# Inflationary Trispectrum of Gauge Fields from Scalar and Tensor Exchanges

(arXiv:2509.11143)



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Inflation 2025

03 December 2025

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- Using the ADM decomposition to next order one can obtain with scalar fluctuations

$$H_{\zeta AA} = n \int d^3x \, \lambda \zeta \left( A_i^{'2} - \frac{1}{2} F_{ij}^2 \right) + O(\epsilon) \quad \text{R.K. Jain and M.S. Sloth, Phys. Rev. D 86 (2012) 123528}_{L. \, Motta \, and \, R.R. \, Caldwell, \, Phys. \, Rev. \, D \, 85 \, (2012) \, 103532}$$

# In-in: Cosmological Diagrams

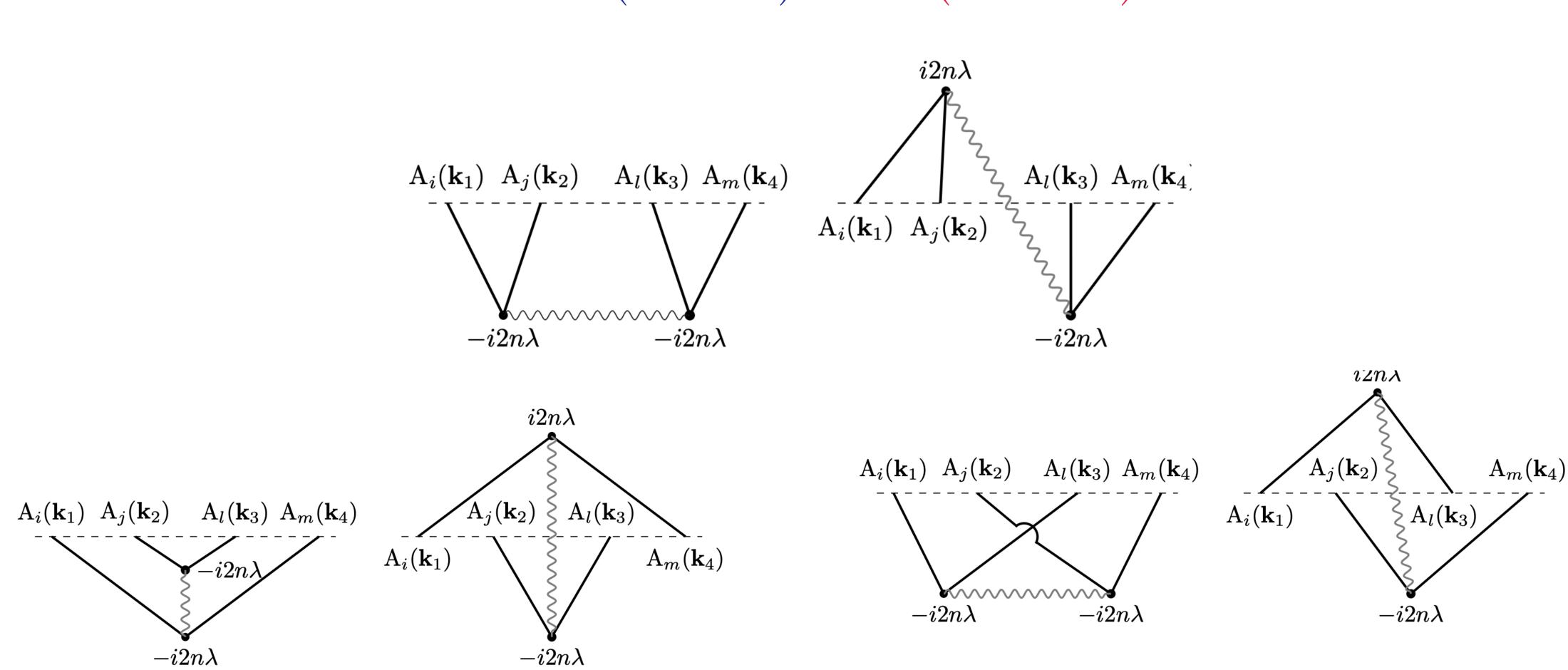
In-in master formula

$$\left\langle \Omega \left| \mathcal{O}(\eta_0) \right| \Omega \right\rangle = \left\langle 0 \left| \tilde{T} \left( e^{i \int_{-\infty}^{\eta_0} d\eta H_{\text{int}}} \right) \mathcal{O}(\eta_0) T \left( e^{-i \int_{-\infty}^{\eta_0} d\eta H_{\text{int}}} \right) \right| 0 \right\rangle$$

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## Scalar Exchange: Magnetic Trispectra

We focus on the index free trispectra

$$\langle \boldsymbol{B}(\boldsymbol{k}_1) \cdot \boldsymbol{B}(\boldsymbol{k}_2) \; \boldsymbol{B}(\boldsymbol{k}_3) \cdot \boldsymbol{B}(\boldsymbol{k}_4) \rangle = (2\pi)^3 \, \delta^{(3)} \left( \sum_{\alpha=1}^4 \mathbf{k}_{\alpha} \right) \mathcal{T}_{\zeta B}(\boldsymbol{k}_1, \boldsymbol{k}_2, \boldsymbol{k}_3, \boldsymbol{k}_4) \qquad \text{With} \quad \mathcal{T}_{\zeta B} = \frac{1}{a^8} \left( \frac{\dot{\lambda}}{H \lambda} \right)^2 \sum_{S \in 1, 2, 3} \sum_{X, Y \in A, A'} \mathbf{Q}_{SXY}^{\zeta} \mathbf{I}^{SXY}$$

Polarisation function is

$$Q_{1AA}^{\zeta} = k_1^2 k_2^2 k_3^2 k_4^2 \left( 1 + \cos^2 \theta_{12} \right) \left( 1 + \cos^2 \theta_{34} \right)$$

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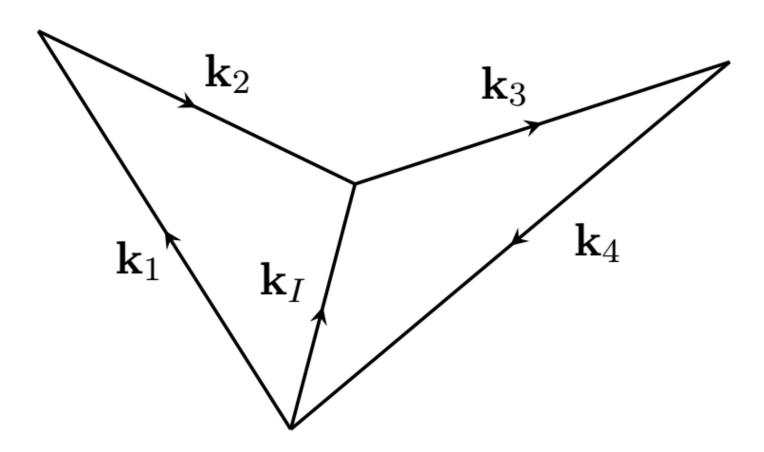
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#### Momentum Quadrilateral

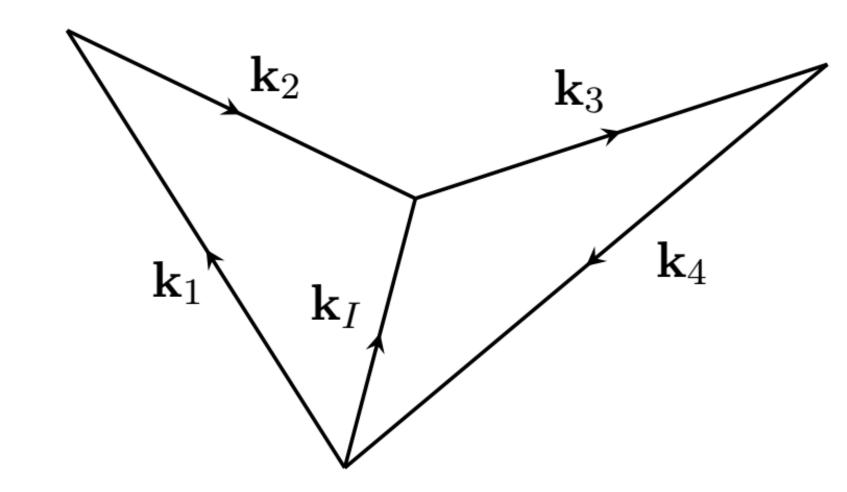
► Momentum conservation ⇒ closed shapes



Shape is not planar

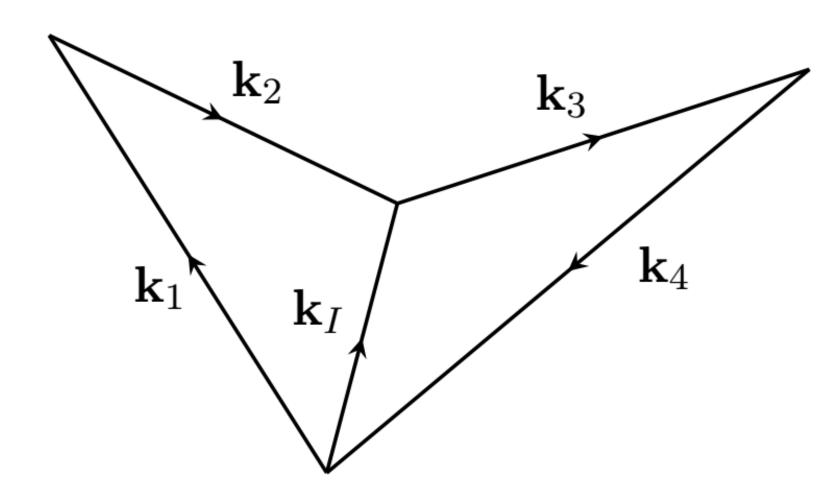
# Equisided Configurations

- Equisided configuration:  $|\mathbf{k}_1| = |\mathbf{k}_2| = |\mathbf{k}_3| = |\mathbf{k}_4|$
- Removes the directional hierarchies
- Shapes can be analysed using two characteristic scale



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Contains all the informations of three channels and the nested time integrals

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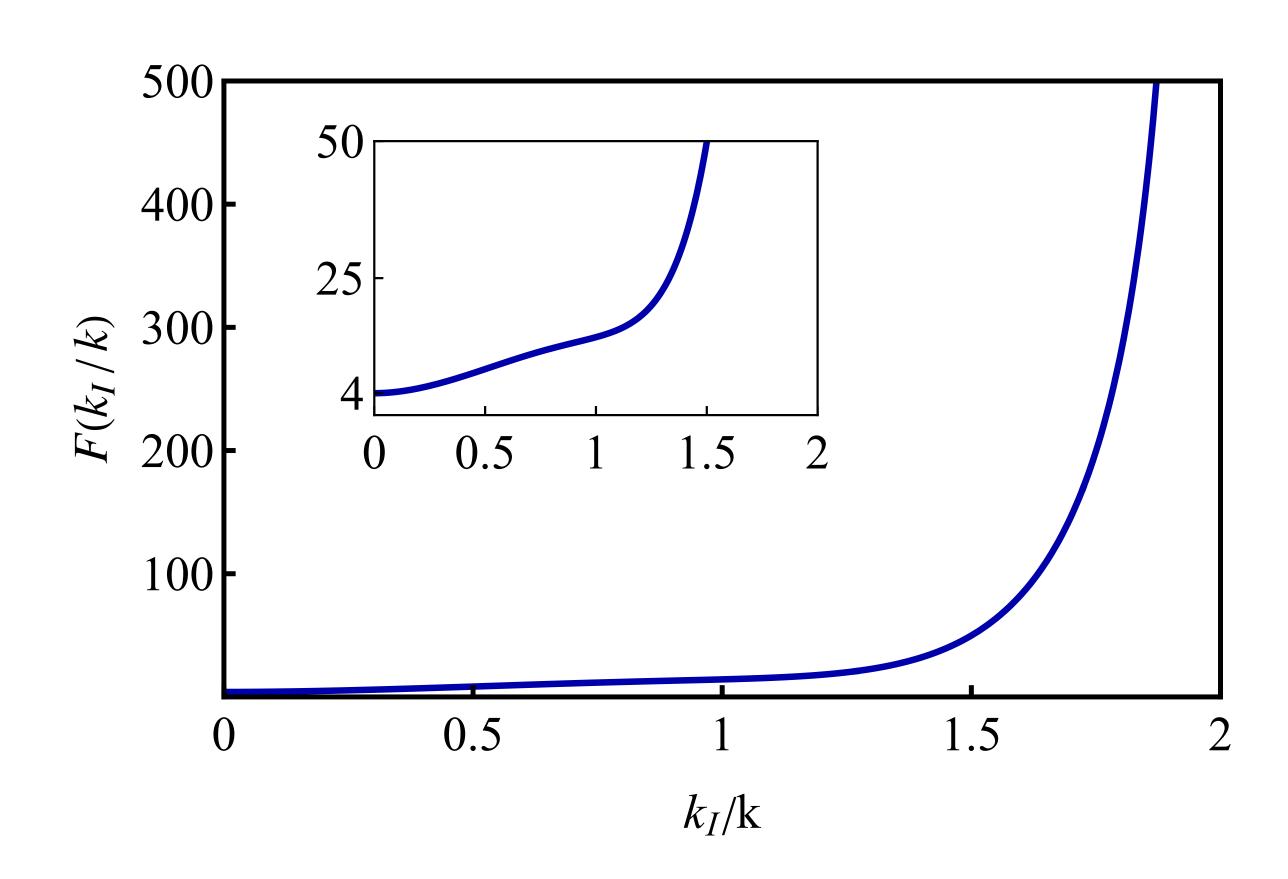
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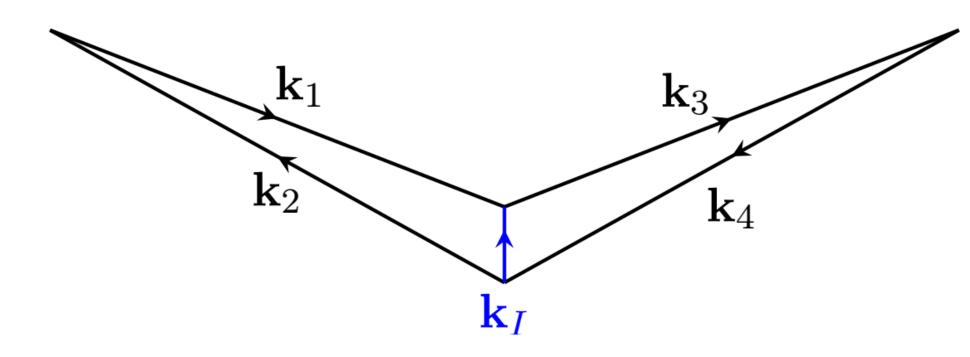
• For  $\phi = 0$ 

- Monotonically increasing function of  $k_I/k$
- Correlation strength  $\propto k_I$
- $F(0) \neq 0$ , is the counter collinear limit



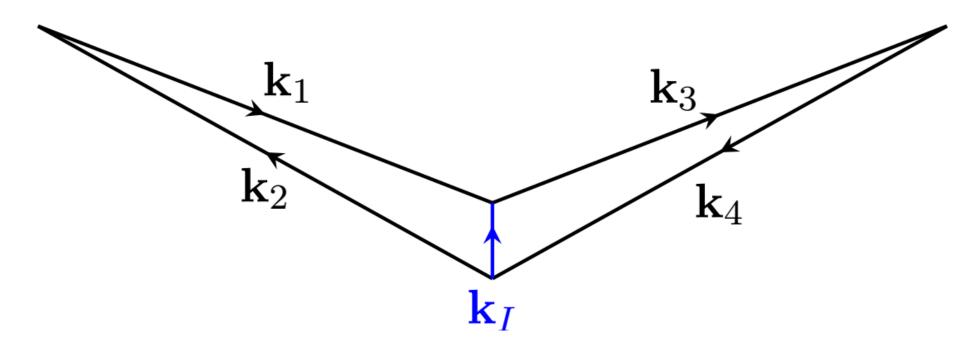
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Novel consistency relation similar to Suyama-Yamaguchi relation

$$eta_{NL}^{\zeta} = \left(\frac{\dot{\lambda}}{H\lambda}\right)^2 = \left(b_{NL}^{\zeta}\right)^2$$

T. Suyama and M. Yamaguchi, Phys. Rev. D 77 (2008) 023505 [0709.2545]

- . The equality is a lower bound, in general  $\beta_{NL}^{\zeta} \geq \left(b_{NL}^{\zeta}\right)^{z}$ 
  - The presence of  $\gamma$  exchange validate this inequality

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- Correlating the azimuthal orientation of  $k_I$  with k
- Higher order angular functions
- In counter collinear limit: enhanced correlation along  $k_I$
- Equisided configuration: oscillatory angular dependence (spin -2 nature)

# Summary and conclusion

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- We did the systematic computation of four point autocorrelation function of U(1) gauge field with full in-in formalism upto cubic order interation with both tensor and scalar perturbation.
- From the auxiliary gauge field we also obtain the trispectrum for B and E, in the contracted form which contributes to energy density.
- In the equisided configurations the non-Gaussian signal is monotonically increasing as the shape approaches the flattened configurations.
- In the counter collinear limit, we showed that the magnetic field trispectrum follows a novel consistency relation.

 $\beta_{NL}^{\zeta} = \left(b_{NL}^{\zeta}\right)^2$ 

• Magnetic trispectrum sourced by tensor exchange, inherits a richer angular dependence leads to distinctive signatures. But suppressed by tensor to scale ratio

#### Thank You

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