

How to obtain slow roll inflation driven by non linear electrodynamics

- Do we need a scalar field for slow roll?
- Can we use gravity plus electromagnetism as a alternative mechanism to inflation?
- We derive the general Lagrangian for NLED satisfying conditions for slow roll.
- With D. Malafarina and H. Chakrabarty
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Ilia Musco

University of Nova Gorica (Slovenia)



- **Action:**

$$S = \int (16\pi R - \mathcal{L}(F)) \sqrt{-g} d^4x, \quad F = F_{\mu\nu} F^{\mu\nu}$$

$$\mathcal{L} = -F/4, \quad F = 2(B^2 - E^2) \quad \text{Maxwell limit}$$

- **Stress-energy tensor:**

$$T_{\alpha\beta} = -4(\partial_F \mathcal{L}) F_{\rho\alpha} F_{\beta}^{\rho} + g_{\alpha\beta} \mathcal{L} = (\rho + p) u_{\alpha} u_{\beta} + p g_{\alpha\beta}$$

$$\rho = -4\partial_F \mathcal{L} E^2 - \mathcal{L} \quad p = -4\partial_F \mathcal{L} \left(\frac{2B^2 - E^2}{3} \right) + \mathcal{L}$$

- **Magnetic Universe:** in the radiation dominated Universe the electromagnetic field interacts strongly with the plasma \Rightarrow the electric field is screened in the rest frame of the fluid ($E^2 = 0$).

- **Equation of state:** $\mathcal{A}(F) := \log \mathcal{L}(F)$

$$p = w\rho \Rightarrow w = -1 + \frac{8B_0^2}{3a^4} \partial_F \mathcal{A} = -1 + \frac{4}{3} F(\partial_F \mathcal{A})$$

- **Accelerated expansion:** $\ddot{a} > 0, \quad \rho(1 + 3\omega) < 0 \Rightarrow \frac{4B_0^2}{a^4} \partial_F \mathcal{A} < 1$

- **Slow roll parameters:**

$$\epsilon \equiv -\frac{\dot{H}}{H^2} = \frac{3}{2}(1 + w) = \frac{4B_0^2}{a^4} \partial_F \mathcal{A} \quad \eta \equiv \frac{\dot{\epsilon}}{\epsilon H} = \frac{w_{,a} a}{1 + w} = \frac{(\partial_F \mathcal{A})_{,a}}{\partial_F \mathcal{A}} a - 4$$

- **De Sitter initial state:** $(\epsilon, \eta) \rightarrow 0 \Rightarrow F(\partial_F \mathcal{A}) \rightarrow 0, \quad \frac{F(\partial_{FF} \mathcal{A})}{\partial_F \mathcal{A}} \rightarrow -1$

- **Lagrangian NLED:**

$$\mathcal{L} = -\alpha \log(1 + \beta F + h(F))$$