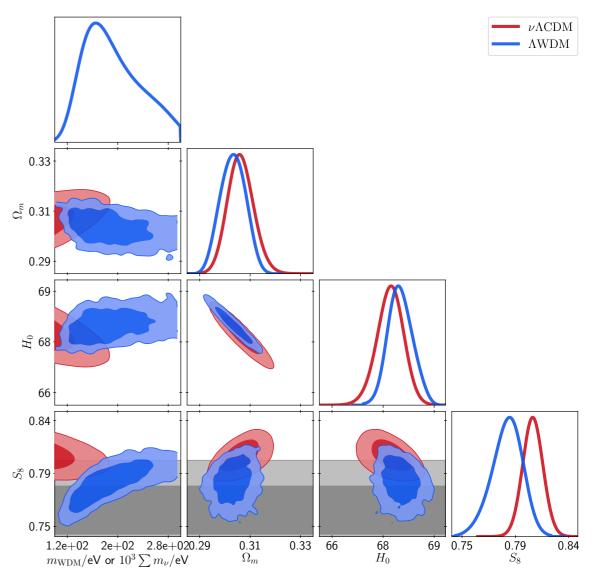
#### The case of warm dark matter

For WDM: suppression at  $k \ge k_{\rm FS} \sim 5~{\rm Mpc^{-1}} \bigg(\frac{m_{\rm WDM}}{1{\rm keV}}\bigg) \bigg(\frac{T_{\nu}}{T_{\rm WDM}}\bigg)$ 

*Viel++ astro-ph/0501562* 

• Requiring  $k_{\rm FS} \ge 0.1 h/{\rm Mpc}$  and  $\omega_{\rm ncdm} = \omega_{\rm cdm}$  leads to  $m_{\rm WDM} \sim 100 {\rm eV}$ .

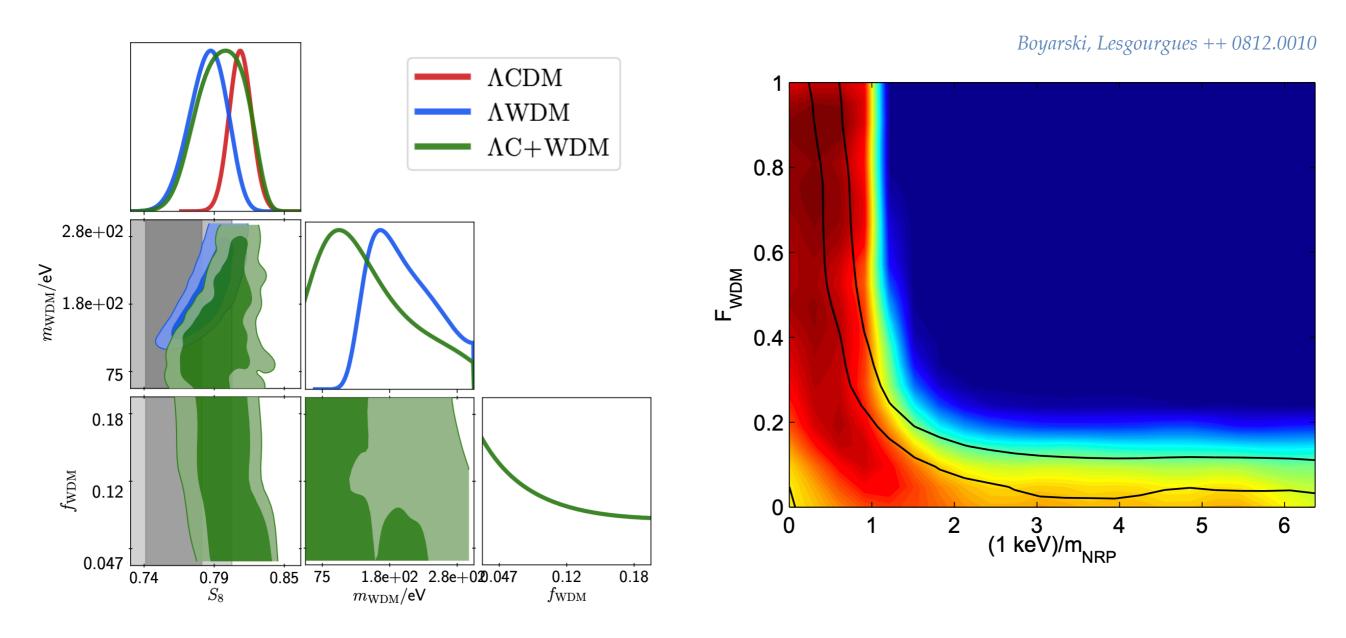




- CMB + BAO +  $S_8 m_{\text{WDM}} \simeq 192.2^{+31}_{-60} \text{ eV}.$
- Yet, exponential cutoff 'a la WDM' are excluded:  $m_{\rm WDM} > 5$  keV.

Palanque-Delabrouille++ 1911.09073

### Fraction of Warm Dark Matter and $S_8$



- A fraction of thermal WDM with  $m \sim \mathcal{O}(100 \text{eV})$  could explain low  $S_8$ : Ly- $\alpha$  restricts  $f_{\text{WDM}} \lesssim 0.2$
- How to generate a fraction of WDM?
- Nb: similar results with fraction of fuzzy dark matter / tightly coupled fraction of DM

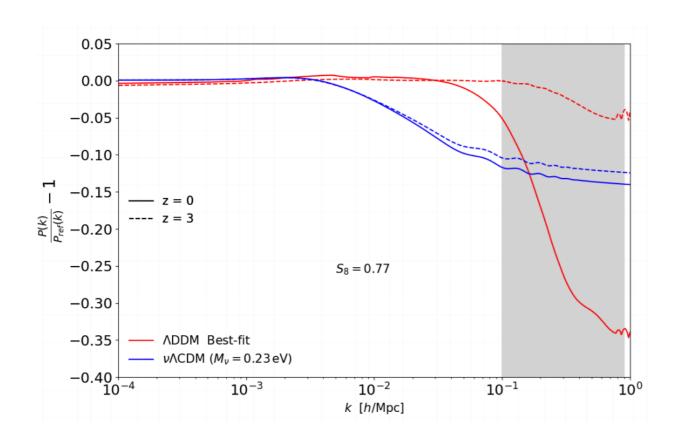
Chacko++ 1609.03569, Allali++ 2104.12798, Laguë++ 2104.07802

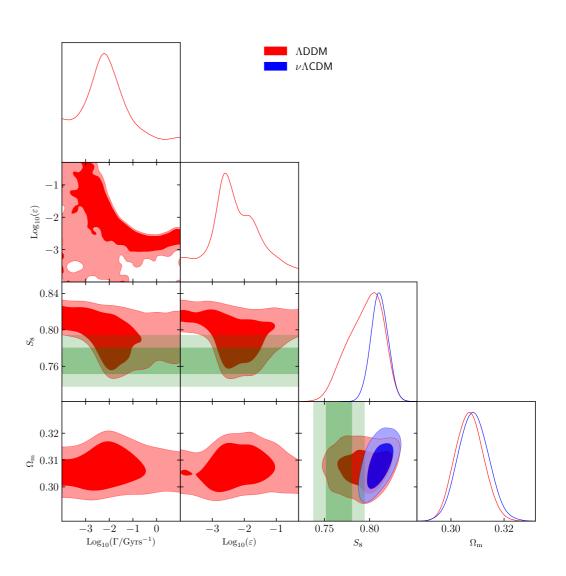
See also Gariazzo++ 1704.02991

### DM decays generate a late-time suppression

- Standard neutrinos are largely excluded as a solution: need  $M_{\nu} \simeq 0.23$  eV but  $M_{\nu} \lesssim 0.08 0.12$  eV Palanque-Delabrouille++ 1911.09073; Brieden++ 2204.11868; Simon (VP)++ 2210.14931
- ΛCDM + two free parameters: the CDM decay rate Γ and the branching ratio to DR  $\varepsilon = \left(\frac{1}{2} \frac{m_{\text{WDM}}^2}{m_{\text{DCDM}}^2}\right) \simeq v_{\text{wdm}}/c$

$$\begin{split} \dot{\bar{\rho}}_{\rm dcdm} &= -3\mathcal{H}\bar{\rho}_{\rm dcdm} - a\Gamma\bar{\rho}_{\rm dcdm}, \\ \dot{\bar{\rho}}_{\rm dr} &= -4\mathcal{H}\bar{\rho}_{\rm dr} + \varepsilon a\Gamma\bar{\rho}_{\rm dcdm}, \\ \dot{\bar{\rho}}_{\rm wdm} &= -3(1+w)\mathcal{H}\bar{\rho}_{\rm wdm} + (1-\varepsilon)a\Gamma\bar{\rho}_{\rm dcdm}. \end{split}$$





• DM with  $\Gamma^{-1} \simeq 55(\varepsilon/0.007)^{1.4}$  Gyrs can explain low  $S_8$  (1.3 $\sigma$  agreement)

Abellan++ 2008.09615 & 2104.03329

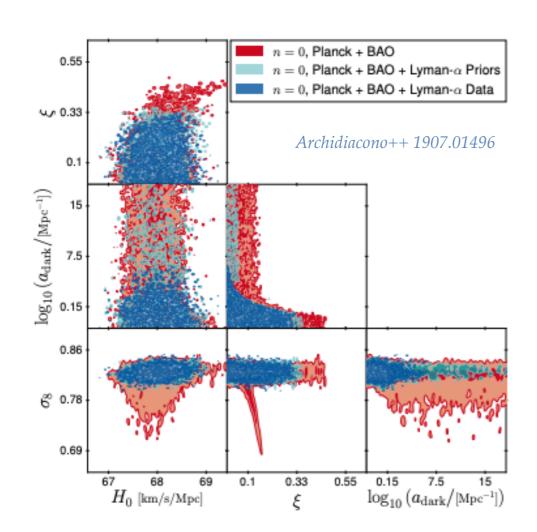
See also Davari&khosravi 2203.09439, Clark++ 2110.09562

### The $S_8$ tension is a drag...

$$DM \le DR$$

$$\dot{\delta}_{\mathrm{DM}} + heta_{\mathrm{DM}} - 3\dot{\phi} = 0,$$
  $\dot{ heta}_{\mathrm{DM}} - k^2 c_{\mathrm{DM}}^2 \delta_{\mathrm{DM}} + \mathcal{H} heta_{\mathrm{DM}} - k^2 \psi = \Gamma_{\mathrm{DM-DR}} \left( heta_{\mathrm{DM}} - heta_{\mathrm{DR}} \right),$ 

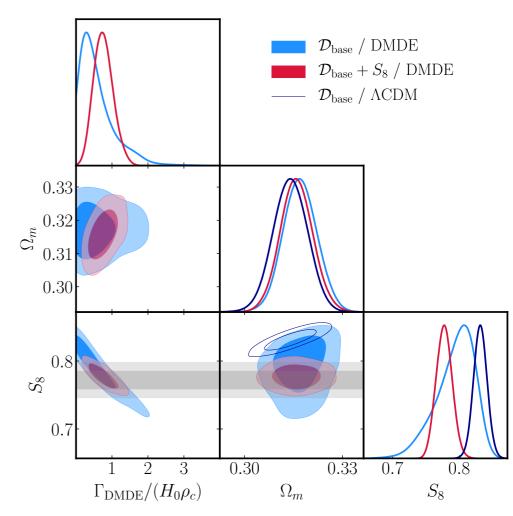
$$\Gamma_{
m DR-DM} = -\Omega_{
m DM} h^2 a_{
m dark} \left(rac{1+z}{1+z_d}
ight)^n, \qquad \quad \xi = T_{
m DR}/T_{\gamma}$$



#### $DM \le DE$

VP, Bernal, Kovetz, Kamionkowski 2209.06217

$$\begin{aligned} \theta_{\mathrm{DM}}' &= -\frac{a'}{a} \theta_{\mathrm{DM}} + k^2 \psi + \Gamma_{\mathrm{DMDE}}(a) (\theta_{\mathrm{DE}} - \theta_{\mathrm{DM}}), \\ \theta_{\mathrm{DE}}' &= -(1 - 3c_{s,\mathrm{DE}}^2) \frac{a'}{a} \theta_{\mathrm{DE}} + \frac{k^2 c_{s,\mathrm{DE}}^2}{(1 + w_{\mathrm{DE}})} \delta_{\mathrm{DE}} \\ &+ k^2 \psi - \Gamma_{\mathrm{DMDE}}(a) R(\theta_{\mathrm{DE}} - \theta_{\mathrm{DM}}), \end{aligned}$$

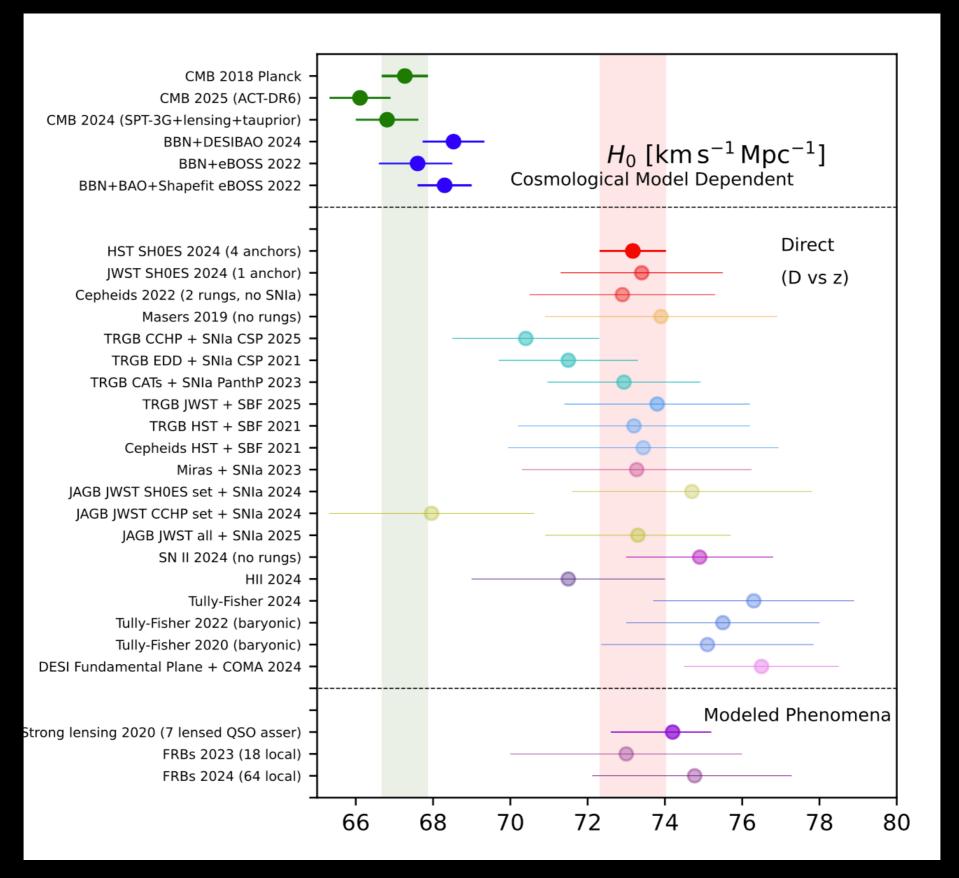


See also Di Valentino++ 1908.04281

Non-Abelian dark matter model, Cannibal dark matter, also with sub-component of strongly interacting DM

Buen-Abad++1505.03542, Lesgourgues++1507.04351, Heimersheim++ 2008.08486, Chacko++1609.03569, Buen-Abad++ 1708.09406, Raveri++ 1709.04877

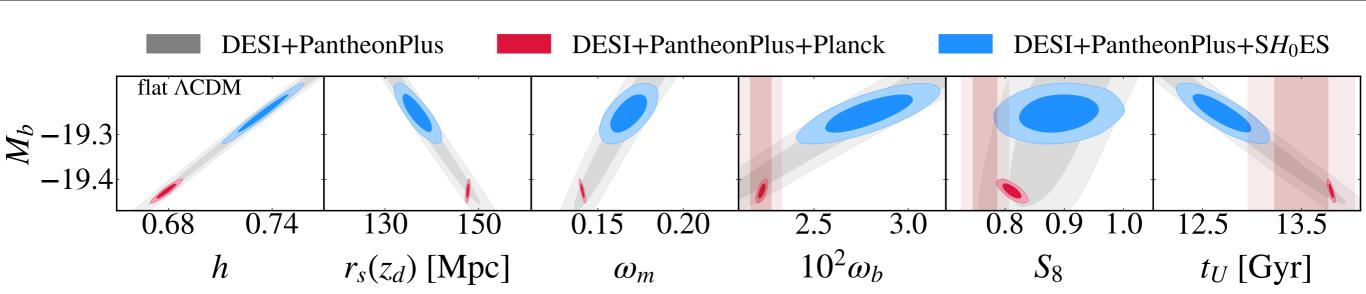
#### The Hubble tension



We are still tracking "unknown unknowns"

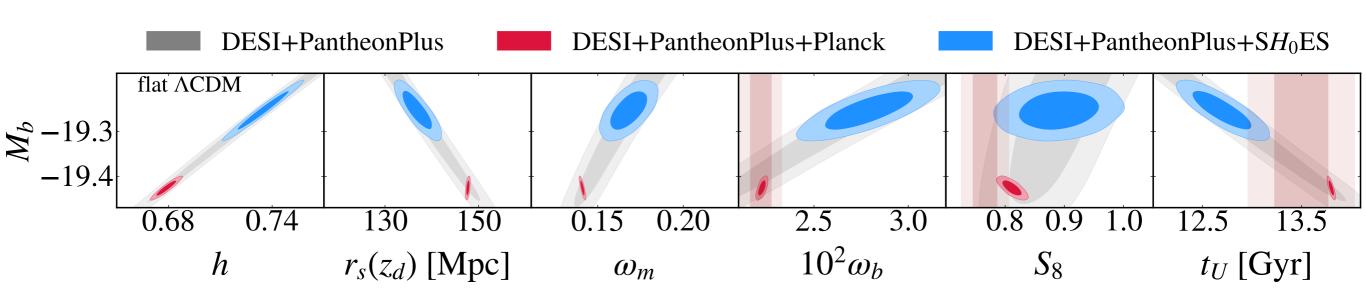
## Stop calling it $H_0$ tension!

- The " $H_0$  tension" should really be called cosmic calibration tension
- Calibrating the ladder has implications beyond the value of  $H_0$ : smaller  $t_U$ , larger  $\omega_m$ ,  $\omega_b$  and  $S_8$



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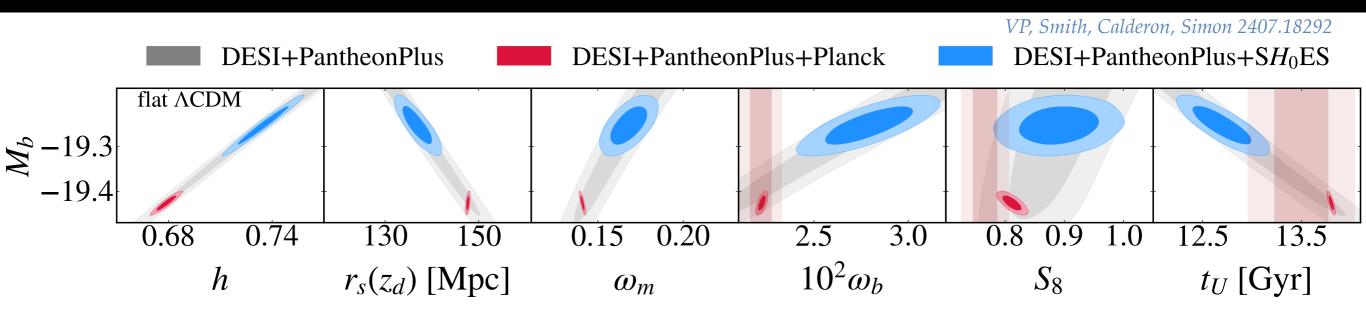
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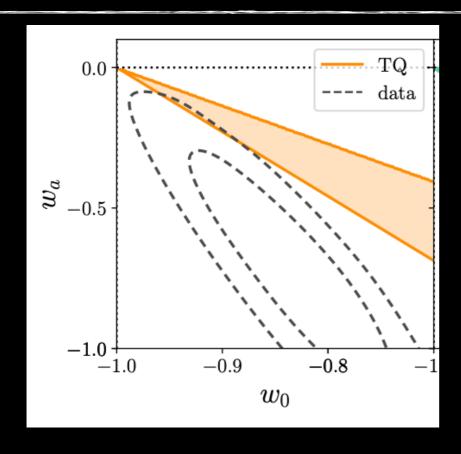
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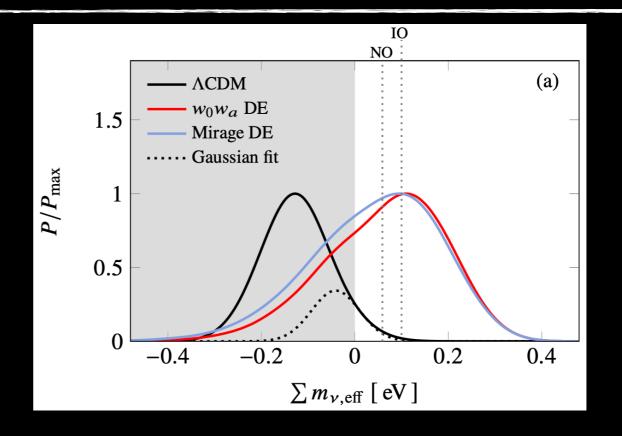


- ullet New physics must reduce the sound horizon  $r_{s}$  and accommodate larger  $\omega_{m}$ .
- Models affecting the pre-recombination expansion history favored but not fully successful
- Or maybe need new degrees of freedom at both early- and late-times?
  - We haven't found the solution yet, but there is a lot we understand!

#### A new DESI tension?

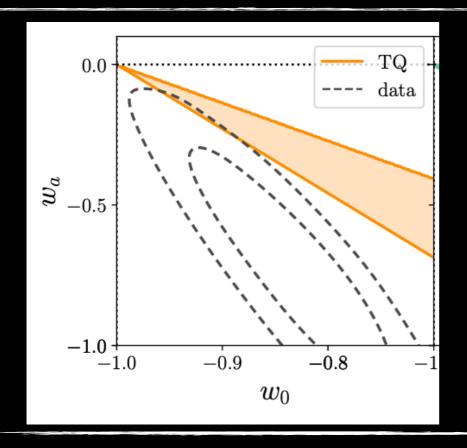
- DESI + CMB in 2.3  $\sigma$  tension, DESI+CMB in tension with SN1a at 2-3 $\sigma$
- This can be explained by Dynamical Dark Energy ==> non-canonical quintessence
- Mirage Dark Energy is favored: weird?

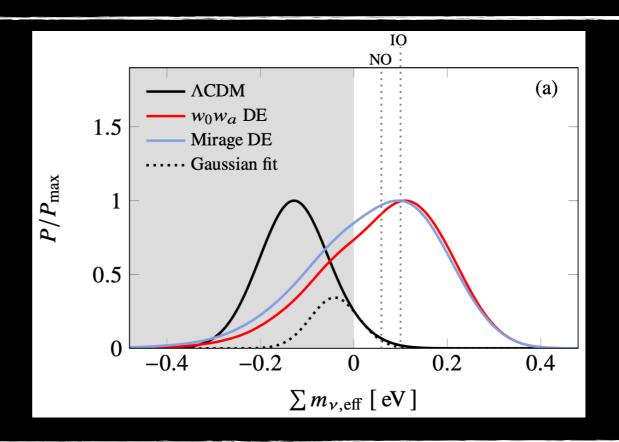




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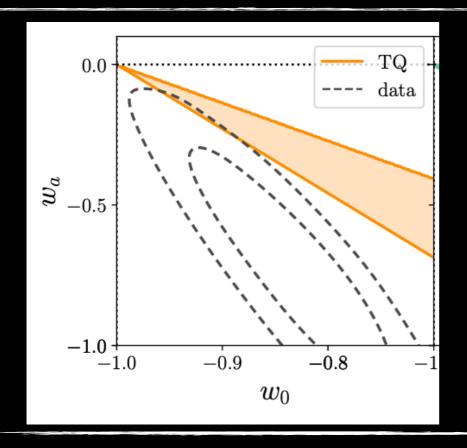


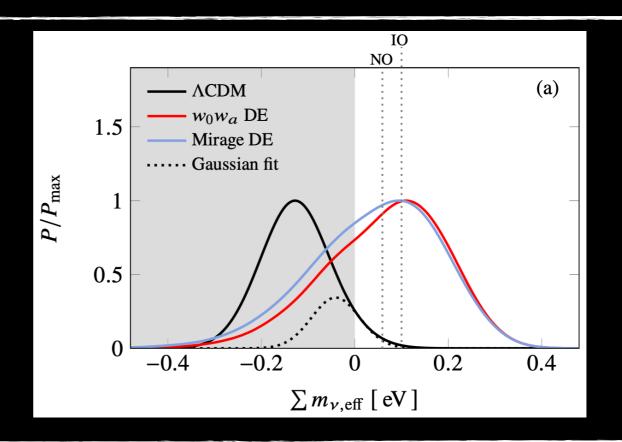


- Very strong bounds to  $M_{\nu}$  due to impact on  $d_{A}(z)$ : P-ACT-SPT +DESI  $M_{\nu}$ < 0.048 eV
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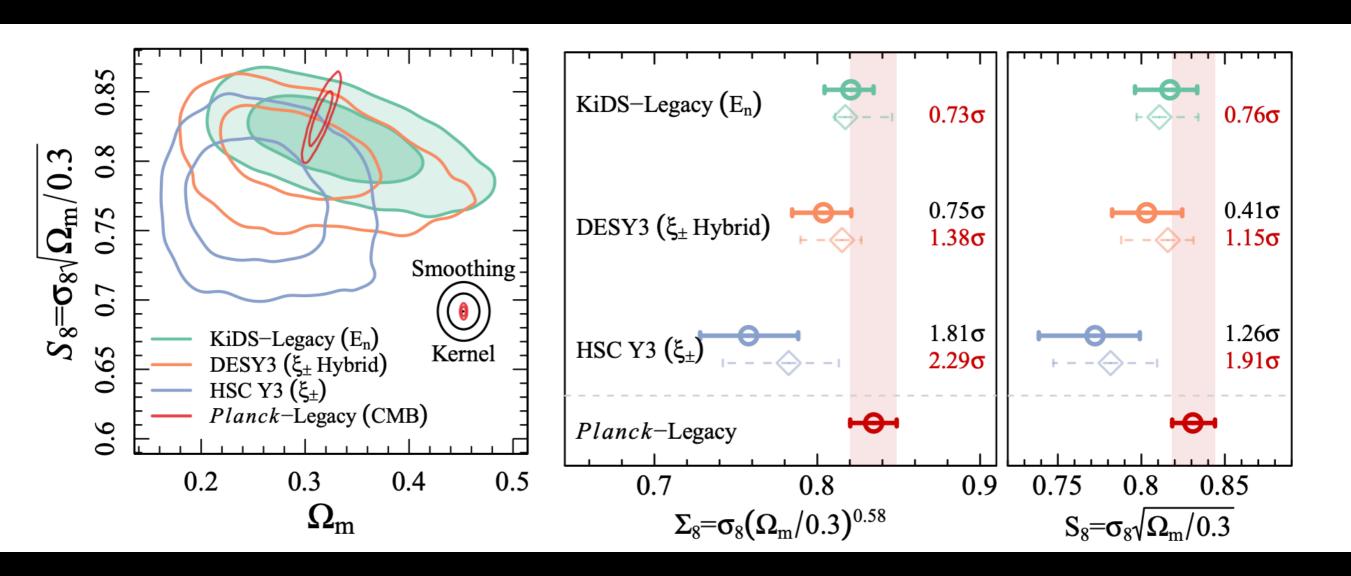
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  - ullet Early universe models could explain BAO data but not large  $\Omega_m$  from SNe

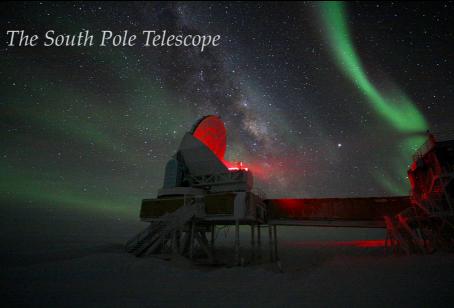
# The $S_8$ tension is gone?



- ullet But remember that the  $H_0$  tension may imply some residual tension ...
- See DM talks for more ideas for solutions

### Cosmic tension: where are we going next?





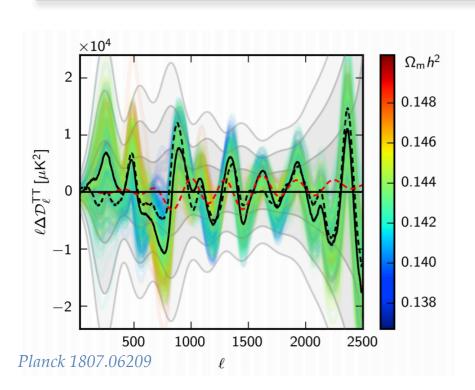






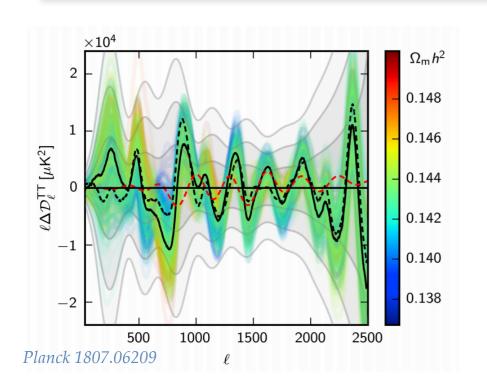


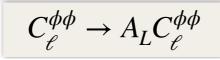
- New CMB data are coming: very sensitive to new physics around recombination!
- New LSS data are coming: measure P(k, z), constrain baryonic feedback, neutrino masses?
- New JWST data + age of the universe? + gravitational wave measurements of  $H_0$ .

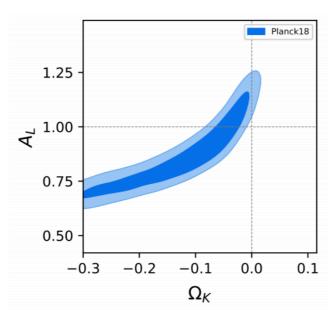


$$C_\ell^{\phi\phi} \to A_L C_\ell^{\phi\phi}$$

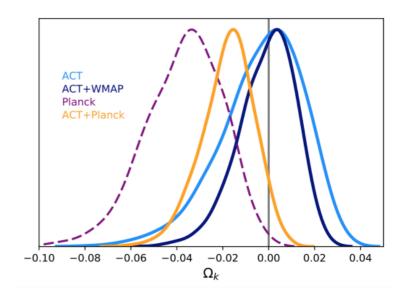
Unless specified, Figs. from Di Valentino++ 1911.02087

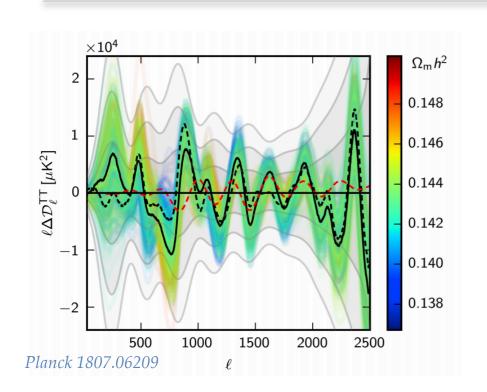


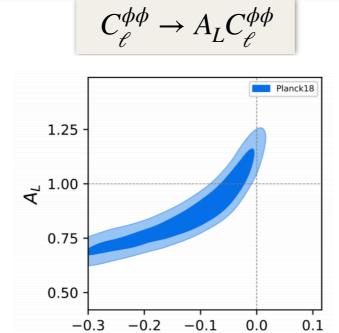




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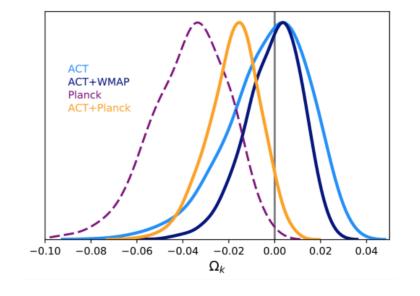


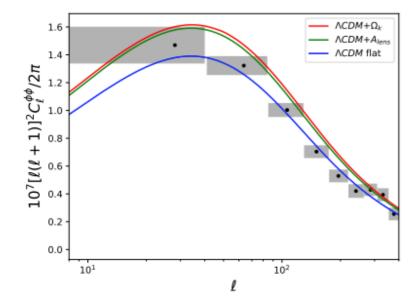


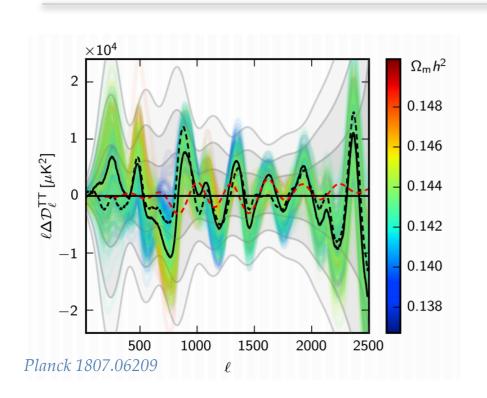


 $\Omega_K$ 

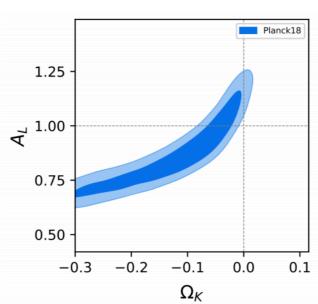




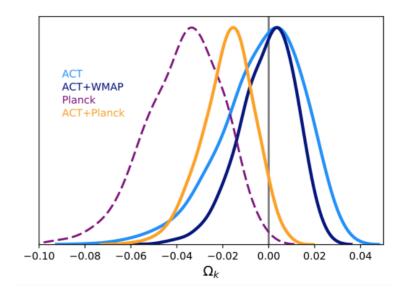


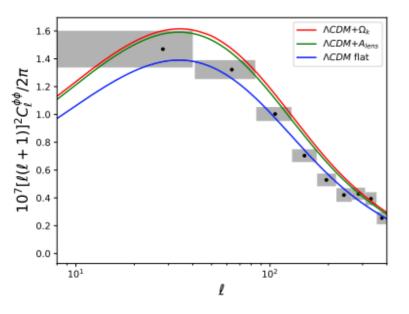


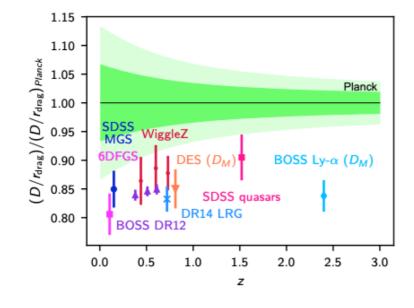


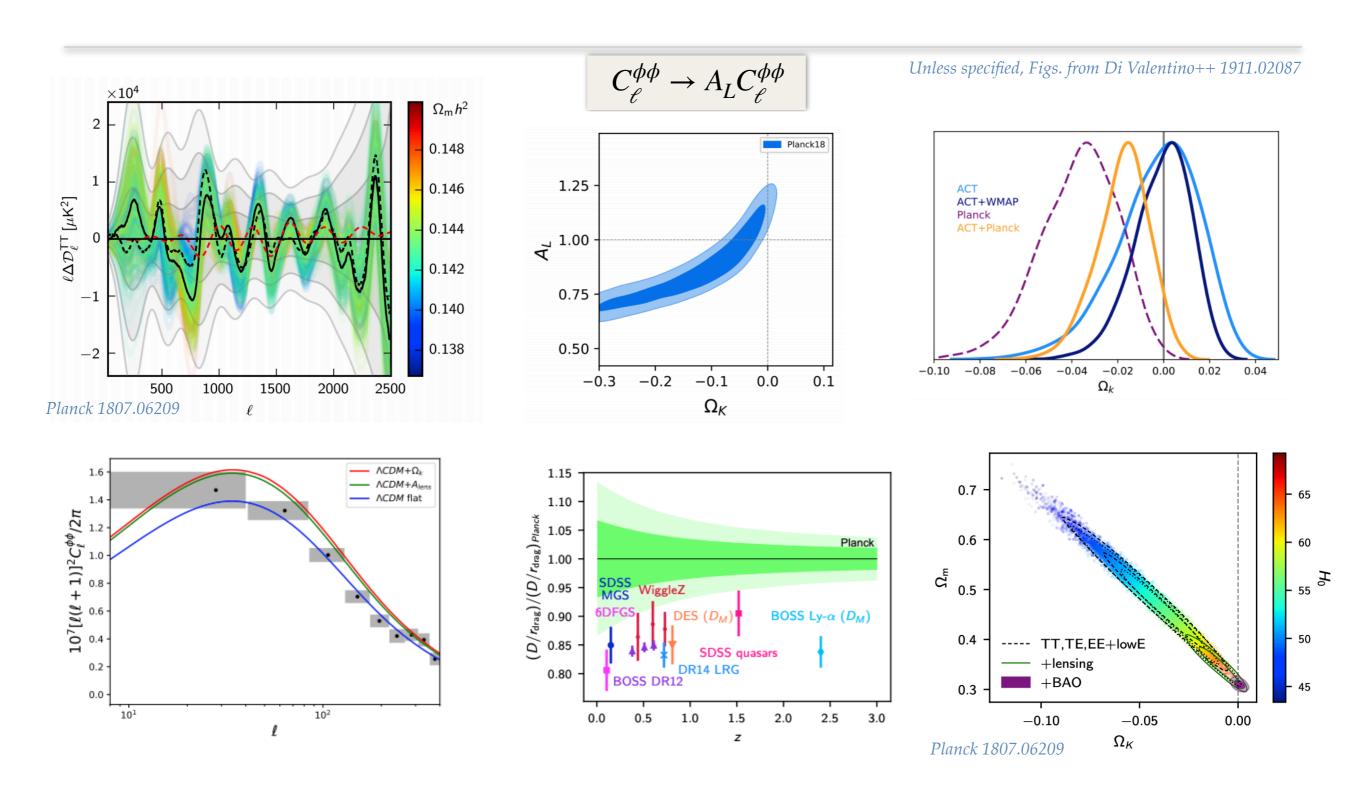


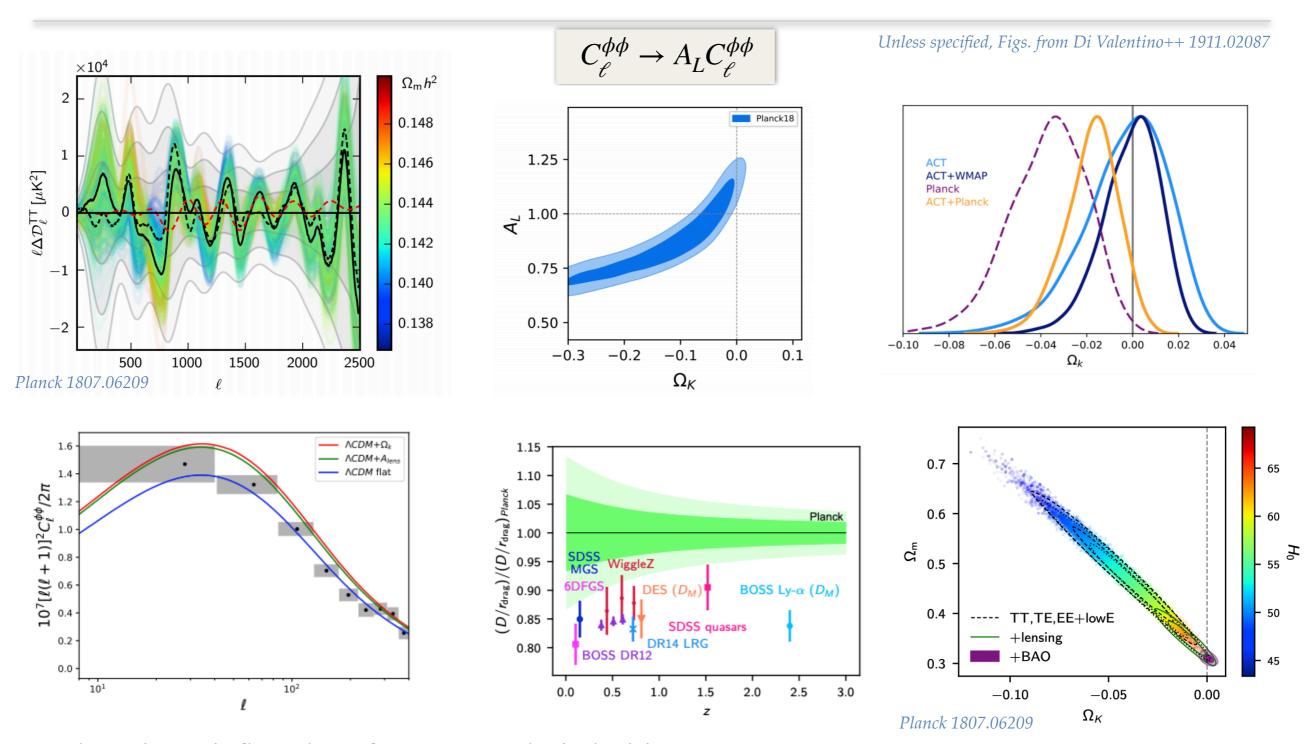
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- The Universe is flat unless of a true 'cosmological crisis'.
- Flat universe is also supported by BOSS and Cosmic Chronometers. *Vagnozzi*++2010.02230, 2011.11645
- Nb:  $A_L$  could also be explained in modified gravity framework, it suffers from the same issues.