

# Colliders in the Sky

Constraining Primordial Non-Gaussianity with CMB and LSS Observations

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## Primordial non-Gaussianity

START: Quantum Fluctuations in Inflaton  $\phi$ 

Single-field, slow-roll inflation, with Bunch-Davies vacuum

**Vanilla inflation**  $\Rightarrow$  **Gaussian** fluctuations in  $\zeta$ 

New physics  $\Rightarrow$  non-Gaussian fluctuations in  $\zeta$ 

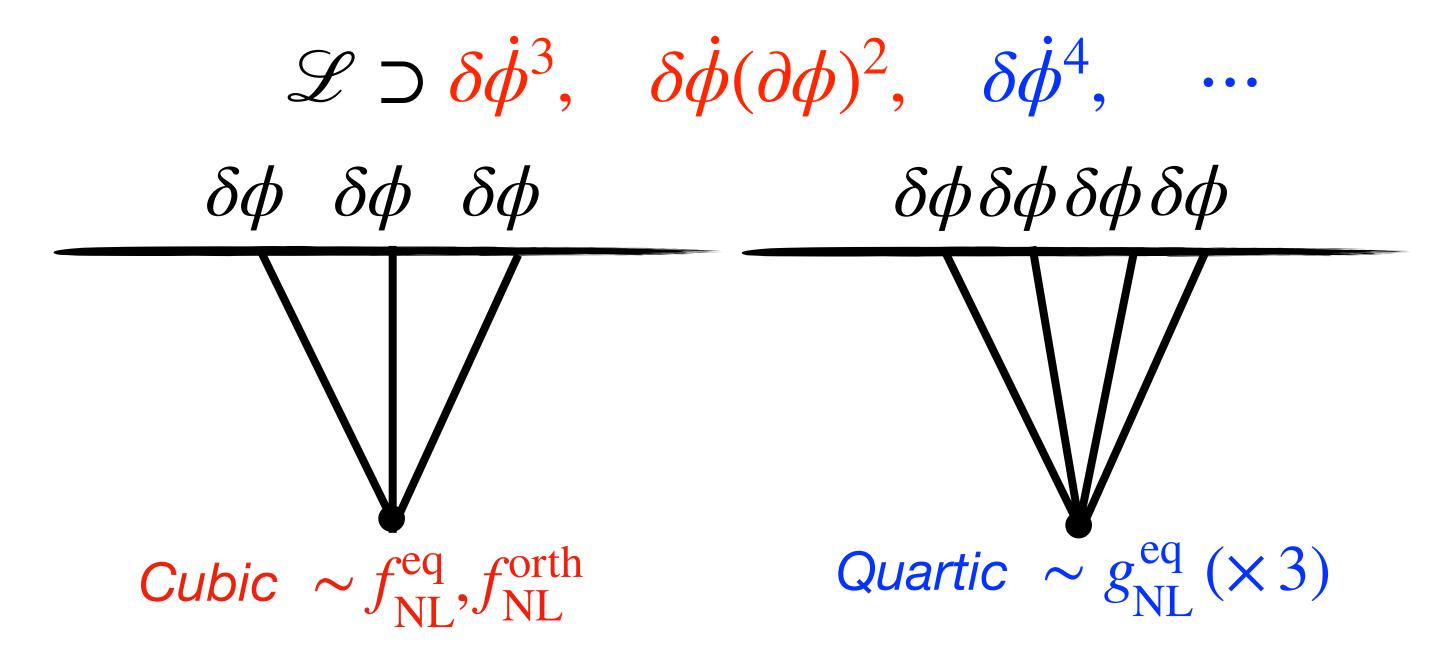
Self-interactions, new fields, new vacuum states, thermal dissipation, ...

END: Classical Fluctuations in Curvature  $\zeta$ 

By searching for non-Gaussianity, we can constrain inflationary physics!

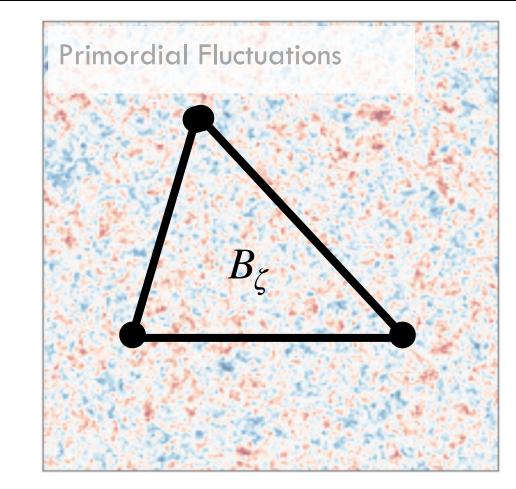
## Self-Interactions

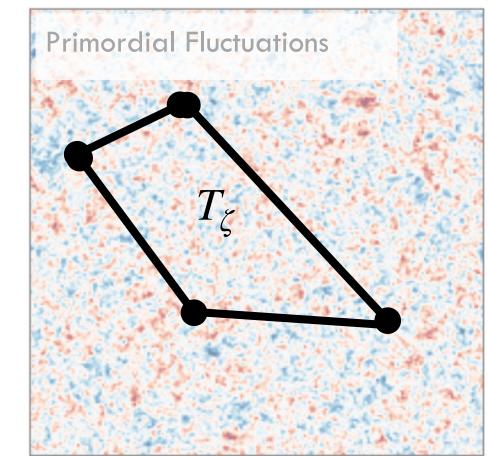
Many models of inflation feature self-interactions:



- These lead to three- and four-point functions at the end of inflation
- The shape encodes the interaction vertex, the amplitude encodes the microphysics

e.g. 
$$\langle \zeta(\mathbf{k}_1)\zeta(\mathbf{k}_2)\zeta(\mathbf{k}_3)\rangle \sim f_{\mathrm{NL}}^{\mathrm{eq}} \times \mathrm{shape}$$





### Self-Interactions

• Other models feature **new particles**,  $\sigma$ :

- Primordial Fluctuations  $B_{\zeta}$
- Primordial Fluctuations  $T_{\zeta}$

- These lead to three- and four-point functions at the end of inflation
- The shape encodes the interaction vertex, the amplitude encodes the microphysics

e.g. 
$$\langle \zeta(\mathbf{k}_1)\zeta(\mathbf{k}_2)\zeta(\mathbf{k}_3)\rangle \sim f_{\rm NL}^{\rm loc} \times {\rm shape}$$

Arkani-Hamed, Maldacena, Lee, Moradinezhad, Cabass, Pajer, Jazayeri, Baumann...

# The Cosmological Collider

- The three- and four-point functions track the exchange of a particle  $\sigma_{\mu_1\cdots\mu_s}$  of mass  $m_\sigma\sim H$  and spin  $s=0,1,2,\cdots$
- In the **collapsed limit** (low exchange momentum), the inflationary signatures are set by **symmetry** and depend **only** on the mass  $m_{\sigma}$ , the spin, s, and the speed  $c_{\sigma}$

$$\langle \zeta^4 \rangle \sim \tau_{\text{NL}} \times \frac{1}{k_1^3 k_3^3 k_{12}^3} \left[ \left( \frac{k_{12}}{k_1 k_3} \right)^{3/2 + i\mu_s} + \left( \frac{k_{12}}{k_1 k_3} \right)^{3/2 - i\mu_s} \right] \mathcal{L}_s(\hat{\mathbf{k}}_1 \cdot \hat{\mathbf{k}}_3)$$

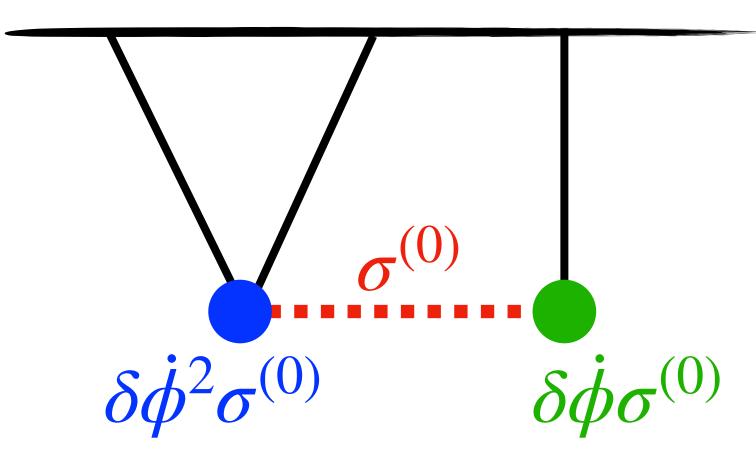
Amplitude  $\sim e^{-m_{\sigma}/H}$ 

Shape (mass dependent) Angle (spin-dependent)

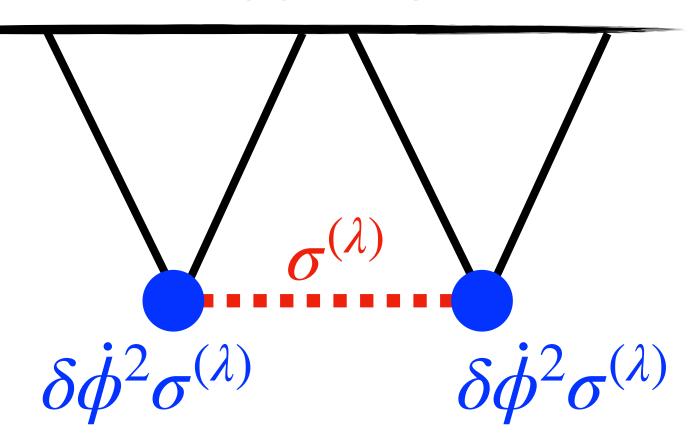
for mass parameter 
$$\mu_{\rm S} = \sqrt{m_{\sigma}^2/H^2 - 9/4}$$

We get oscillations for particles with  $m_{\sigma} \gtrsim H$ 

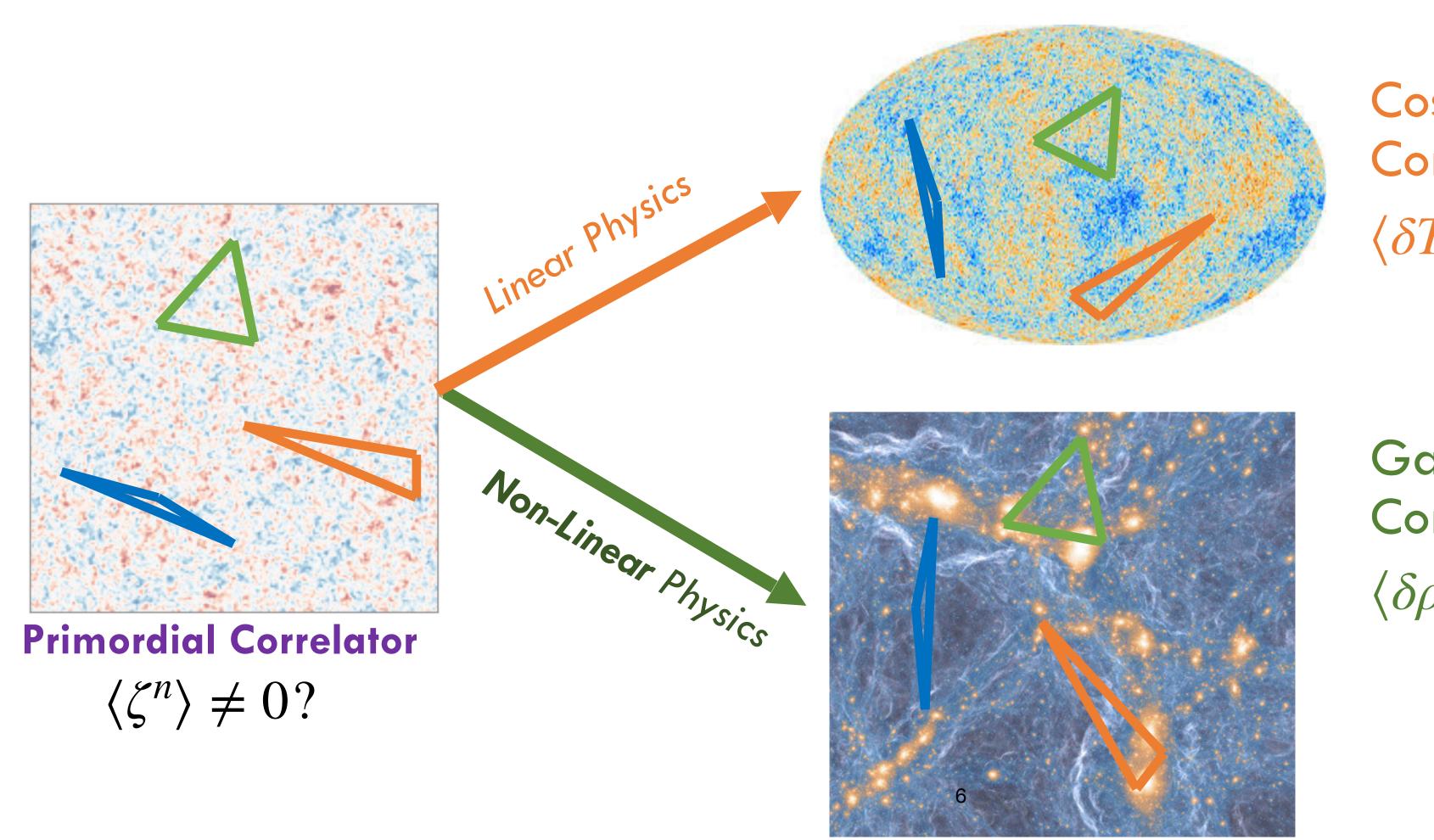
### **Three-Point**



### **Four-Point**



## How to Measure Primordial Non-Gaussianity



Cosmic Microwave Background Correlator

 $\langle \delta T^n \rangle \neq 0$ ?

Galaxy Distribution
Correlator

$$\langle \delta \rho_{\text{galaxy}}^n \rangle \neq 0$$
?

## How to Measure a Three-Point Function

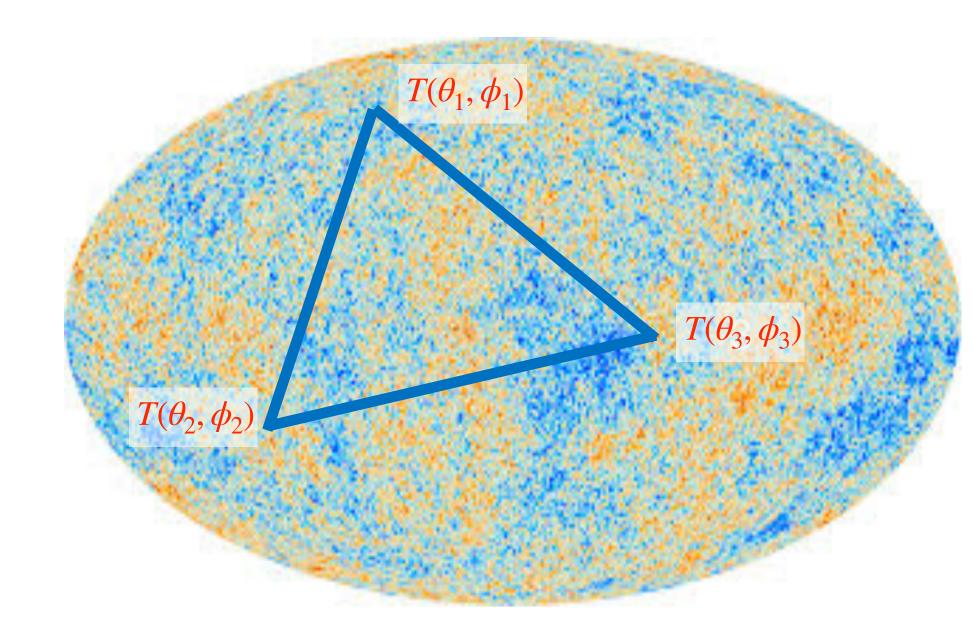
 CMB experiments measure the temperature and polarization across the whole sky

$$T(\theta, \phi), \quad E(\theta, \phi) \quad \leftrightarrow \quad a_{\ell m}^T, \quad a_{\ell m}^E$$

 Since the physics is linear we just need to correlate the CMB at three angles

$$\langle T(\theta_1, \phi_1) T(\theta_2, \phi_2) T(\theta_3, \phi_3) \rangle \leftrightarrow \langle a_{\ell_1 m_1}^T a_{\ell_2 m_2}^T a_{\ell_3 m_3}^T \rangle$$

- This is computationally expensive:
  - The bispectrum is 3-dimensional [after symmetries]
  - There's  $N_{\rm pix}^3 \sim 10^{21}$  combinations of points!



## How to Measure a Three-Point Function

Most CMB analyses use two tricks:

### 1. Compression:

• We compress all  $10^{21}$  elements into a single number, encoding the amplitude of a specific model, e.g.,  $f_{\rm NL}^{\rm loc}$ 

$$\widehat{f_{\mathrm{NL}}^{\mathrm{loc}}} \sim \sum_{\ell_1 m_1 \ell_2 m_2 \ell_3 m_3} \langle a_{\ell_1 m_1} a_{\ell_2 m_2} a_{\ell_3 m_3} \rangle_{\mathrm{theory}}^{\dagger} \times (C^{-1}a)_{\ell_1 m_1} (C^{-1}a)_{\ell_2 m_2} (C^{-1}a)_{\ell_3 m_3}$$

$$\mathsf{Model}$$

$$\mathsf{Data}$$

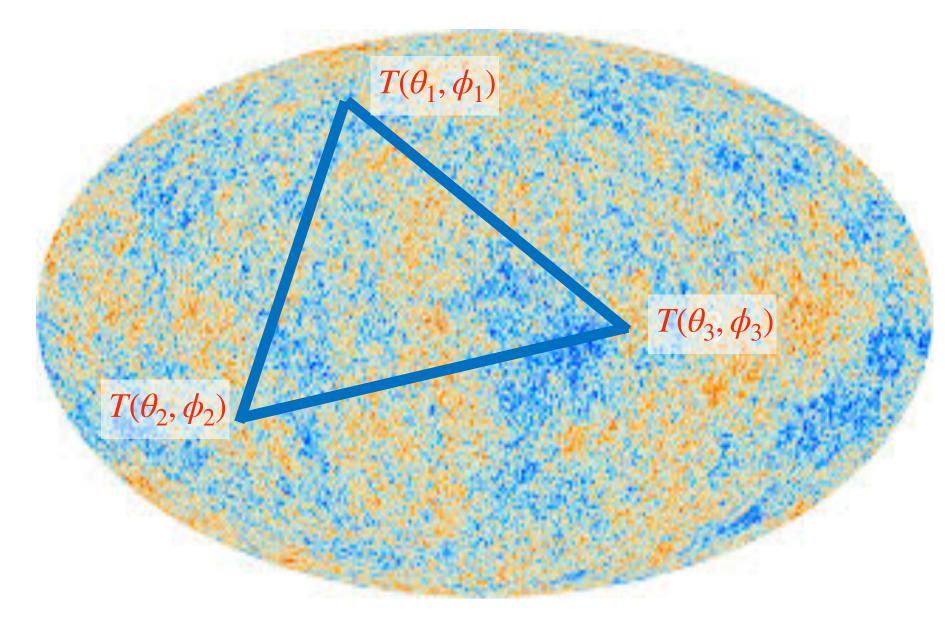
• This is an **optimal estimator** for  $f_{\rm NL}$ , *i.e.* it is **lossless** 

### 2. Separability:

• If the **theory model** is **separable**, we can rewrite the  $\ell$ , m sum using spherical harmonic transforms!

$$B_{\zeta}(k_1, k_2, k_3) \sim \sum_{n} \alpha_n(k_1) \beta_n(k_2) \gamma_n(k_3)$$

- This reduces the complexity from  $\mathcal{O}(N_{\rm pix}^3)$  to  $\mathcal{O}(N_{\rm pix}\log N_{\rm pix})$ 



(**Note**: binned/modal analyses use a similar trick, but compress to a lower-dimensional basis, rather than a single amplitude)

# CMB Bispectrum Constraints

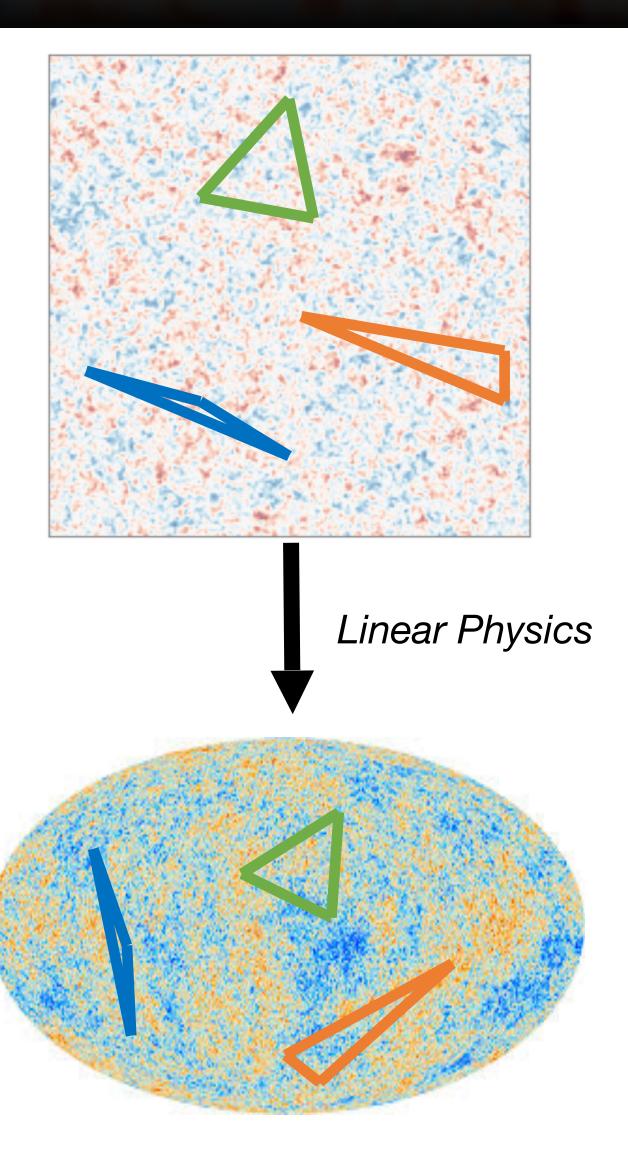
• Planck placed strong constraints on scalar three-point functions, e.g.,

Planck 2018 Local . . . . . . . 
$$-0.9 \pm 5.1$$
 New light scalars  $-26 \pm 47$  Orthogonal . . . .  $-38 \pm 24$  Self-interactions

- These span many phenomenological templates
- Planck uses both separable shapes and modal/binned approximations
- Recent work has also constrained **tensor** three-point functions, e.g.,  $\langle \zeta \zeta h \rangle$  and **cosmological collider** bispectra

**Conclusion**: Scalar primordial non-Gaussianity is **small**:  $10^{-5} \, |f_{\rm NL}| \ll 1$ 

However, we are still far from the (rough) theory targets:  $\sigma(f_{\rm NL}) \sim 1$ 



### What's Next for PNG?

### 1. More models

Folded NG? Excited states? Slow colliders? Strongly-mixed colliders?

### 2. Higher-orders

Four-point functions? Five-point functions? Non-perturbative effects?

### 3. Other datasets

• Next-generation CMB? Galaxy clustering? Weak lensing? 21cm observations?

# The CMB Trispectrum

Very few previous works have considered four-point functions!

Are they worth investigating?

#### Yes!

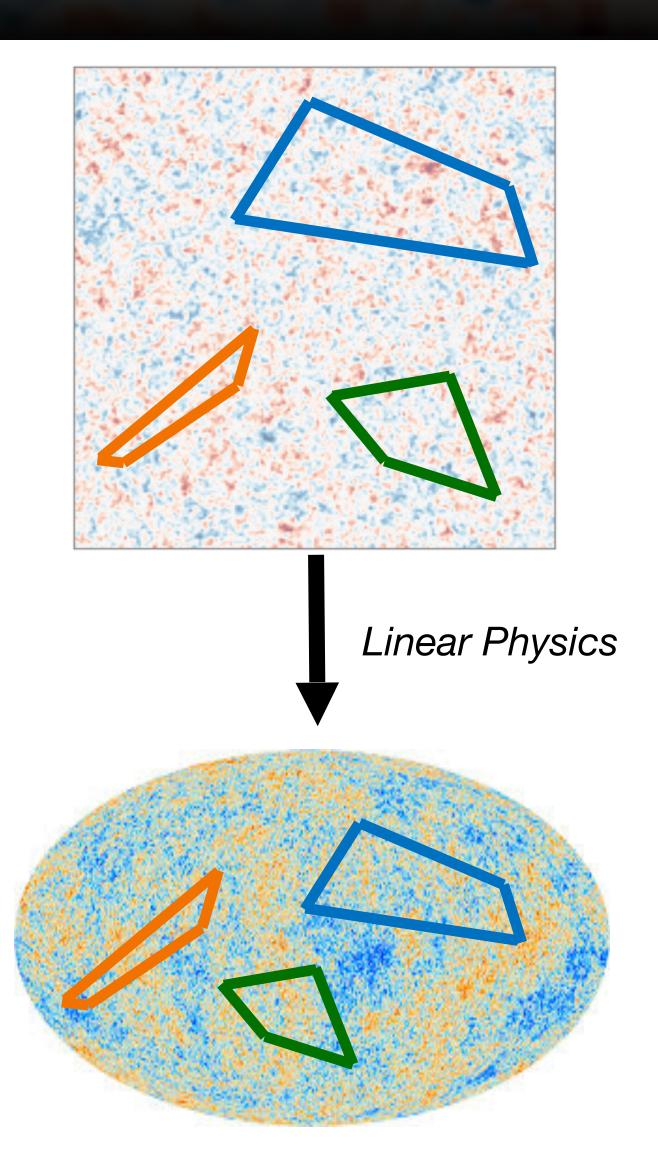
Cubic-terms in the Lagrangian could be protected by symmetry

$$\mathcal{L} \sim \frac{1}{2} (\partial \sigma)^2 + \dot{\sigma}^3 + \dot{\sigma}(\partial \sigma)^2 + \delta \sigma^4 + \cdots$$

(for a general light scalar  $\sigma$ , ignoring coupling amplitudes)

Killed by  $\mathbb{Z}_2$  symmetry ( $\sigma \to -\sigma$ ), or some supersymmetries

- Four-point functions can reveal hidden particle physics, e.g, helicities
- Collider trispectra don't require a linear mixing with the inflaton
- Until recently, we only had constraints on
  - Local effects ( $g_{
    m NL}^{
    m loc}, au_{
    m NL}^{
    m loc}$ )
  - Self-interactions (from the EFT of inflation:  $g_{\mathrm{NL}}^{\mathrm{eq}} \times 3$ )



## How to Measure a Four-Point Function

### Measuring the CMB trispectrum is a challenge!

• The trispectrum is five-dimensional [after symmetries] and depends on  $10^{28}$  sets of points!

$$\langle T(\theta_1, \phi_1) T(\theta_2, \phi_2) T(\theta_3, \phi_3) T(\theta_4, \phi_4) \rangle \leftrightarrow \langle a_{\ell_1 m_1}^T a_{\ell_2 m_2}^T a_{\ell_3 m_3}^T a_{\ell_4 m_4}^T \rangle$$

• We can use **compression** as for the bispectrum:

$$\widehat{g_{\rm NL}} \sim \sum_{\ell_1 m_1 \ell_2 m_2 \ell_3 m_3 \ell_4 m_4} \langle a_{\ell_1 m_1} a_{\ell_2 m_2} a_{\ell_3 m_3} a_{\ell_4 m_4} \rangle_{\rm theory}^{\dagger} \times (C^{-1}a)_{\ell_1 m_1} (C^{-1}a)_{\ell_2 m_2} (C^{-1}a)_{\ell_3 m_3} (C^{-1}a)_{\ell_4 m_4}$$

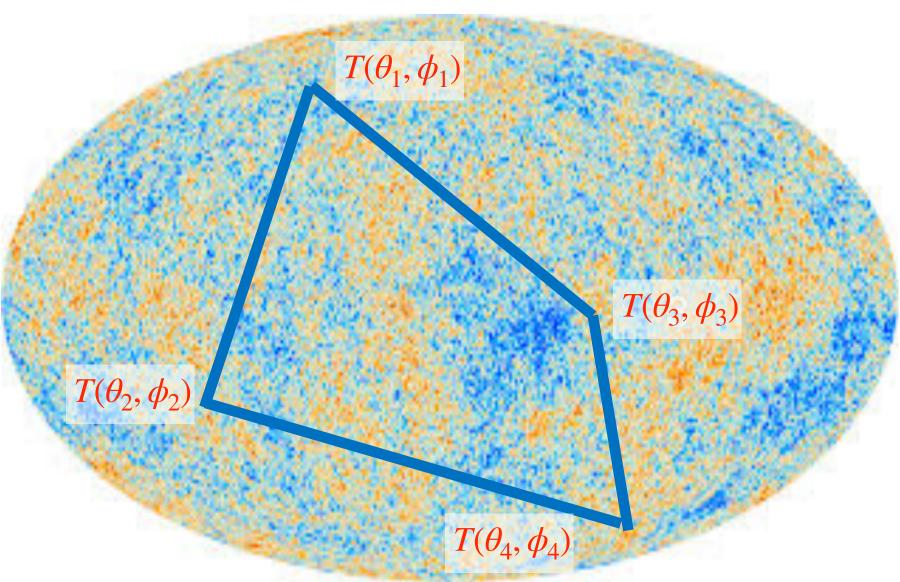
#### Model

#### Data

• To **compute** the  $\ell, m$  sum we use a variety of tricks, including low-dimensional integrals, harmonic transforms, and Monte Carlo summation

$$T_{\zeta}(k_1, k_2, k_3, k_4, s, t, u) \sim F(k_1)G(k_2)H(k_3)I(k_4)J(s^{1/2}) + \cdots$$

• If the trispectrum can be (integral-)**factorized**, this reduces the complexity from  $\mathcal{O}(N_{\rm pix}^4)$  to  $\mathcal{O}(N_{\rm pix}\log N_{\rm pix})$ 



## Optimal Trispectrum Analyses



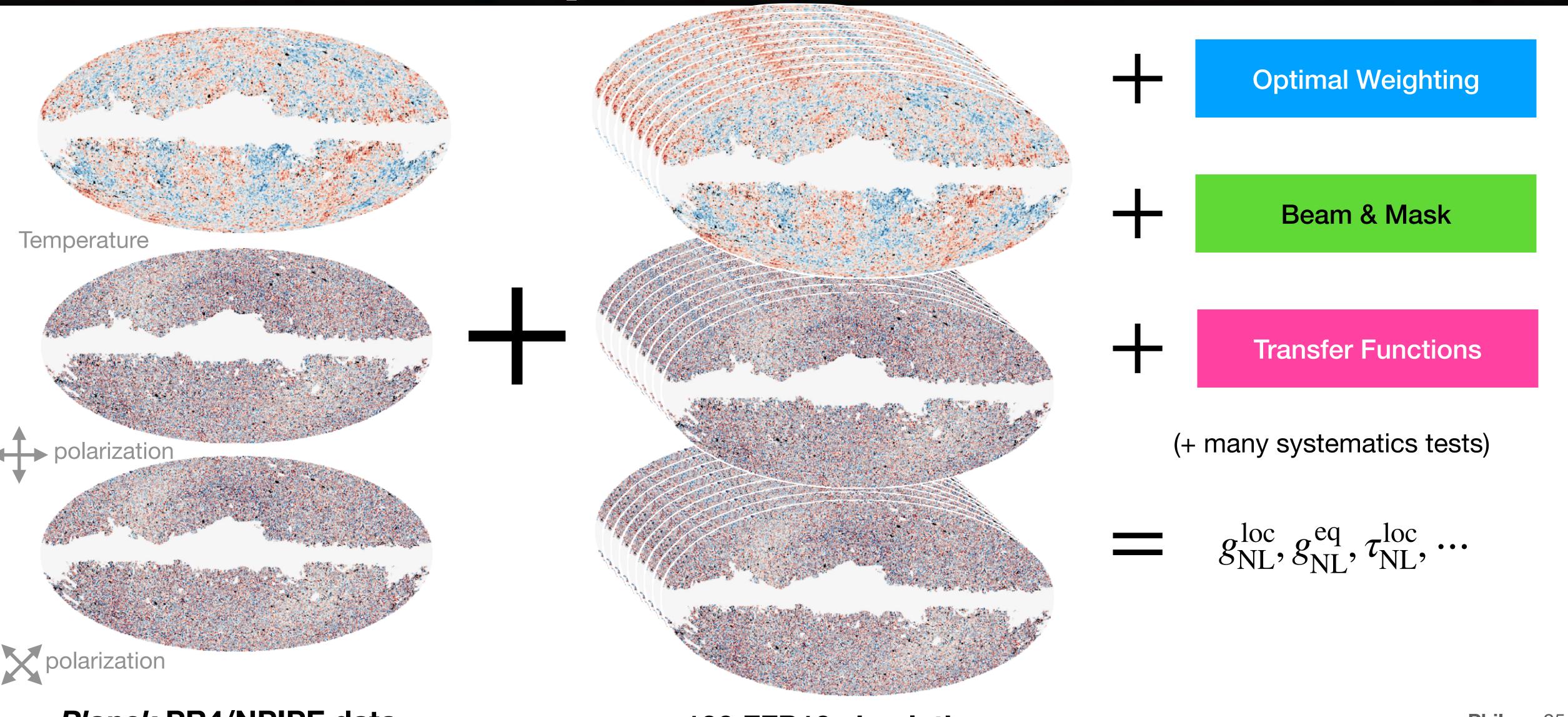
The result: **fast** estimation of four-point amplitudes!

#### The estimators are

- *Unbiased* (by the mask, geometry, beams, lensing, ...)
- Efficient (limited by spherical harmonic transforms)
- Minimum-Variance (they saturate the Cramer-Rao bound)
- Open-Source (entirely written in Python/Cython)
- General (17 classes of factorizable model included so far)

inflation parameters

# The Planck Trispectrum



Planck PR4/NPIPE data

**100 FFP10 simulations** 





### What did we try to detect?

- 1. Cubic local shape  $(g_{NL}^{loc})$
- 2. Quadratic<sup>2</sup> local shape ( $\tau_{\rm NL}^{\rm loc}$ )
- 3. Constant shape  $(g_{NL}^{con})$
- 4. Effective Field Theory of Inflation shapes ( $\times 3$ )
- 5. Direction-dependent shapes
- 6. Cosmological Collider shapes [non-analytic part]
- 7. Weak Gravitational Lensing
- 8. Unresolved **Point-Sources**
- 9. ISW-lensing Trispectra

All of these can be integral-factorized!

## Detecting Non-Gaussianity?



### What did we try to detect?

1. Cubic local shape  $(g_{NI}^{loc})$ 

2. Quadratic<sup>2</sup> local shape  $(\tau_{NI}^{loc})$ 

3. Constant shape  $(g_{NL}^{con})$ 

4. Effective Field Theory of Inflation shapes ( $\times 3$ )

**Direction-dependent** shapes

Cosmological Collider shapes [non-analytic part]

7. Weak Gravitational Lensing

Unresolved **Point-Sources** 

ISW-lensing Trispectra

#### Did we detect it?

No

No

No

No  $(\times 3)$ 

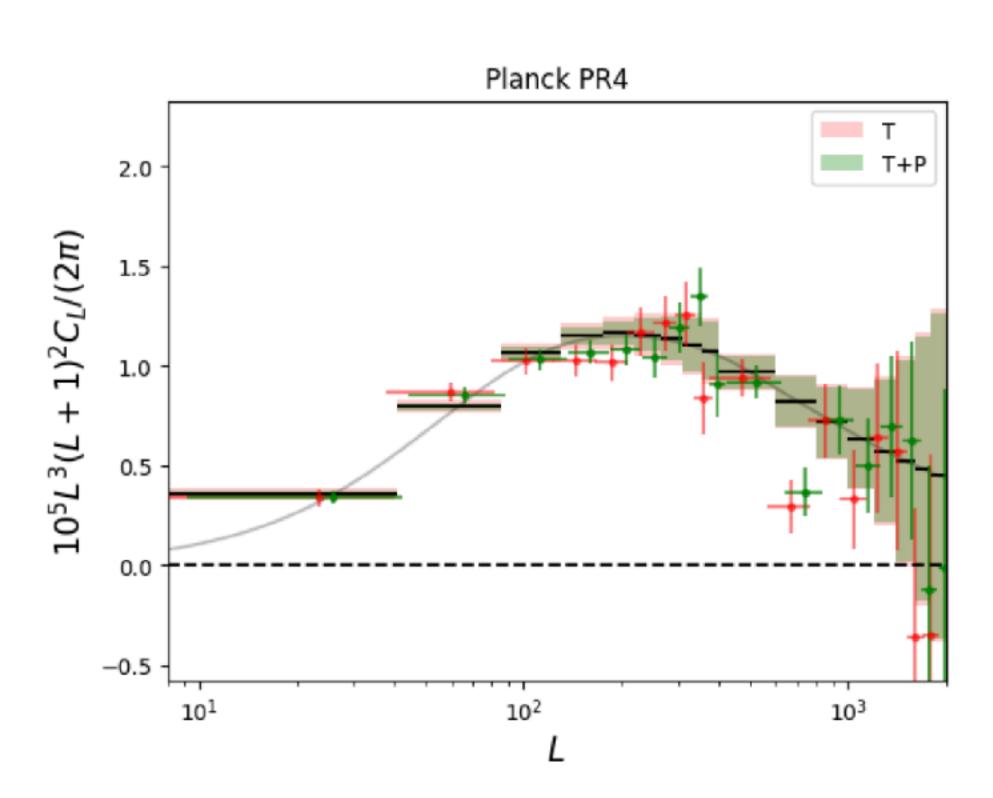
No  $(\times 8)$ 

No  $(\times 17)$ 

Yes!!!

No

No



### Gravitational Lensing $(43\sigma, but not a new detection)$

All of these can be integral-factorized!

## Equilateral Non-Gaussianity

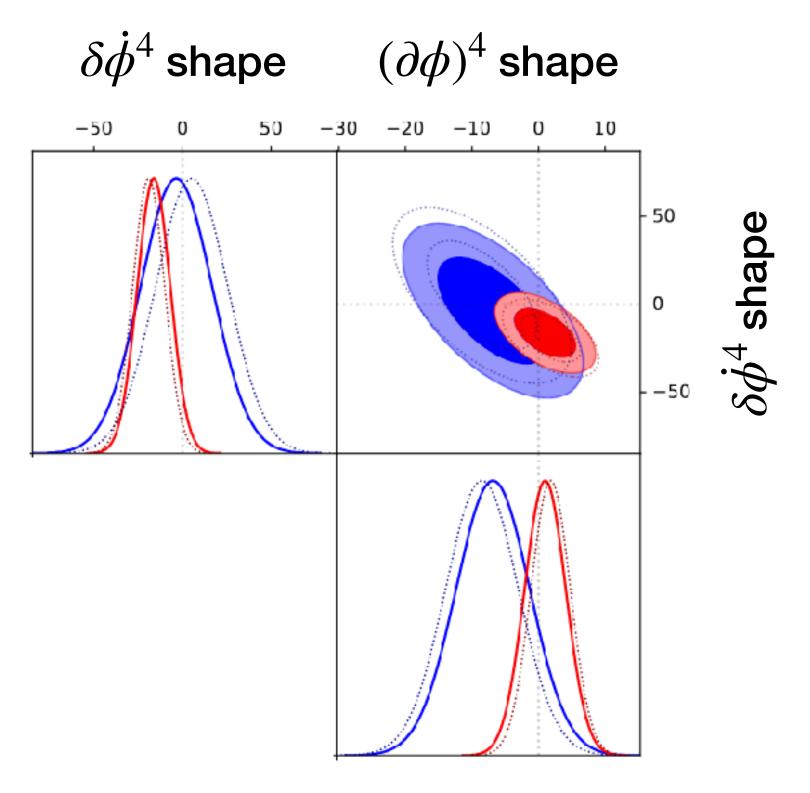
We can constrain cubic self-interactions in inflation

- Constrains models such as:
  - Effective Field Theory couplings
  - **DBI** inflation (string theory + small sound-speed)
  - Generic single-field inflation (including Lorentz Invariant models)
  - Ghost inflation, k-inflation, and beyond...

Outcome: Consistent with zero!

• (50 - 150%) better than any previous constraints!

### T+Pol >>> T-only



The third shape  $-\delta\dot{\phi}^2(\partial\phi)^2$  — is very correlated, so we don't plot it [but we don't detect it]

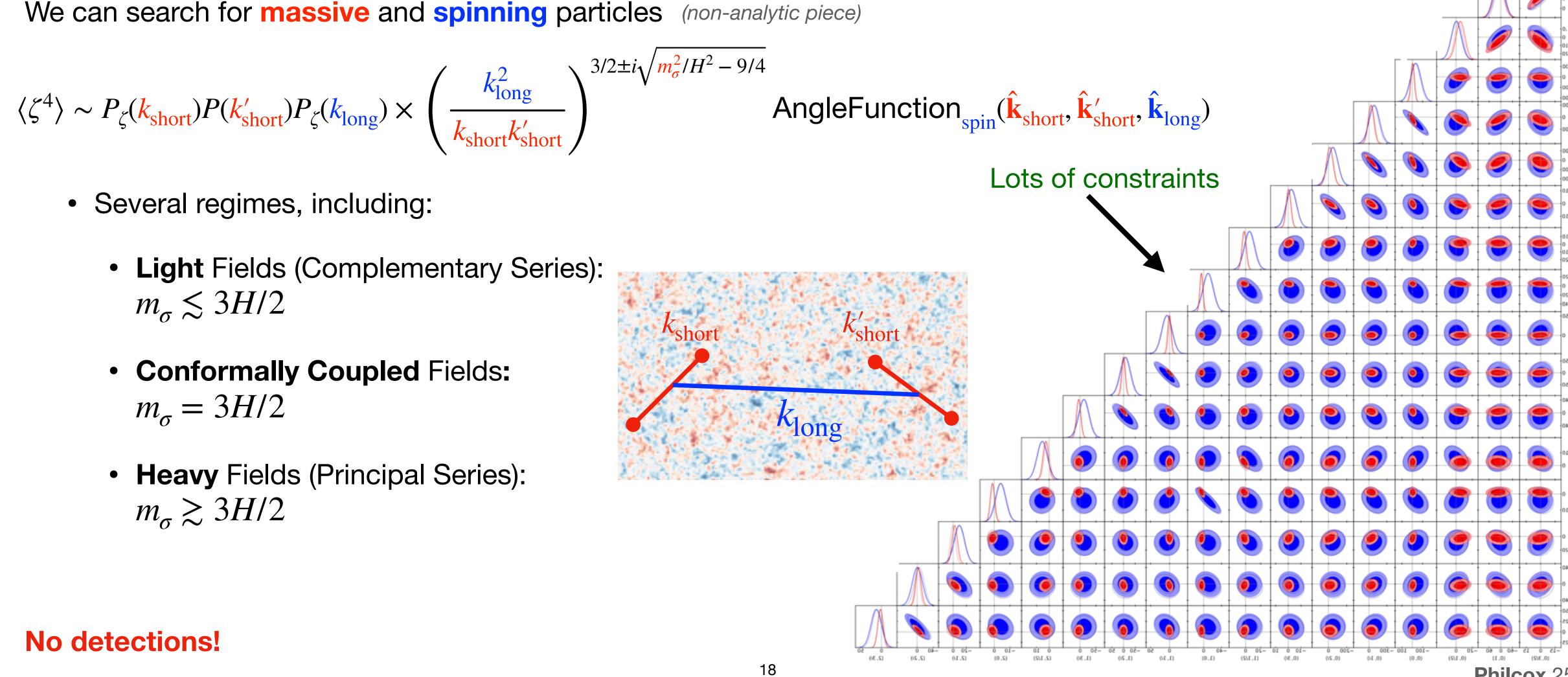
# Cosmological Colliders

We can search for massive and spinning particles (non-analytic piece)

$$\langle \zeta^4 \rangle \sim P_{\zeta}(k_{\text{short}})P(k'_{\text{short}})P_{\zeta}(k_{\text{long}}) \times \left(\frac{k_{\text{long}}^2}{k_{\text{short}}k'_{\text{short}}}\right)^{\frac{1}{2}}$$

• Several regimes, including:

- Light Fields (Complementary Series):  $m_{\sigma} \lesssim 3H/2$
- Conformally Coupled Fields:  $m_{\sigma} = 3H/2$
- Heavy Fields (Principal Series):  $m_{\sigma} \gtrsim 3H/2$

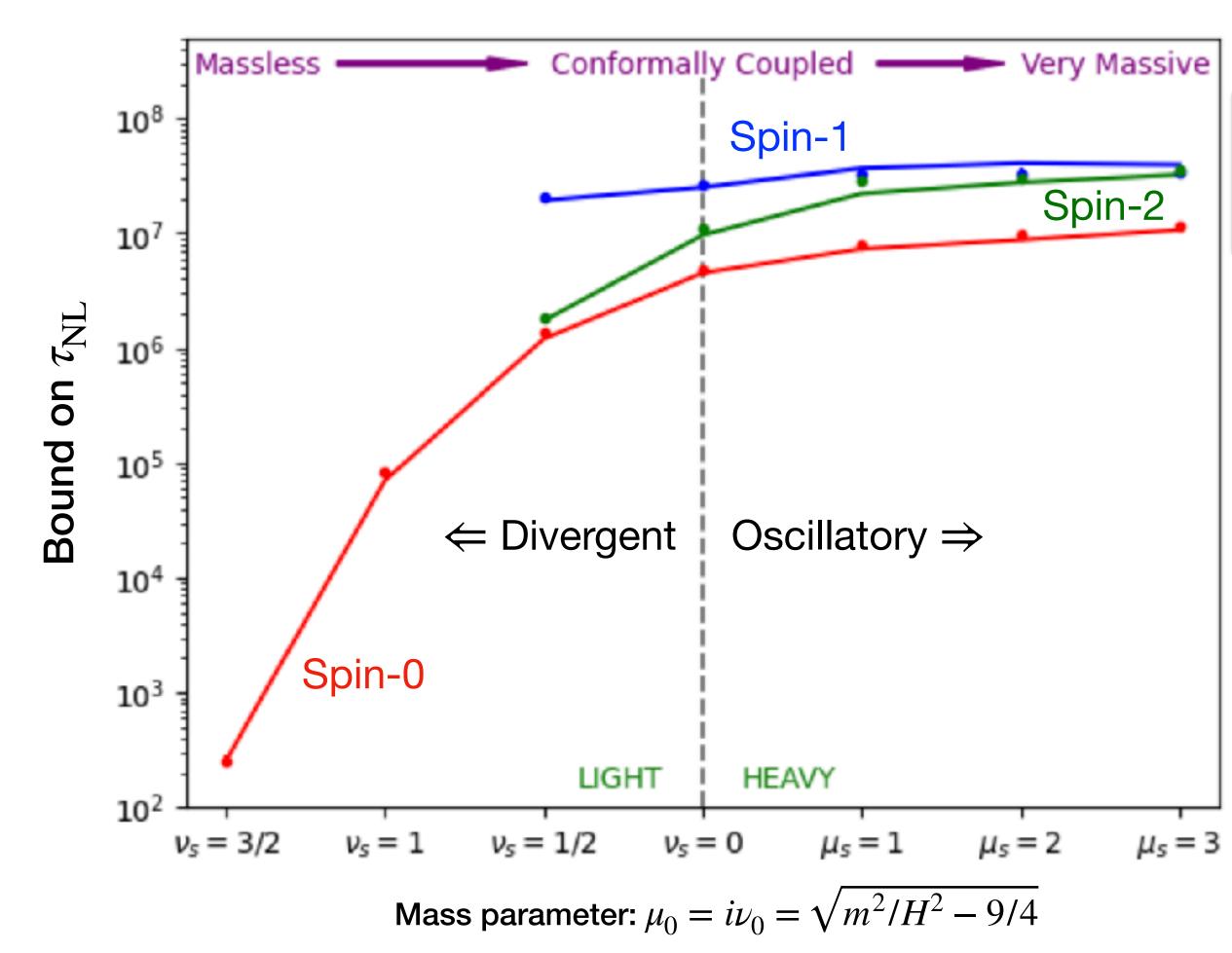


# Cosmological Colliders

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- Several regimes, including:
  - **Light** Fields (Complementary Series):  $m_{\sigma} \lesssim 3H/2$
  - Conformally Coupled Fields:  $m_{\sigma} = 3H/2$
  - Heavy Fields (Principal Series):  $m_{\sigma} \gtrsim 3H/2$

- As expected, light fields are easiest to constrain since their trispectrum diverges
- Odd-spins are hard to constrain due to cancellations!
- **Note**: many of the collider signals are **orthogonal** to the standard templates! [Suman+25, Sohn+24]



## What's Next For the Trispectrum?

There are many ways to extend.

1. More Data

$$\sigma(\tau_{\rm NL}^{\rm loc}) \sim \ell_{\rm max}^{-2}$$

- ACT, SPT, Simons Observatory, LiteBird, CMB-HD, ... will provide data down to much smaller scales!
- Polarization will be particularly useful and could benefit from delensing

#### 2. More Models

- Lighter particles? Heavier particles? Unparticles?
- Tensor non-Gaussianity?
- Collider physics beyond the collapsed limit?
- Thermal baths? Higher-spin particles? Modified sound speeds? Loops? Fermions?
- Scale-dependence? Isocurvature? Primordial magnetic fields?

# Separable Inflationary Correlators

• Efficient bispectrum and trispectrum analyses require factorizable primordial signals.

 $B_{\zeta}(k_1, k_2, k_3) \sim F(k_1)G(k_2)H(k_3)$ 

 $T_{\zeta}(k_1, k_2, k_3, k_4, s, t, u) \sim F(k_1)G(k_2)H(k_3)I(k_4)J(s^{1/2}) + \cdots$ 

- Separable N-point function  $\rightarrow \mathcal{O}(N_{\rm pix} \log N_{\rm pix})$  algorithm
- Non-separable N-point function  $\rightarrow \mathcal{O}(N_{\rm pix}^N)$  algorithm
- Many models of interest are not separable
  - Some require complex oscillatory integrals (via in-in)
  - Others cannot be expressed analytically (e.g., numerical methods)
- To analyze these models, we have two options:
  - 1. Bin the statistic [lossy, and expensive to compute theory predictions!]
  - 2. Create a separable approximation [e.g., modal decompositions]

$$B_{\zeta}(k_1, k_2, k_3) \sim \frac{k_1 k_2 k_3}{(k_1 + k_2 + k_3)^3}$$

# Separable Inflationary Bispectra

• Modal approach: represent the bispectrum as a sum of polynomials:

$$(k_1k_2k_3)^2B_{\zeta}(k_1,k_2,k_3) \sim \sum_{p+q+r=0} \alpha_{pqr}k_1^pk_2^qk_3^r$$
 (or Legendre polynomials)

- Given a target bispectrum, the coefficients  $\alpha_{pqr}$  are computed with linear algebra
- Each term is **factorizable**, so we can build an efficient CMB estimator!
- However, this basis is **big** (  $\approx 5000$  terms used in the *Planck* collider analysis) and does not represent all shapes of interest.
- Alternative approach: learn the basis from the theory itself

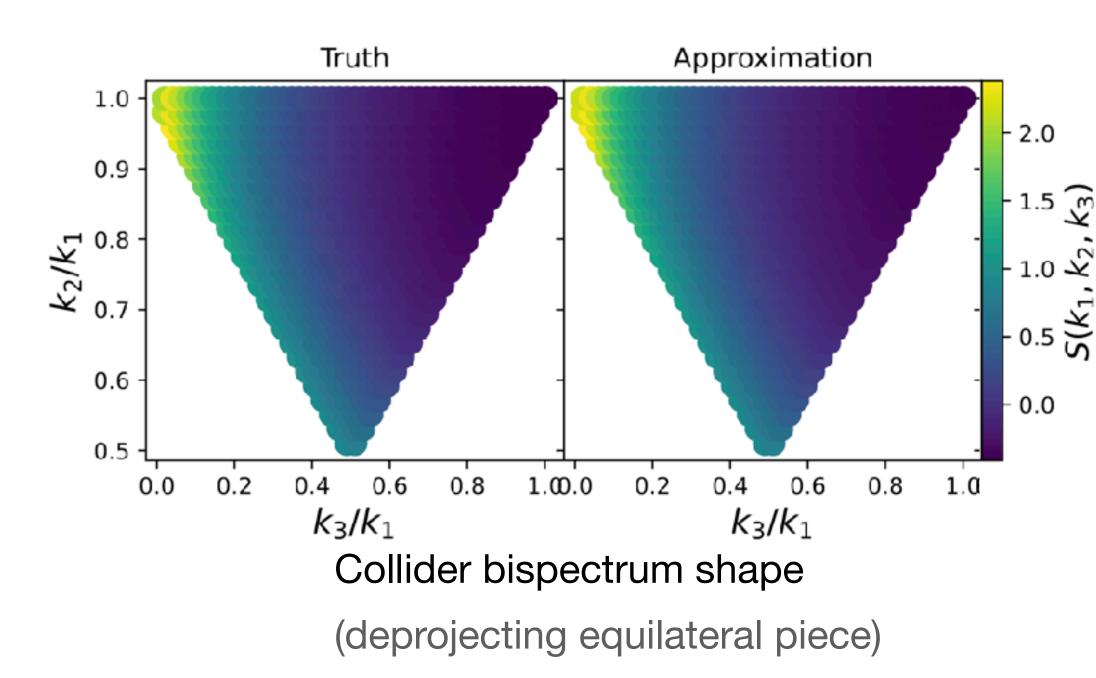
$$(k_1k_2k_3)^2B_{\zeta}(k_1,k_2,k_3) \sim \sum_n w_n\alpha_n(k_1)\beta_n(k_2)\gamma_n(k_3) + \text{perms.}$$

- Given a target bispectrum, the functions  $\alpha_n, \beta_n, \gamma_n$  and weights  $w_n$  are computed using machine learning
- By carefully choosing the loss function, we can optimize the decomposition for the task of interest, e.g., Planck CMB analysis
- This typically requires far fewer terms ( $N \leq 3$ ) to compute the bispectra!

## Separable Bispectra in Practice

- This is implemented in the separable\_bk code, which includes:
  - Simple **neural network** architecture, supplemented with permutation symmetries
  - Training with stochastic gradient descent
  - Fast pytorch implementation, giving basis functions in  $\mathcal{O}(\text{minutes})$
- We test separable\_bk using numerical bispectra obtained with the CosmoFlow code
  - We use a strongly-mixed collider shape that cannot be computed analytically
  - With just three terms, we find approximations with > 99.9% accuracy!

 $(k_1k_2k_3)^2B_{\zeta}(k_1,k_2,k_3) \sim \sum_n w_n\alpha_n(k_1)\beta_n(k_2)\gamma_n(k_3) + \text{perms.}$ 

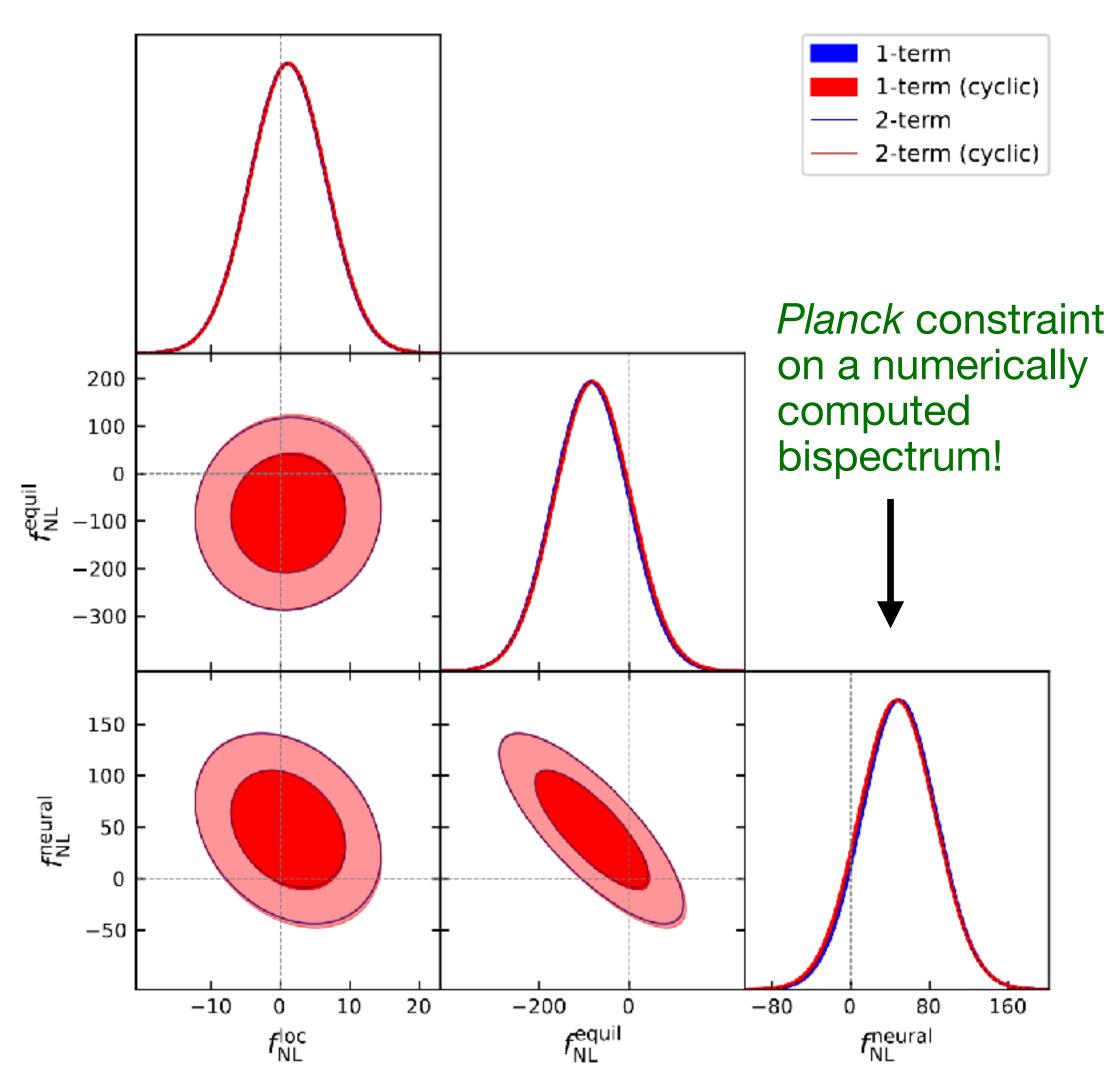


Available at <a href="https://github.com/KunhaoZhong/separable\_bk">https://github.com/KunhaoZhong/separable\_bk</a>

## Separable Bispectra in Practice

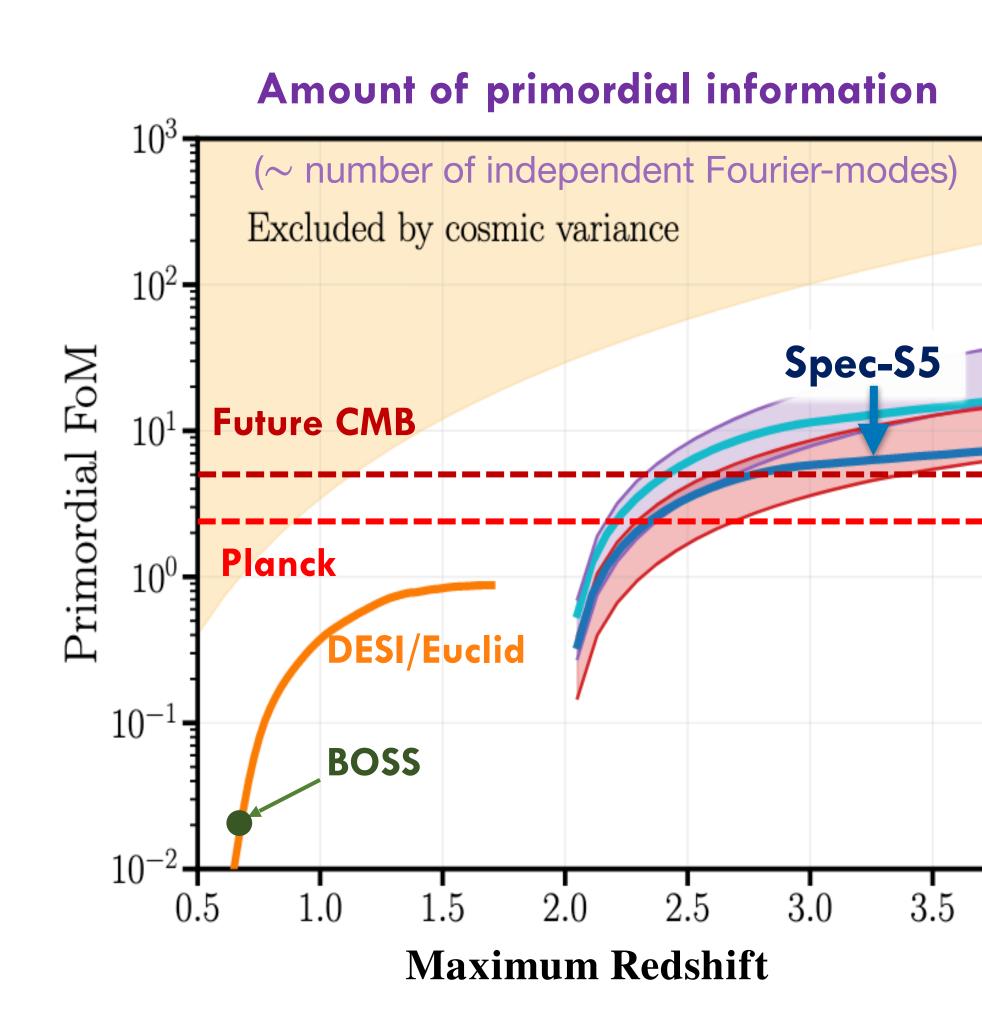
- The output of separable\_bk is a set of **basis functions** describing a particular primordial bispectrum.
- These can be interfaced with the PolySpec code to compute  $f_{\rm NL}$  bounds, e.g., from Planck
- Since the number of separable terms is **small**, we can analyze **collider** bispectra in similar computation time to standard shapes, such as local and equilateral!
- This will allow us to constrain arbitrary primordial bispectra, including those that can only be computed numerically.
- There is **lots** to explore, e.g., analysis of strongly-mixed colliders, and extension to trispectra

Available at <a href="https://github.com/KunhaoZhong/separable\_bk">https://github.com/KunhaoZhong/separable\_bk</a>



## The Future of Non-Gaussianity

- Future CMB experiments will improve bounds on PNG by  $\lesssim 3 \times 10^{-2}$ 
  - This is a two-dimensional field
  - We're running out of modes to look at
    - Large-scales are cosmic-variance-limited
    - Small-scales are limited by secondaries and Silk damping
- What about galaxy surveys?
  - The data precision is rapidly increasing
    - Legacy surveys map a million galaxies [BOSS]
    - New surveys map  $\sim 100 \times$  more! [DESI, Euclid, Rubin, Roman, SphereX,...]
  - This is a three-dimensional field
    - We aren't limited by projection effects
  - There are new observables e.g., galaxy **shapes**, kSZ cross-correlations, ...



# Inflation from Galaxy Surveys

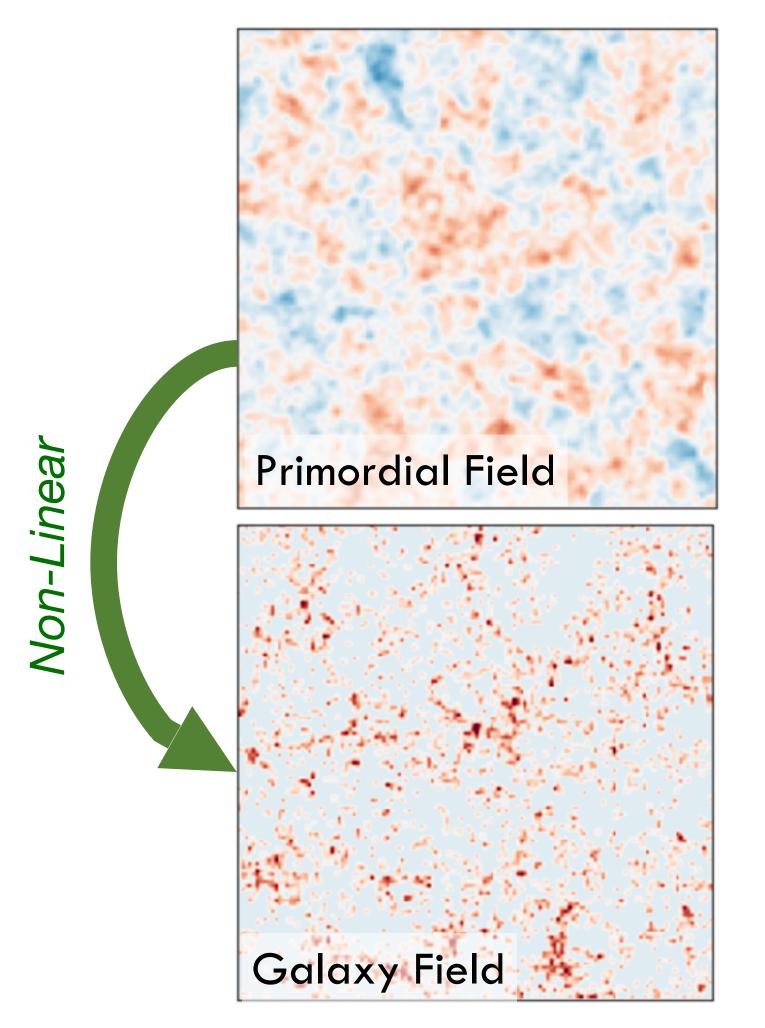
- Modern galaxy surveys map of the distribution of galaxies in three-dimensions:  $\delta_g(\mathbf{x},z)$
- This traces dark matter evolution and the initial conditions

 To extract inflationary information, we need a joint model of all effects:

 $\langle \delta_g \delta_g \delta_g \rangle \sim \text{Primordial Physics} + \text{Gravity} + \text{cross-terms}$ 

State-of-the-art method:

Effective Field Theory of Large Scale Structure (EFTofLSS)

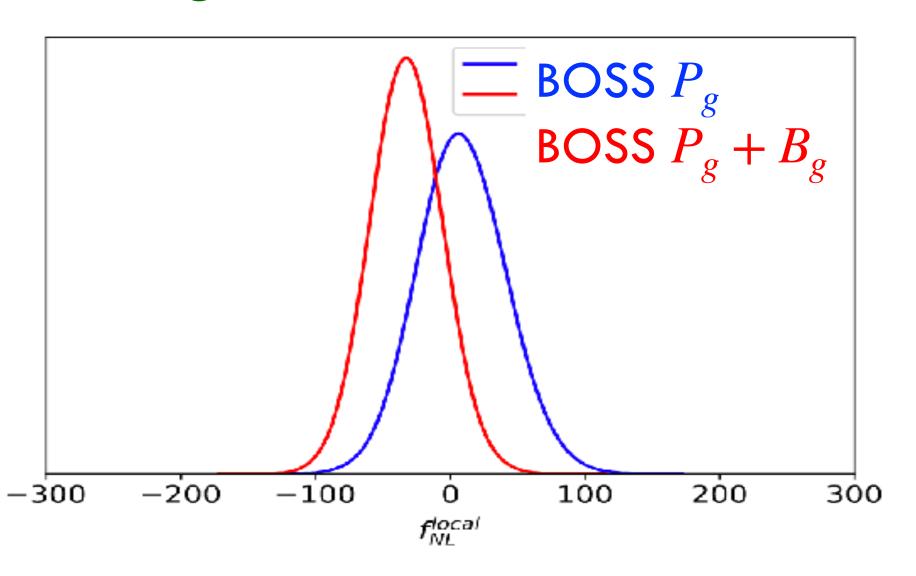


# Inflation from Galaxy Surveys

- Recent works have constrained inflationary bispectra with legacy galaxy survey data (SDSS-BOSS):
  - $f_{
    m NL}^{
    m loc}$ : Local three-point functions from additional light fields
  - $f_{
    m NL}^{
    m eq,orth}$ : **Equilateral** three-point functions from cubic interactions in single-field inflation
  - $f_{
    m NL}^{
    m coll}(m_{\sigma},c_{\sigma})$ : Collider three-point functions from the exchange of massive scalar fields

- For now, the constraints are **much** worse than the CMB  $(5-20\times)$
- Much better data is coming soon!

### Light Field Constraints



$$f_{\rm NL}^{\rm loc} = -33 \pm 28$$
 (9 ± 34 w/o bispectra)

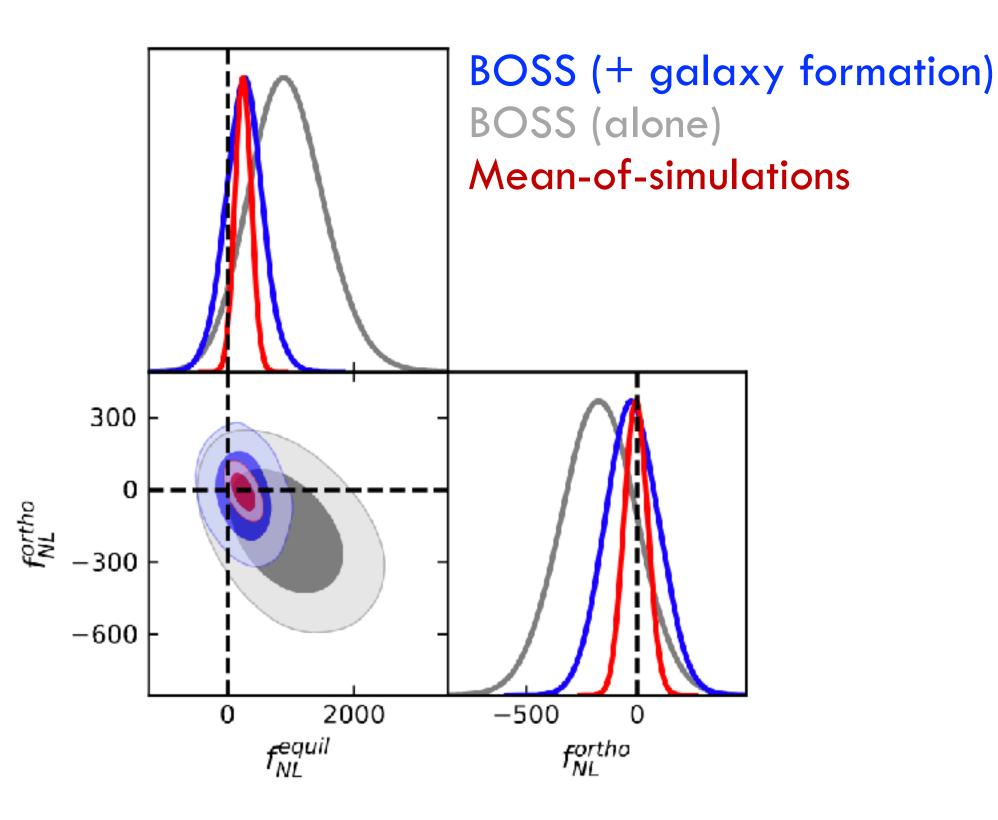
(CMB:  $\pm 5$ , Target:  $\pm 1$ )

# Inflation from Galaxy Surveys

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  - $f_{
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    m loc}$ : Local three-point functions from additional light fields
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### Self-Interaction Constraints



$$f_{\rm NL}^{\rm eq} = 940 \pm 600, f_{\rm NL}^{\rm orth} = -170 \pm 170$$
  
(CMB:  $\pm 50, \pm 25, Target: \pm 1$ )

## Inflation from DESI

### The first year of DESI data is now public!

- We have developed an independent pipeline for analyzing the power spectrum and bispectrum
- This has been used to constrain:  $\Lambda$ CDM ( $\Omega_m, H_0, \sigma_8$ ), dark energy ( $w_0 w_a$ ), curvature ( $\Omega_k$ ), neutrino masses ( $\sum_{n} m_{\nu}$ )

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### New constraints on inflation!

• Multi-field: 
$$f_{\rm NL}^{\rm loc}=0\pm7$$

• Single-Field: 
$$f_{\rm NL}^{\rm eq} = 200 \pm 230, f_{\rm NL}^{\rm orth} = -24 \pm 86$$

(Using the DESI DR1 one-loop power spectrum and bispectrum, plus the high-z, quasar sample)

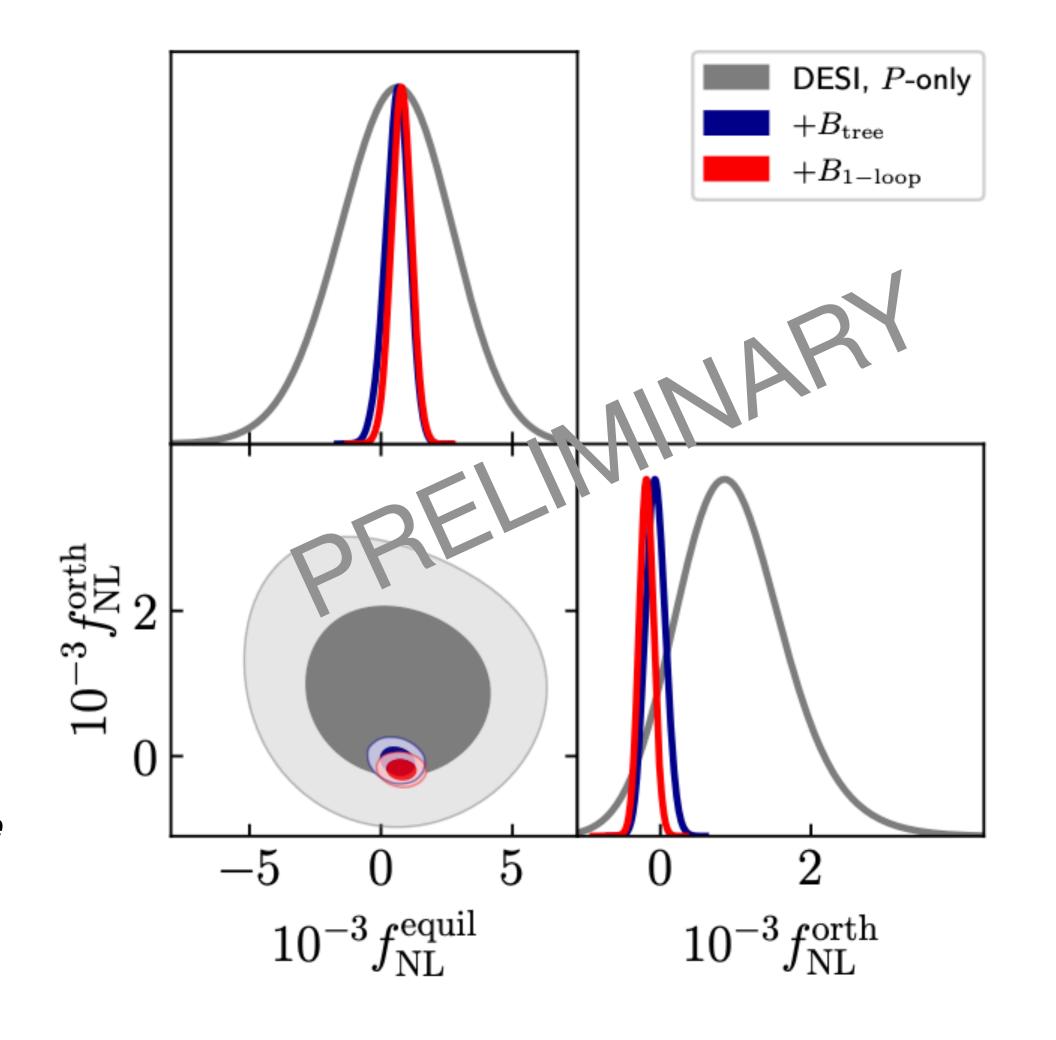
Compare to official DESI

results:

 $f_{\rm NL}^{\rm loc} = -2 \pm 10$ 

• Adding Planck, we obtain the tightest constraint on local PNG yet!!

$$f_{\rm NL}^{\rm loc} = 0 \pm 4$$



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#### **New constraints on inflation!**

• Multi-field: 
$$f_{\rm NL}^{\rm loc} = 0 \pm 7$$

 $f_{\rm NL}^{\rm loc} = -2 \pm 10$ • Single-Field:  $f_{\rm NL}^{\rm eq} = 200 \pm 230, f_{\rm NL}^{\rm orth} = -24 \pm 86$ 

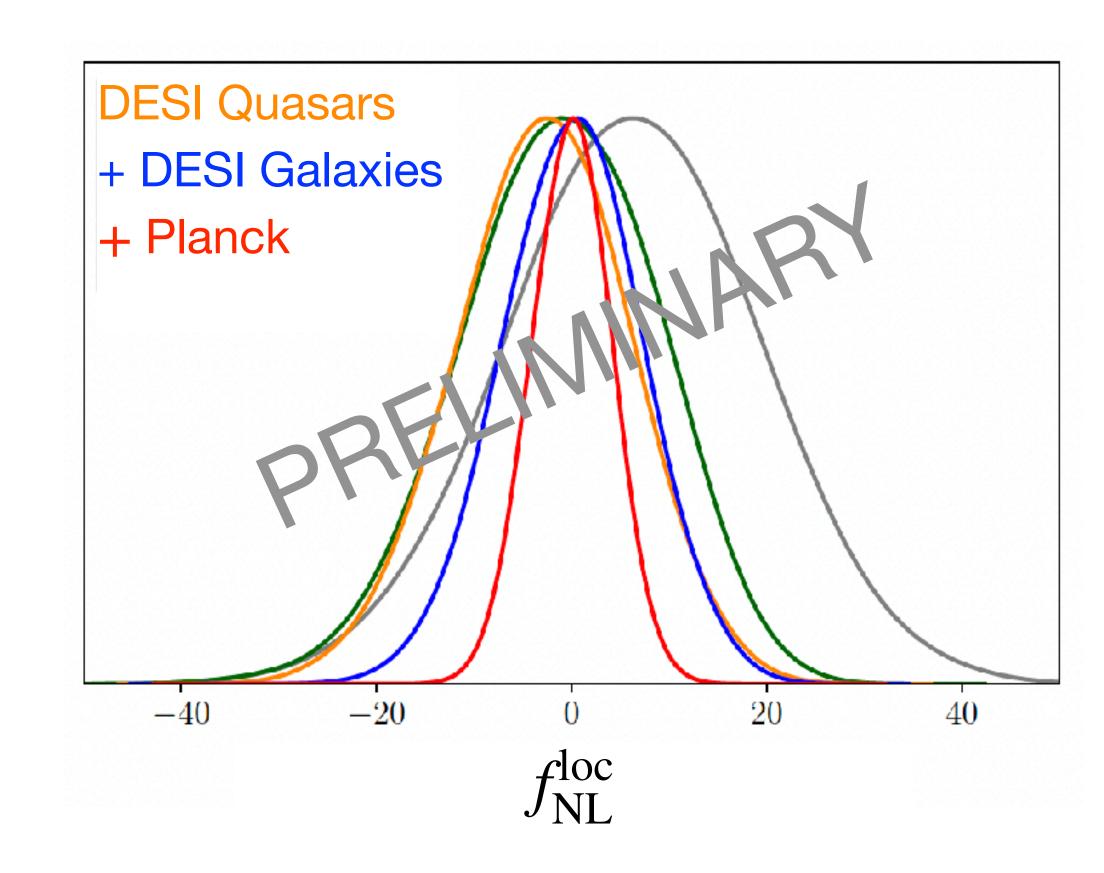
(Using the DESI DR1 one-loop power spectrum and bispectrum, plus the high-z, quasar sample)

Compare to official DESI

results:

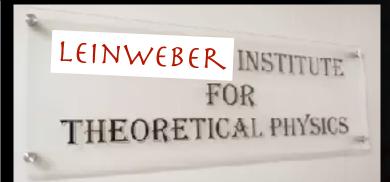
Adding Planck, we obtain the tightest constraint on local PNG yet!!

$$f_{\rm NL}^{\rm loc} = 0 \pm 4$$









# Summary

PNG analysis is a very active field!

- We can now constrain inflationary four-point functions in the CMB, including the cosmological collider!
- We can probe arbitrary non-separable bispectrum models with CMB data and machine-learning
- Galaxy surveys are providing exciting new insights into inflation and are starting to rival the CMB

